Prequantum field theory

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exposition, details and references at
ncatlab.org/schreiber/show/Local+prequantum+field+theory

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Let Σ be a $(p+1)$-dimensional smooth manifold

Given a smooth bundle E over Σ,
think of its sections $\phi \in \Gamma_\Sigma(E)$ as physical fields.

Task: Describe field theory **fully-local to fully-global**.

1. Classical field theory
2. BV-BFV field theory
3. Prequantum field theory
4. Prequantum ∞-CS theories

1)
Classical field theory

A *differential operator* $D : \Gamma_\Sigma(E) \to \Gamma_\Sigma(F)$ comes from bundle map out of the jet bundle

$$\tilde{D} : J^\infty_\Sigma E \longrightarrow F$$

Composition of diff ops $D_2 \circ D_1$ comes from

$$J^\infty_\Sigma E \longrightarrow J^\infty_\Sigma J^\infty_\Sigma E \xrightarrow{J^\infty_\Sigma \tilde{D}_1} J^\infty_\Sigma F \xrightarrow{\tilde{D}_2} G$$

here the first map witnesses comonad structure on jets, this is composition in the coKleisli category $\mathrm{Kl}(J^\infty_\Sigma)$.

A differential equation is an equalizer of two differential operators.
Exists in the Eilenberg-Moore category $\mathrm{EM}(J^{\infty}_{\Sigma})$

$$\mathcal{E} \longrightarrow E \underset{\tilde{D}_2}{\overset{\tilde{D}_1}{\rightrightarrows}} F$$

Theorem [Marvan 86]:

$$\mathrm{EM}(J^{\infty}_{\Sigma} E) \simeq \mathrm{PDE}_{\Sigma}$$

PDE solutions are sections:

$$\begin{array}{ccc} & \mathcal{E} & \\ {}^{\phi_{\mathrm{sol}}}\nearrow & & \searrow \\ \Sigma & \xrightarrow[\phi]{} & E \end{array}$$

Def. A *horizontal differential form* on jet bundle $\alpha \in \Omega^k_H(J^\infty_\Sigma E)$ is diff op of the form

$$\tilde{\alpha} : E \longrightarrow \wedge^k T^*\Sigma$$

This induces vertical/horizontal bigrading $\Omega^{\bullet,\bullet}(J^\infty_\Sigma E)$

Def: A *local Lagrangian* is $L \in \Omega^{p+1}_H(J^\infty_\Sigma E)$

Prop. Unique decomposition

$$d_{\mathrm{dR}} L = \mathrm{EL} - d_H(\Theta + d_H(\cdots))$$

with $\mathrm{EL} \in \Omega^{p+1,1}_S \hookrightarrow \Omega^{p+1,1}$
depending only on vector fields along 0-jets.

This is the local incarnation of the variational principle.

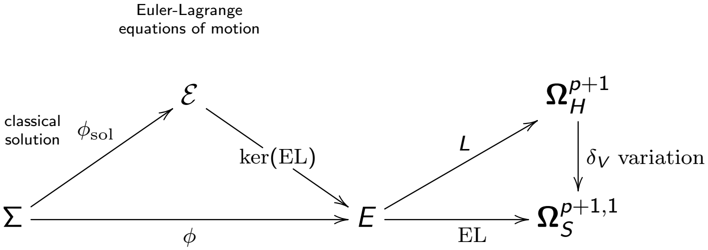
Prop. the bigrading is preserved by pullback along diff ops

This means that

$$\mathbf{\Omega}^{\bullet,\bullet} \in \mathrm{Sh}(\mathrm{DiffOp}_{\mathbf{\Sigma}}) \xrightarrow{\text{Kan ext}} \mathrm{Sh}(\mathrm{PDE}_{\mathbf{\Sigma}})$$

is a bicomplex of sheaves on the category of PDEs.

In this sheaf topos, classical field theory looks like so:



By *transgression* of this local data we get

1) the **global action functional**

$$[\Sigma, E]_\Sigma \xrightarrow{S = \int_\Sigma L} \mathbb{R}$$

2) the **covariant phase space** [Zuckerman 87]:

$$\begin{array}{ccc} & & \mathbf{\Omega}^1 \\ & \nearrow^{\theta := \int_{\Sigma_p} \Theta} & \downarrow^{d} \\ [N^\infty \Sigma_p, \mathcal{E}]_\Sigma & \xrightarrow[\omega]{} & \mathbf{\Omega}^2 \end{array}$$

2)
BV-BFV field theory

$\mathcal{E}$ need not be representable by a sub*manifold* of $J^\infty_\Sigma E$
if EL-equation is singular.

idea of BV-theory/derived geometry:

1. replace base category of smooth manifolds
   by smooth dg-manifolds in non-positive degree.
2. resolve singular $\mathcal{E}$ by realizing it
   as 0-cohomology of smooth dg-manifold.

If $\{\Phi^i\}$ are local fiber coordinates on E (field coordinates), then the derived shell $\mathcal{E}$ has dg-algebra of functions the algebra $C^\infty(\mathcal{E})$ with degree-(-1) generators Φ_i^* added ("antifields") and with differential given by

$$d_{BV} : \Phi_i^* \mapsto \mathrm{EL}_i \,.$$

Usual to write Q for d_{BV} regarded as vector field on the dg-manifold. Then this is

$$\mathcal{L}_Q \Phi_i^* = \mathrm{EL}_i \,.$$

This gives a third grading (BV antifield grading) on differential forms on the jet bundle

$$\Omega^{\bullet,\bullet;-\bullet}(\mathcal{E}_d)$$

Observation: $\mathcal{E}_d$ carries 2-form locally given by

$$\Omega_{\mathrm{BV}} = d\Phi^*_i \wedge d\Phi^i \quad \in \quad \Omega^{p+1,2;-1}(\mathcal{E}_d)$$

which satisfies

$$\begin{aligned}\iota_Q \Omega_{\mathrm{BV}} &= \mathrm{EL} \\ &= dL + d_H \Theta \qquad \in \Omega^{p+1,1;0}(\mathcal{E}_d)\,.\end{aligned}$$

under transgression to the space of fields this becomes

$$\iota_Q \omega_{\mathrm{BV}} = dS + \pi^* \theta$$

This is the central compatibility postulate for BV-BFV field theory in [Cattaneo-Mnev-Reshetikhin 12, eq. (7)]

3)
Prequantum field theory

For many field theories of interest, L is not in fact globally defined.

Simple examples:

- electron in EM-field with non-trivial magentic charge;
- 3d $U(1)$-Chern-Simons theory;

Large classes of examples:

- higher **WZW-type** models
  (super p-branes, topological phases of matter)
- higher **Chern-Simons-type** models
  (AKSZ, 7dCS on String-2-connections, 11dCS,...)

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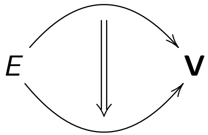
Claim: There is a *systematic* solution to this problem by

1. passing to the derived category of $\mathrm{Sh}(\mathrm{PDE}_\Sigma)$;
2. generalizing Lagrangian forms to differential cocycles ("gerbes with connection")

If
$$\mathbf{V} := [\cdots \stackrel{\partial_V}{\to} \mathbf{V}^2 \stackrel{\partial_V}{\to} \mathbf{V}^1 \stackrel{\partial_V}{\to} \mathbf{V}^0]$$
is a sheaf of chain complexes, then a map in the derived category

$$E \longrightarrow \mathbf{V}$$

is equivalently a cocycle
in the sheaf hypercohomology of E with coefficients in $\mathbf{V}$ [Brow73].
Homotopy is coboundary:



hence generalize sheaf of horizontal forms $\mathbf{\Omega}_H^{p+1}$ to:

Def. The "variational Deligne complex"

$$\mathbf{B}_H^{p+1}(\mathbb{R}/\mathbb{Z})_{\mathrm{conn}} := [\mathbb{Z} \to \mathbf{\Omega}_H^0 \stackrel{d_H}{\to} \mathbf{\Omega}_H^1 \stackrel{d_H}{\to} \cdots \stackrel{d_H}{\to} \mathbf{\Omega}_H^{p+1}]$$

Def. A local prequantum Lagrangian is

$$\mathbf{L} : E \longrightarrow \mathbf{B}_H^{p+1}(\mathbb{R}/\mathbb{Z})_{\mathrm{conn}}$$

Consequence: Via fiber integration in differential cohomology $\mathbf{L}$ transgresses to globally well defined action functional

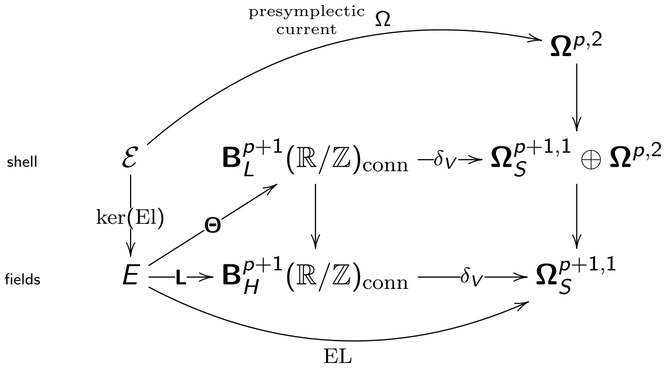
$$S := \int_\Sigma [\Sigma, \mathbf{L}] \; : \; [\Sigma, E]_\Sigma \longrightarrow \mathbb{R}/\mathbb{Z}$$

Theorem: the curving of such Euler-Lagrange p-gerbes is given by the Euler variational operator

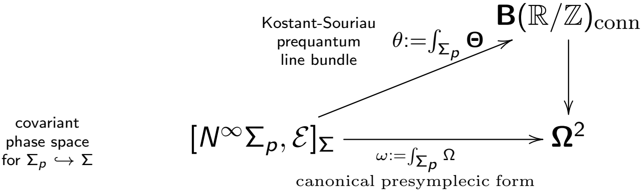
$$\delta_V : \mathbf{B}_H^{p+1}(\mathbb{R}/\mathbb{Z})_{\mathrm{conn}} \to \mathbf{\Omega}_S^{p+1,1}$$

principle of extremal action $\leftrightarrow$ flatness of EL-p-gerbes

Def./Prop. Prequantization of Θ is via *Lepage p-gerbes* $\mathbf{\Theta}$ whose curvature in degree $(p, 2)$
is the *pre-symplectic current*.



Theorem: Transgressions of single Lepage p-gerbe to all codimension-1 (Cauchy-)hypersurfaces $\Sigma_p \hookrightarrow \Sigma$ gives *natural* Kostant-Souriau prequantizations of all covariant phase spaces:



Theorem: Transgression of higher prequantum gerbes ∇ to fields on manifold Σ with boundary looks like so:

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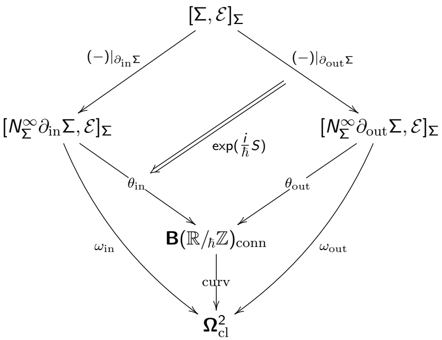
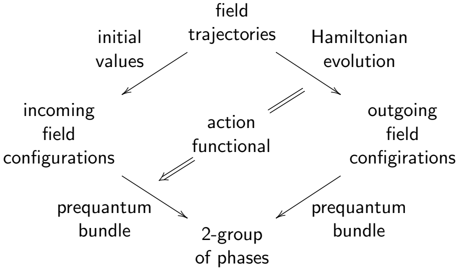
\begin{array}{ccccc}
[\Sigma, X] & \xrightarrow{[\Sigma,\nabla]} & [\Sigma, \mathbf{B}^{p+1}(\mathbb{R}/\hbar\mathbb{Z})_{\mathrm{conn}}] & \xrightarrow{\int_\Sigma \mathrm{curv}} & \mathbf{\Omega}^{p+2-d} \\
\downarrow{\scriptstyle (-)|_{\partial\Sigma}} & & \downarrow{\scriptstyle (-)|_{\partial\Sigma}} & \swarrow\!\!\swarrow{\scriptstyle \int_\Sigma} & \downarrow \\
[\partial\Sigma, X] & \xrightarrow[{[\partial\Sigma,\nabla]}]{} & [\partial\Sigma, \mathbf{B}^{p+1}(\mathbb{R}/\hbar\mathbb{Z})_{\mathrm{conn}}] & \xrightarrow[\int_{\partial\Sigma}]{} & \mathbf{B}^{p+2-d}(\mathbb{R}/\hbar\mathbb{Z})_{\mathrm{conn}}
\end{array}
$$

Corollary: Transgressing Lepage p-gerbe to spacetime Σ with incoming and outgoing boundary yields the prequantized Lagrangian correspondence that exhibits dynamical evolution:





4)
Prequantum ∞-CS theories

Chern-Simons is nonabelian gauge theory

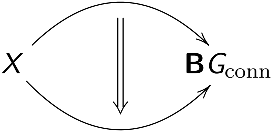
So now pass to the "nonabelian derived category" (aka ∞-topos)
where sheaves of chain complexes
are generalized to
sheaves of Kan complexes [Brow73].

This serves to describe nonabelian gauge fields.

For instance there is sheaf of Kan complexes $\mathbf{B}G_{\mathrm{conn}}$ such that a G-principal connection on X is a map

$$X \longrightarrow \mathbf{B}G_{\mathrm{conn}}$$

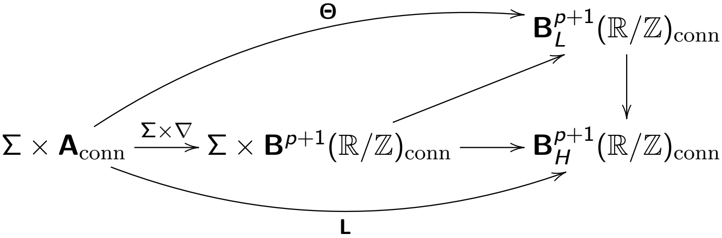
and a gauge transformation is a homotopy



One way to construct prequantum field theories is:
construct a p-gerbe connection on some moduli stack

$$\nabla : \mathbf{A}_{\mathrm{conn}} \longrightarrow \mathbf{B}^{p+1}(\mathbb{R}/\mathbb{Z})_{\mathrm{conn}}$$

and consider the stacky field bundle $E := \Sigma \times \mathbf{A}_{\mathrm{conn}}$
then pullback and project to get Euler-Lagrange and Lepage p-gerbe



For instance for G a simply-connected compact simple Lie group, there is a unique differential refinement

$$\nabla : \mathbf{B}G_{\mathrm{conn}} \longrightarrow \mathbf{B}^3 U(1)_{\mathrm{conn}}$$

of the canonical universal characteristic 4-class

$$c_2 : BG \longrightarrow K(\mathbb{Z}, 4)$$

This induces the standard 3d Chern-Simons Lagrangian and universally prequantizes it:

codim 0	$[\Sigma_3, \mathbf{B}G_{\mathrm{conn}}] \longrightarrow \mathbb{R}/\mathbb{Z}$	CS invariant
codim 1	$[\Sigma_2, \mathbf{B}G_{\mathrm{conn}}] \longrightarrow \mathbf{B}(\mathbb{R}/\mathbb{Z})_{\mathrm{conn}}$	CS prequantum line
codim 2	$[\Sigma_1, \mathbf{B}G_{\mathrm{conn}}] \longrightarrow \mathbf{B}^2(\mathbb{R}/\mathbb{Z})_{\mathrm{conn}}$	WZW gerbe
codim 3	$[\Sigma_0, \mathbf{B}G_{\mathrm{conn}}] \longrightarrow \mathbf{B}^3(\mathbb{R}/\mathbb{Z})_{\mathrm{conn}}$	Chern-Weil map

construct this and other examples from Lie integration of L_∞-data:

Def. L_∞-algebroid $\mathfrak{a}$ is dg-manifold

Chevalley-Eilenberg algebra $\mathrm{CE}(\mathfrak{a})$ is the dg-algebra of functions

$\mathrm{W}(\mathfrak{a})$ is dg-algebra of differential forms

cocycle is

$$\frac{\mathfrak{a} \xrightarrow{\mu} \mathbf{B}^{p+2}\mathbb{R}}{\mathrm{CE}(\mathfrak{a}) \longleftarrow \mathrm{CE}(\mathbf{B}^{p+2}\mathbb{R})}$$

Def. universal Lie integration to the derived category over $\mathrm{SmoothMfd}$ is the sheaf of Kan complexes

$$\exp(\mathfrak{a}) : (U, k) \mapsto \{\Omega^\bullet_{\mathrm{vert}}(U \times \Delta^k) \longleftarrow \mathrm{CE}(\mathfrak{a})\}$$

Examples:

- for $\mathfrak{g}$ semisimple Lie algebra, then $\tau_1 \exp(\mathbf{B}\mathfrak{g}) \simeq \mathbf{B}G$
- for $\mathfrak{P}$ a Poisson Lie algebroid then $\tau_2 \exp(\mathfrak{P})$ is symplectic Lie groupoid.
- for $\mathfrak{string}$ the String Lie 2-algebra, then $\tau_2 \exp(\mathbf{B}\mathfrak{string}) \simeq \mathbf{B}\mathrm{String}$

an *invariant polynomial* $\langle - \rangle$ on $\mathfrak{a}$ is closed differential form on the dg-manifold.
If $\langle - \rangle$ is binary and non-degenerate,
this became also called "shifted symplectic form"

Def. differential Lie integration $\exp(\mathfrak{g})_{\mathrm{conn}}$ is

$$\exp(\mathfrak{a})_{\mathrm{conn}} : (U, k) \mapsto \left\{ \begin{array}{ccc} \Omega^\bullet_{\mathrm{vert}}(U \times \Delta^k) & \longleftarrow & \mathrm{CE}(\mathfrak{a}) \\ \uparrow & & \uparrow \\ \Omega^\bullet(U \times \Delta^k) & \longleftarrow & \mathrm{W}(\mathfrak{a}) \\ \uparrow & & \uparrow \\ \Omega^\bullet(U) & \longleftarrow & \mathrm{inv}(\mathfrak{a}) \end{array} \right\}$$

Def. A cocycle μ is in transgression with an invariant polynomial $\langle - \rangle$ if there is a diagram of the form

$$\begin{array}{ccc} \mathrm{CE}(\mathfrak{a}) & \stackrel{\mu}{\longleftarrow} & \mathrm{CE}(\mathbf{B}^{p+2}\mathbb{R}) \\ \uparrow & & \uparrow \\ \mathrm{W}(\mathfrak{a}) & \stackrel{\mathrm{cs}}{\longleftarrow} & \mathrm{W}(\mathbf{B}^{p+2}\mathbb{R}) \\ \uparrow & & \uparrow \\ \mathrm{inv}(\mathfrak{a}) & \stackrel{\langle - \rangle}{\longleftarrow} & \mathrm{inv}(\mathbf{B}^{p+2}\mathbb{R}) \end{array}$$

By pasting of diagrams, this data defines a map

$$\exp(\mathrm{cs}) \ : \ \exp(\mathfrak{a})_{\mathrm{conn}} \longrightarrow \exp(\mathbf{B}^{p+2}\mathbb{R})_{\mathrm{conn}}$$

Theorem:
under truncation this descends to

$$\begin{array}{ccc} \exp(\mathfrak{a})_{\mathrm{conn}} & \xrightarrow{\exp(\mathrm{cs})} & \exp(\mathbf{B}^{p+2}\mathbb{R})_{\mathrm{conn}} \\ \downarrow{\scriptstyle \mathrm{cosk}_{p+2}} & & \downarrow \\ \mathbf{A}_{\mathrm{conn}} & \xrightarrow{\nabla} & \mathbf{B}^{p+2}(\mathbb{R}/\Gamma)_{\mathrm{conn}} \end{array}$$

where $\Gamma \hookrightarrow \mathbb{R}$ is the group of periods of μ.

This gives large supply of **examples** of
prequantum field theories induced from "shifted *n*-plectic forms":

AKSZ including 3dCS, PSM (hence A-model/B-model), CSM, ...
7dCS on String 2-connections, 11dCS on 5brane 6-connections, ...

Outlook

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Our formulation of prequantum field theory
works also with *generalized* differential cohomology.

For instance the “topological” Lagrangian term
for super Dp-brane sigma models
for all even or all odd p at once
needs to be a cocycle in differential K-theory.

And “U-duality” predicts
that the topological terms for the M2 and M5 brane
needs to be unified in single generalized differential cocycle
with Chern character in the rational 4-sphere
[Fiorenza-Sati-Schreiber 15].

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References

exposition, details and full bibliography is at

ncatlab.org/schreiber/show/Local+prequantum+field+theory

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