

# Obstruction theory for parameterized higher WZW terms

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June 11, 2015

Talk at AMS-EMS-SPM meeting 2015 in Porto

<http://ncatlab.org/schreiber/show/Obstruction+theory+for+parameterized+higher+WZW+terms>



FCT Fundação para a Ciência e a Tecnologia  
MINISTÉRIO DA EDUCAÇÃO E CIÊNCIA

funded by UID/MAT/00297/2013

# Motivation

**Classical fact:** Obstruction to globalizing a form  $\omega \in \Omega_{\text{cl}}^{p+2}(\mathbb{R}^n)$  over an  $n$ -manifold is existence of  $\text{Stab}_{\text{GL}(n)}(\omega)$ -structure.

Example:  $\omega \in \Omega^3(\mathbb{R}^7)$  the associative 3-form  $\Rightarrow G_2$ -structure.

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**Questions:** What happens as...

- 1.) ...forms are prequantized to Deligne cocycles  $\mathbf{L}$ ?
  - 2.) ...base is allowed to be a higher étale stack?
  - A) What are the obstructions to existence of a globalization?
  - B) What is group stack of symmetries of any given globalization?
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**Application:**

All Green-Schwarz-type super  $p$ -brane sigma-models are controlled by  $\mathbf{L}_{\text{WZW}}^{p+2}$  globalized over a super-spacetime;

for the D-branes and for the M5-brane the base is a higher stack modeled on the homotopy fiber of  $\mathbf{L}_{\text{WZW}}^{\text{F1}}$ ,  $\mathbf{L}_{\text{WZW}}^{\text{M2}}$ , respectively.

- A) Obstruction to global existence:  
classical anomalies (was completely open)
- B) Symmetries of given globalization:  
BPS charge extended superisometries (was only known rationally)

## Blueprint: ordinary geometric prequantization

Consider  $\omega := dp_i \wedge dq^i \in \Omega_{\text{cl}}^2(\mathbb{R}^{2n})$  and  $\mathbf{L} := p_i \wedge dq^i$ , then:

- |   |                            |
|---|----------------------------|
| 1. definite globalization $(X, \omega^X)$ :     | alm. symplectic structure; |
| 2. definite globalization $(X; \mathbf{L}^X)$ : | prequantum line bundle;    |
| 3. symmetry group of $(X, \omega^X)$ :          | symplectomorphism group;   |
| 4. symmetry group of $(X, \mathbf{L}^X)$ :      | quantomorphism group.      |
- 

Since  $\mathbf{L}$  has automorphisms, where  $\omega$  does not, quantomorphisms form (central) extension:

$$\bigoplus_{\pi_0(X)} (\mathbb{R}/\Gamma)^{\text{c}} \longrightarrow \text{QuantMorph}(X, \mathbf{L}^X) \longrightarrow \text{HamSympl}(X, \omega^X)$$

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On the Lie algebra level this is the Poisson bracket extension:

$$0 \rightarrow H^0(X) \rightarrow \text{pois}(X, \omega^X) \rightarrow \text{HamVect}(X, \omega^X) \rightarrow 0.$$

For  $(X, \omega^X)$  a symplectic vector space, this is the Heisenberg extension.

## Higher differential geometry

Lift classical theory from the category `SmoothMfd` to the homotopy theory  $\mathbf{H}$  of simplicial sheaves over smooth manifolds (“smooth  $\infty$ -groupoids”, “higher smooth stacks”).

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**Theorem** (classical+Lurie'12):  $A_\infty$ -group stacks are equivalently loopings of pointed connected higher stacks:

$$\mathrm{Grp}(\mathbf{H}) \begin{array}{c} \xleftarrow{\Omega} \\ \xrightarrow[\mathbf{B}]{\simeq} \end{array} \mathbf{H}_{\geq 1}^*$$

---

**Theorem** [dcct]:  $\mathbf{H}$  is *cohesive*, the derived global section coreflection  $\flat := \mathrm{Lconst} \circ \Gamma$  produces moduli stacks of flat  $G$ -principal connections for any  $A_\infty$ -group stack  $G$ .

The double homotopy fiber of the  $\flat$ -counit is the higher

*Maurer-Cartan form*:  $G \xrightarrow{\theta_G} \flat_{\mathrm{dR}} \mathbf{B}G$

$$\begin{array}{ccc} & \downarrow & \\ & \flat \mathbf{B}G & \longrightarrow \mathbf{B}G \end{array}$$

# Higher prequantization

The Dold-Kan correspondence  $DK : \text{Ch}_{\bullet \geq 0} \xrightarrow{\simeq} \text{Ab}^{\Delta^{\text{op}}} \rightarrow \text{Set}^{\Delta^{\text{op}}}$  includes traditional sheaf hypercohomology into  $\mathbf{H}$ .

Write:  $\mathbf{B}^{p+1}(\mathbb{R}/\Gamma) := DK((\underline{\mathbb{R}}/\Gamma)[p+1]);$   
 $b_{\text{dR}} \mathbf{B}^{p+2} \mathbb{R} := DK(\Omega^1 \xrightarrow{d} \dots \xrightarrow{d} \Omega_{\text{cl}}^{p+2})$

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**Proposition** [dcct]: Homotopy pullback of MC-form  $\theta_{\mathbf{B}^{p+1}(\mathbb{R}/\Gamma)}$  along the global differential form inclusion

$$\begin{array}{ccc}
 \mathbf{B}^{p+1}(\mathbb{R}/\Gamma)_{\text{conn}} & \longrightarrow & \mathbf{B}^{p+1}(\mathbb{R}/\Gamma) \\
 \uparrow \mathbf{L} & & \downarrow \theta_{\mathbf{B}^{p+1}(\mathbb{R}/\Gamma)} \\
 X & \xrightarrow{\omega} & \Omega_{\text{cl}}^{p+2} \longrightarrow b_{\text{dR}} \mathbf{B}^{p+2} \mathbb{R} \\
 & & \downarrow F_{(-)}
 \end{array}$$

is given by the Deligne complex:

$$\mathbf{B}^{p+1}(\mathbb{R}/\Gamma)_{\text{conn}} \simeq DK[\Gamma \hookrightarrow \Omega^0 \xrightarrow{d} \Omega^1 \xrightarrow{d} \dots \xrightarrow{d} \Omega^{p+1}]$$


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**Fact:**  $X \xrightarrow{\mathbf{L}} \mathbf{B}^{p+1}(\mathbb{R}/\Gamma)_{\text{conn}}$  has holonomy over closed mfd.  $\Sigma$ :

$$[\Sigma, X] \xrightarrow{[\Sigma, \mathbf{L}]} [\Sigma, \mathbf{B}^{p+1}(\mathbb{R}/\Gamma)_{\text{conn}}] \xrightarrow{\text{fiber integration}} \mathbf{B}^{p+1-\dim(\Sigma)}(\mathbb{R}/\Gamma)_{\text{conn}}$$

$\xrightarrow{\text{WZW action functional}}$

## Infinitesimal symmetries

**Theorem** [FRS13b]: For  $X \in \text{SmthMfd}$ , infinitesimal symmetries of  $(X, \mathbf{L})$  form  $L_\infty$ -algebra extension of vector fields by the abelian  $L_\infty$ -algebra on the de Rham complex, classified by the  $L_\infty$ -cocycle given by  $\iota_{(-)}\omega$ : there is a homotopy fiber sequence

$$\begin{array}{ccc} \Omega^\bullet[\rho] & \longrightarrow & \text{poiss}(X, \mathbf{L}) \\ & & \downarrow \\ & & \text{HamVect}(X, \omega) \xrightarrow{\iota_{(-)}\omega} \Omega^\bullet[\rho + 1] \end{array}$$

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**Corollary.** Under 0-truncation  $\tau_0$  (chain homology) this becomes a central extension of Lie algebras

$$0 \longrightarrow H_{\text{dR}}^p(X) \longrightarrow \tau_0 \text{poiss}(X, \mathbf{L}) \longrightarrow \text{HamVect}(X, \omega) \rightarrow 0.$$

Regarding  $\mathbf{L}$  as a WZW term, then  $\tau_0 \text{poiss}(X, \mathbf{L})$  is the Dickey bracket on Noether currents<sup>1</sup> for target space symmetries of the sigma-model with WZW term  $\mathbf{L}$  (compare AGIT'89).

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<sup>1</sup>With Igor Khavkine.

# Finite symmetries

## Definition:

- 1.)  $\text{conc} : [X, \mathbf{B}^{p+1}(\mathbb{R}/\Gamma)_{\text{conn}}] \longrightarrow (\mathbf{B}^p(\mathbb{R}/\Gamma))\text{Conn}(X)$   
projection on moduli of vertical differential forms;
  - 2.)  $\text{QuantMorph}(X, \mathbf{L}) := \text{Stab}_{\text{Aut}(X)}(\text{conc}(\mathbf{L}))$   
homotopy stabilizer group;
  - 3.)  $\text{HamSymp}(X, \mathbf{L}) := \text{im}_1(\text{QuantMorph}(X, \mathbf{L}) \rightarrow \text{Aut}(X))$   
1-image of quantomorphisms in automorphisms;
  - 4.)  $\text{Heis}_G(X, \mathbf{L}) := \rho^* \text{QuantMorph}(X, \mathbf{L})$   
its pullback along any  $G$ -action  $\rho : G \rightarrow \text{Aut}(X)$ .
- 

**Theorem** [FRS13a]: There is a homotopy fiber sequence:

$$\begin{array}{ccc} (\mathbf{B}^{p-1}(\mathbb{R}/\Gamma))\text{FlatConn}(X) & \rightarrow & \text{QuantMorph}(X, \mathbf{L}) \\ & & \downarrow \\ & & \text{HamSymp}(X, \mathbf{L}) \xrightarrow{\text{KS}} \mathbf{B}((\mathbf{B}^{p-1}(\mathbb{R}/\Gamma))\text{FlatConn}(X)) \end{array}$$

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**Example:** For  $p = 0$  this reduces to the traditional Heisenberg-Kostant-Souriau quantomorphism group extension.

## Obstruction theory, part I

**Theorem** [dcct]:

Let  $G$  be an  $A_\infty$ -group stack  
and  $\mathbf{L} : G \rightarrow \mathbf{B}^{p+1}(\mathbb{R}/\Gamma)_{\text{conn}}$   
with  $G \hookrightarrow \mathbf{HamSymp}(X, \mathbf{L})$ .

Then the obstruction to a definite parameterization of  $\mathbf{L}$  over the fibers of a  $G$ -principal  $\infty$ -bundle  $P \rightarrow X$  [NSS12a] is equivalently:

- 1.) a lift of the structure group through  $\mathbf{Heis}_G(X, \mathbf{L}) \rightarrow G$ ;
  - 2.) a trivialization of  $\mathbf{KS}(P)$ .
- 

**Example** [FRS13a]: For  $G = \text{Spin}$  and  $\mathbf{L}_{\text{WZW}}^{\langle -, [-, -] \rangle}$  the traditional WZW term, then

$$\mathbf{Heis}_{\text{Spin}}(\text{Spin}, \mathbf{L}_{\text{WZW}}^{\langle -, [-, -] \rangle}) \simeq \text{String}$$

is the smooth String-2-group, and hence a definite parameterization here is precisely a smooth String-structure.

This gives the geometric interpretation of the cancellation of the Green-Schwarz anomaly due to Distler-Sharpe'07 .



## Higher WZW terms

For  $\mathfrak{g}$  an  $L_\infty$ -algebra and  $\mu_{p+2} \in \text{CE}^{p+2}(\mathfrak{g})$  an  $L_\infty$ -cocycle, write  $\mathbf{B}G \in \mathbf{H}$  for the  $(p+2)$ -coskeleton of the simplicial sheaf of flat  $\mathfrak{g}$ -valued forms on simplices, parameterized by manifolds  $U$ :

$$\mathbf{B}G : (U \in \text{SmoothMfd}) \mapsto \text{cosk}_{p+2}(\Omega_{\text{flat}}^{\bullet}(\underset{\text{vert}}{U} \times \Delta_{\text{smth}}^{\bullet}, \mathfrak{g}))$$

**Theorem [FSS10]:**  $\mu$  Lie integrates to  $\mathbf{c} : \mathbf{B}G \rightarrow \mathbf{B}^{p+2}(\mathbb{R}/\Gamma)$

---

**Definition.** Denote by

$$\begin{array}{ccc} \tilde{G} & \longrightarrow & G \\ \theta_{\tilde{G}} \downarrow & \text{(pb)} & \downarrow \theta_G \\ \Omega_{\text{flat}}(-, \mathfrak{g}) & \longrightarrow & \mathfrak{b}_{\text{dR}} \mathbf{B}G \end{array}$$

the homotopy

pullback of the MC form on  $G$  to globally defined forms.

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**Theorem [dcct]:**  $\Omega \mathbf{c}$  underlies a unique prequantization

$$\mathbf{L}_{\text{WZW}}^{\mu} : \tilde{G} \rightarrow \mathbf{B}^{p+1}(\mathbb{R}/\Gamma)_{\text{conn}}$$

of  $\mu(\theta_{\tilde{G}}) \in \Omega^{p+2}(\tilde{G})$ .

This is the higher WZW term of  $\mu$ .

## Higher étale stacks (higher orbifolds)

Let now  $\mathbf{H}$  be simplicial sheaves over *formal* manifolds.

**Theorem** [dcct]: This is *differentially cohesive*, reduction  $\mathfrak{R}$  of infinitesimals has a right adjoint  $\mathfrak{S}$  (“de Rham stack functor”).

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**Definition.** A morphism is infinitesimally étale  $\dashrightarrow_{\text{et}}$  if its  $\mathfrak{S}$ -unit is a homotopy pullback square.

**Definition.** For  $V$  an  $A_\infty$ -group stack, a  $V$ -étale stack is an  $X \in \mathbf{H}$  such that there exists a  $V$ -cover:  $V \leftarrow_{\text{et}} U \dashrightarrow_{\text{et}} X$ .

**Definition.** The *infinitesimal disk bundle* is:

$$\begin{array}{ccc} T_{\text{inf}}X & \xrightarrow{\text{ev}} & X \\ \downarrow p & \text{(pb)} & \downarrow \\ X & \longrightarrow & \mathfrak{S}X \end{array}$$

---

**Theorem** [dcct]: 1.) The infinitesimal disk bundle of  $V$  trivializes via left translation, with typical fiber the infinitesimal disk  $\mathbb{D}_e^V$ ; 2.) the infinitesimal disk bundle of any  $V$ -étale stack is locally trivial and associated to a  $\text{GL}(V) := \mathbf{Aut}(\mathbb{D}_e^V)$ -principal  $\infty$ -bundle: the *frame bundle*  $\text{Fr}(X) \rightarrow X$ .

## Obstruction theory, part II

**Definition:** Given  $\mathbf{L} : V \rightarrow \mathbf{B}^{p+1}(\mathbb{R}/\Gamma)_{\text{conn}}$ , a *definite globalization* over a  $V$ -étale stack  $X$  is  $\mathbf{L}^X : X \rightarrow \mathbf{B}^{p+1}(\mathbb{R}/\Gamma)_{\text{conn}}$  such that its restriction to infinitesimal disks along  $T_{\text{inf}}X \xrightarrow{\text{ev}} X$  is a parameterization (as above) definite on  $\mathbf{L}|_{\mathbb{D}_e^V}$ .

---

**Corollary [dcct]:** An obstruction to definite globalization is  $\mathbf{Heis}_{\text{GL}(V)}(\mathbb{D}_e^V, \mathbf{L}|_{\mathbb{D}_e^V})$ -structure, hence trivialization of the **KS**-class of the frame bundle.

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**Example:** For  $X$  an  $\mathbb{R}^{2n}$ -manifold and  $\mathbf{L} = p^i \wedge dq^i$ , and for second order infinitesimals, then  $\mathbf{Heis}_{\text{GL}(V)}(\mathbf{L}|_{\mathbb{D}_e^V}) \simeq \text{Mp}^c(2n)$  is the Metaplectic<sup>c</sup> group.

**Example [dcct]<sup>2</sup>:** For  $V$  super-Minkowski spacetime and  $\mathbf{L}$  the  $\kappa$ -WZW term for the GS super-string, then  $\mathbf{Heis}_{\text{Aut}_{\text{grp}}(\mathbb{D}_e^V)}(\mathbf{L}|_{\mathbb{D}_e^V})$  is a  $\mathbf{B}(\mathbb{R}/\Gamma)$ -extension of the Lorentzian Spin group  $\text{Spin}(d-1, 1)$ .

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<sup>2</sup>With John Huerta.

# Higher supergeometry

Let now  $\mathbf{H}$  be  
simplicial sheaves  
over formal  
*supermanifolds*.

**Theorem** [dcct]<sup>3</sup>:

$\exists$  Progression of adjoint  
(co-)localizations:  $\rightarrow$   
satisfying  $\overset{\sim}{\mathfrak{S}} \simeq \mathfrak{S}$ .

**Proposition** [dcct]:

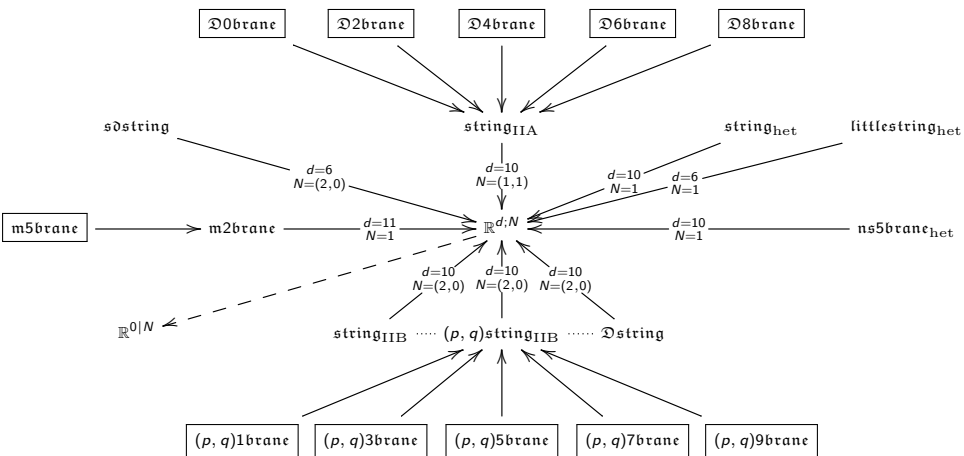
For  
 $X$  a  $V$ -étale stack,  
then  
 $\overset{\sim}{X}$  is  $\overset{\sim}{V}$ -étale stack.

	id	$\dashv$	id
	$\vee$		$\vee$
even	$\rightrightarrows$	$\dashv$	$\rightsquigarrow$
	$\perp$		$\perp$
bosonic	$\rightsquigarrow$	$\dashv$	$\text{loc}_{\mathbb{R}^{0 1}}$
	$\vee$		$\vee$
reduced	$\mathfrak{R}$	$\dashv$	$\mathfrak{S}$
	$\perp$		$\perp$
étale	$\mathfrak{S}$	$\dashv$	$\&$
	$\vee$		$\vee$
shape	$\text{loc}_{\mathbb{R}}$	$\dashv$	$b$
	$\perp$		$\perp$
flat	$b$	$\dashv$	$\#$
	$\vee$		$\vee$
	$\emptyset$	$\dashv$	$*$

<sup>3</sup>With Dave Carchedi.

# Application: GS-WZW terms for super $p$ -branes

**Fact** Azcárraga-Townsend'89+[FSS13]: The iterative super- $L_\infty$   $b^p\mathbb{R}$ -extensions of the superpoint come from the WZW cocycles  $\mu$  of all the Green-Schwarz-type super- $p$ -branes sigma-models:



## M5-brane on M2-brane extended superspacetime I

Regard super-Minkowski spacetime  $\mathbb{R}^{d-1,1|N}$  as a super Lie algebra. Write

$$\mu_{p+2} := \overline{\psi} \Gamma^{a_1 \cdots a_p} \wedge \psi \wedge e_{a_1} \wedge \cdots \wedge e_{a_p} \in \text{CE}(\mathbb{R}^{d-1,1|N}).$$

---

**Proposition** D'Auria-Fré 89: The elements  $\mu_4, \mu_7 \in \text{CE}(\mathbb{R}^{10,1|32})$  satisfy

$$d\mu_4 = 0, \quad d\mu_7 = \mu_4 \wedge \mu_4.$$

**Proposition** [FSS13]: The M2-brane extended super Minkowski spacetime with  $\text{CE}(\hat{\mathbb{R}}^{10,1|32}) := \text{CE}((\mathbb{R}^{10,1|32}) \otimes \langle h_3 \rangle, dh_3 = -\mu_4)$  is the  $L_\infty$ -homotopy fiber of  $\mu_4$  and we have

$$\begin{array}{ccc} \hat{\mathbb{R}}^{10,1|32} & \xrightarrow{h_3 \wedge \mu_4 + \frac{1}{15} \mu_7} & b^6 \mathbb{R} \\ \downarrow & & \\ \mathbb{R}^{10,1|32} & \xrightarrow{\mu_4} & b^3 \mathbb{R} \end{array}$$

matching the proposal in BLNPST'97.

## M5-brane on M2-brane extended superspacetime II

**Theorem** [dcct]: For a definite globalization of consecutive WZW terms such as

$$\begin{array}{ccc} \tilde{X} & \xrightarrow{\mathbf{L}_{\text{WZW}}^{\text{M5}}} & \mathbf{B}^6 U(1)_{\text{conn}} \\ \downarrow & & \\ X & \xrightarrow{\mathbf{L}_{\text{WZW}}^{\text{M2}}} & \mathbf{B}^3 U(1)_{\text{conn}} \end{array}$$

$\tilde{X}$  is a  $\mathbf{B}^2(\mathbb{R}/\Gamma)_{\text{conn}}$ -bundle over  $X$ .

---

Hence a map  $\Sigma \rightarrow \tilde{X}$  is a pair consisting of

1. a sigma-model field  $\phi : \Sigma \rightarrow X$ ;
2. a  $\phi$ -twisted degree-3 Deligne cocycle (twisted 2-gerbe with connection) on  $\Sigma$ .

This is the “tensor multiplet” field content of the M5-brane globalized to a twisted 2-gerbe connection.

## M5-brane on M2-brane extended superspacetime III

**Proposition** above results+ Candiello-Lechner'93: First order integrable definite globalization of  $\mathbf{L}_{WZW}^{M2}$  implies super-Lorentzian structure with vanishing supertorsion, this in turn implies the vacuum equations of motion of 11d Einstein gravity., enhances them by cancelling the obstruction to making the M2-brane and M5-brane WZW terms be globally defined.

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**Proposition** [dcct]:

The isometry action on  $X$  lifts to an  $\infty$ -action on  $\tilde{X}$ .

**Corollary**: The infinitesimal symmetries of  $\mathbf{L}_{WZW}^{M5}$  are an extension of superisometries of spacetime by  $H^5$  of the  $K(\mathbb{Z}, 3)$ -bundle underlying  $\mathbf{L}_{WZW}^{M2}$ . Running the Serre spectral sequence, rationally this is  $H^2(X) \oplus H^5(X)$ .

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This is the traditional result for the *M-theory super Lie algebra*, the extension of the superisometries by BPS charges for the M2-brane and the M5-brane Sorokin-Townsend'97. The above analysis gives the finite global symmetries involving various (torsion) corrections to this.



# M5-brane on M2-brane extended superspacetime - Outlook

**Proposition:** M5-cocycle descends equivariantly down to super-Minkowski spacetime

$$\begin{array}{ccc}
 \vdots & & \vdots \\
 \Downarrow & & \Downarrow \\
 \widehat{\mathbb{R}}^{10,1|32} & \xrightarrow{h_3 \wedge g_4 + \frac{1}{15} \mu_7} & b^6 \mathbb{R} \\
 (dh_3 = -\mu_4) & & \\
 \downarrow & & \downarrow \\
 \mathbb{R}^{10,1|32} & \xrightarrow{h_3 \wedge (g_4 + \mu_4) + \frac{1}{15} \mu_7} & b^6 \mathbb{R} / b^2 \mathbb{R} \\
 (dh_3 = g_4 - \mu_4) & & \\
 \searrow \mu_4 & & \swarrow p_\rho \\
 & b^3 \mathbb{R} &
 \end{array}$$

By [NSS12a] this rational 4-sphere valued cocycle is in degree-7 twisted cohomology, the twist being the degree-4 class of the supergravity C-field.

This structure of the M-theory C-field was conjectured in Sati'13.

# M5-brane on M2-brane extended superspacetime - Outlook

All cocycles here are Spin-invariant. Hence we may ask for extending them from super-Minkowski to super-Poincaré  $\mathfrak{iso}(\mathbb{R}^{10,1|32})$ .

Such extensions are given by shifting  $h_3$  by an  $\mathfrak{so}$ -3-cocycle and  $\mu_7$  by an  $\mathfrak{so}$ -7-cocycle. The only such are  $\propto \langle \omega^{\wedge 3} \rangle, \langle \omega^{\wedge 7} \rangle$ :

$$\begin{array}{ccc} \mathfrak{iso}(\mathbb{R}^{10,1|32}) & & \\ \uparrow & \searrow & \\ \mathbb{R}^{10,1|32} & \xrightarrow{h_3 \wedge (g_4 + \mu_4) + \frac{1}{15} \mu_7} & b^6 \mathbb{R} / b^2 \mathbb{R} \\ & \nearrow (h_3 + \langle \omega^{\wedge 3} \rangle) \wedge (g_4 + \mu_4) + \frac{1}{15} \mu_7 + \langle \omega^{\wedge 7} \rangle & \end{array}$$

This is then to be globalized not just over  $X$ , but over the frame bundle  $\text{Fr}(X)$ . By the above obstruction theory, the parameterization of the WZW terms for  $\langle \omega^{\wedge 3} \rangle$  and  $\langle \omega^{\wedge 7} \rangle$  over the Frame bundle imposes String-structure and Fivebrane structure (cancelling the  $l_8$ -one loop term).

## Conclusion

- ▶ There is good general abstract theory for prequantized definite globalizations of higher degree forms over higher étale stacks.
- ▶ Higher Lie theory provides prequantization of every  $L_\infty$ -cocycle to a higher WZW term.
- ▶ Applying this to the bouquet of cocycles emanating from the superpoint yields super-orbifolds equipped with Lorentzian structure solving the vacuum Einstein equations of 11-dimensional supergravity and equipped with the classical anomaly cancellation that makes the M2-brane and M5-brane sigma models globally well-defined. The group stack of symmetries of these structures encodes various torsion corrections to the BPS charge extension of the superisometries.

The restriction to *vacuum* solutions (vanishing gravitino and C-field strength) could be circumvented by intervening by hand, but it is maybe noteworthy that these are the solutions relevant for realistic phenomenology (e.g. [Acharya'02](#), [Acharya'12](#)).

# Thank you!

For more details see course notes at

[ncatlab.org/schreiber/show/Structure+Theory+for+Higher+WZW+Terms](https://ncatlab.org/schreiber/show/Structure+Theory+for+Higher+WZW+Terms)



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