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Equivariant Cohomotopy and Branes

—

Towards microscopic M-Theory

talk at

String and M-Theory: The New Geometry of the 21st Century

Singapore 2018

joint work with

Hisham Sati & V. Braunack-Mayer

based on [FSS13, FSS15, FSS16a, FSS16b, HS17, HSS18, BSS18]

Motivation.

Nonperturbative QFT and an old Prophecy

Part I.

Some M-Theory from Super homotopy theory

Part II.

Some corners of M-theory

Motivation

[back to Contents](#)

Glaring open problem of contemporary quantum field theory:

All non-perturbative physics.

Such as:

- quark confinement in hadrons (existence of ordinary matter!)
- quark-gluon plasma & nucleosynthesis (becoming of ordinary matter!)
- Higgs field metastability (existence of vacuum spacetime!)
- QCD cosmology (becoming of vacuum spacetime!)

Important non-answers:

lattice QFT numerics is (great but) *not the answer*:

like Bohr-Sommerfeld's "old quantization"

it allows to compute some numbers

but without conceptual understanding

string theory is (great but) *not the answer*:

string scattering series just as perturbative as Feynman series

(vanishing radius of convergence, both)

But string theory is the vehicle with which to glimpse **M-Branes**...

The emerging answer: Intersecting M-branes

Web of plausibility arguments and consistency checks suggests:

*Non-perturbative standard model of particle physics & cosmology
arises on intersecting M-branes
at asymptotic boundary of approximately AdS spacetime.*

In particular

Witten-Sakai-Sugimoto model for QCD:

N_c M5-branes intersecting N_f M9-branes

KK-compactified, breaking all supersymmetry, to

N_c D4-branes intersecting N_f D8-branes

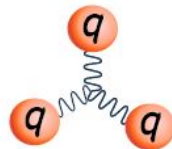
yields QFT at least close to non-perturbative QCD

with transparent interpretation of non-perturbative effects

meson



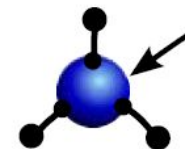
baryon



open string



D4 with N_c strings



D4-brane

Fig. 15.1. mesons and baryons in quark model and string theory

graphics from Sugimoto 16

However...

Glaring open problem of contemporary M-brane theory:

What is it, really?

We still have no fundamental formulation of “M-theory” -

Work on formulating the fundamental principles underlying M-theory has noticeably waned. [...]. If history is a good guide, then we should expect that anything as profound and far-reaching as a fully satisfactory formulation of M-theory is surely going to lead to new and novel mathematics. Regrettably, it is a problem the community seems to have put aside - temporarily. But, ultimately,

Physical Mathematics must return to this grand issue.

G. Moore, *Physical Mathematics and the Future*, at Strings 2014

Important non-answer:

BFSS/IKKT matrix model is (great but) *not the answer*:

like lattice QFT numerics

it allows to compute some numbers

but without conceptual understanding

What is missing?..

An old prophecy

Back in the '70s, the Italian physicist, D. Amati reportedly said that string theory was part of 21st-century physics that fell by chance into the 20th century. I think it was a very wise remark. How wise it was is so clear from the fact that 30 years later we're still trying to understand what string theory really is.

E. Witten, Nova Interview 2003

New development brought by the 21st century:

Homotopy theory & higher topos theory (“higher structures”)

physics	mathematics
gauge principle	homotopy theory
& Pauli exclusion	super-geometry
<hr/> <hr/>	
=	super homotopy theory

Part I.

Some M-theory from Super homotopy theory

1. Super homotopy theory and the Atom of Superspace
Rational
2. Super homotopy theory and the fundamental super p -Branes
Global equivariant
3. Super homotopy theory and the C -field at singularities
4. Super Cartan geometry and 11d orbifold supergravity

Super homotopy theory
and the Atom of Superspace

[back to Part I](#)

Global equivariant Super homotopy theory

Definition. Consider the 2-site

$$\text{SuperSingularities} := \left\{ \underbrace{\mathbb{R}^{d|N}}_{\text{super-space}} \times \underbrace{\mathbb{D}}_{\text{infinitesimal disk}} \times \underbrace{\mathbb{B}G}_{\text{orbifold singularity}} \right\}$$

Global equivariant super homotopy theory is the ∞ -stack ∞ -topos over SuperSingularities:

$$\mathbf{H} := \underbrace{\text{Sh}_\infty(\text{SuperSingularities})}_{\text{generalized spaces}} \quad \begin{array}{l} \nearrow \\ \text{probe-able} \\ \text{by these} \\ \text{local model spaces} \end{array} \quad \begin{array}{l} \text{[Sch13]} \\ \& \text{[Rezk14]} \\ \hline \text{[BMSS19]} \end{array}$$

space-time

SuperManifolds,
SuperOrbifolds
with G -structure

(e.g.

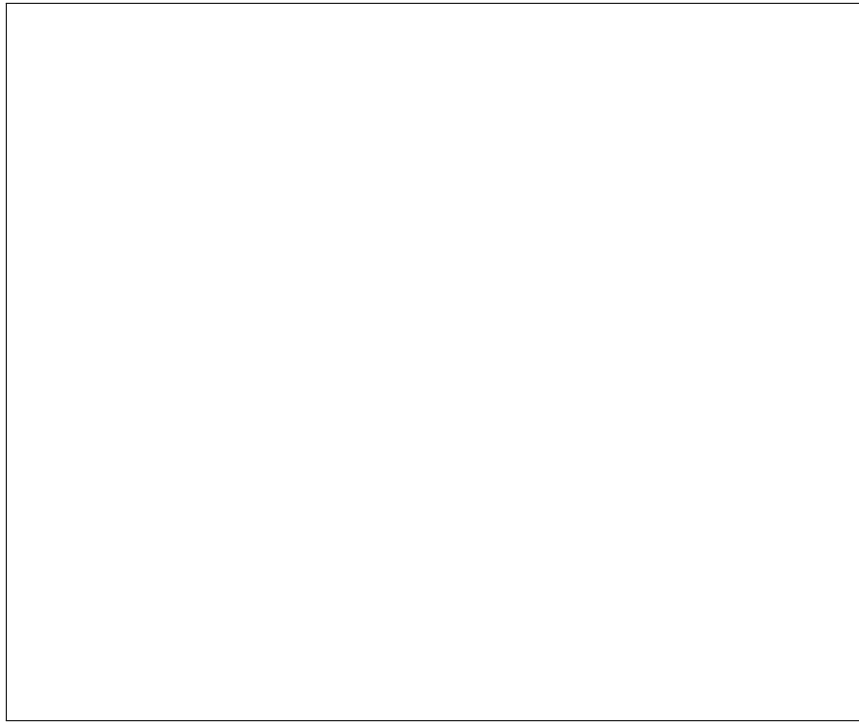
super-Riemannian-, super-conformal-, ...
Spin-, String-, Fivebrane-, ... ([SSS09]
structure)

$\longleftrightarrow \mathbf{H} \longleftrightarrow$

higher gauge fields

classifying spaces for
equivariant, differential
generalized cohomology theories

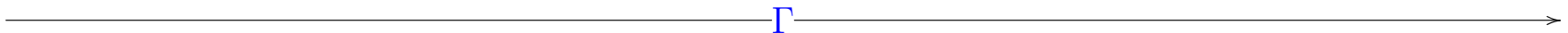
The Modalities of Super homotopy theory



The **terminal functor** factors into a system of dualities = adjunctions.

supergeometric
 ∞ -groupoids

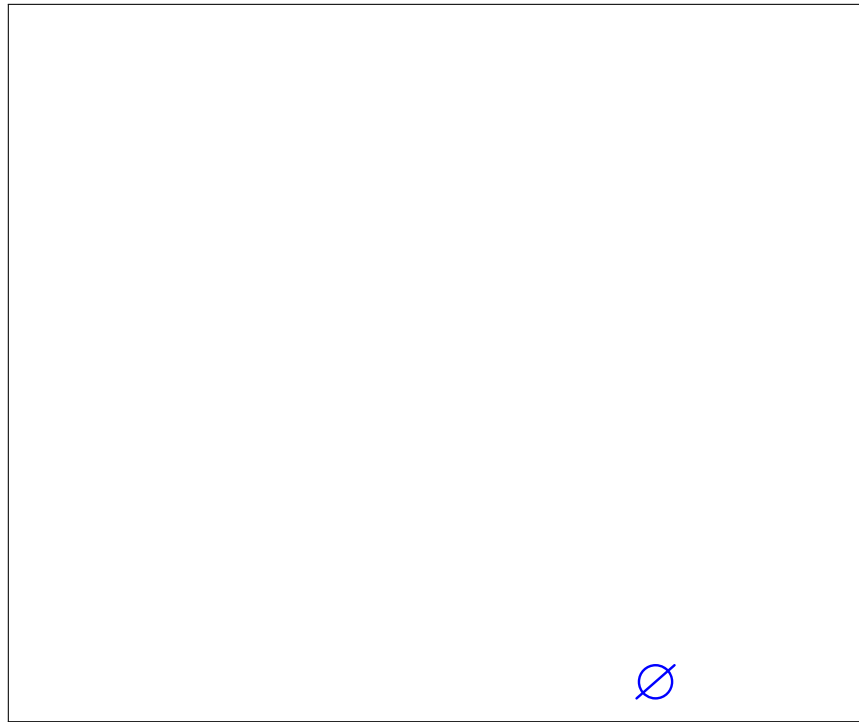
\mathcal{H}



Γ

*

The Modalities of Super homotopy theory

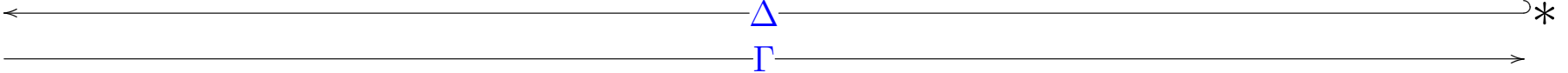


nothing

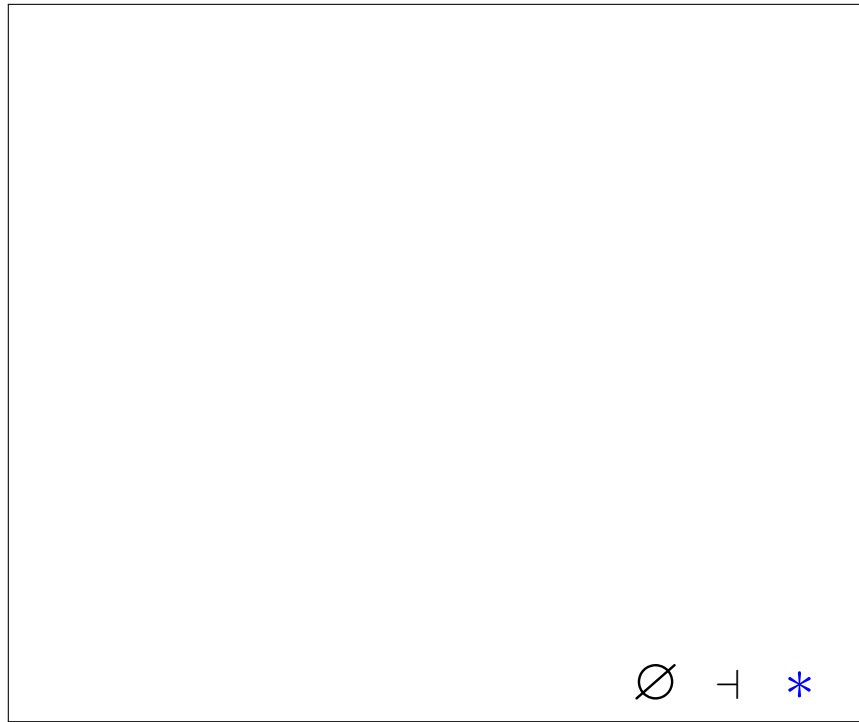
The terminal functor factors into a system of dualities = adjunctions.

supergeometric
 ∞ -groupoids

\mathbf{H}



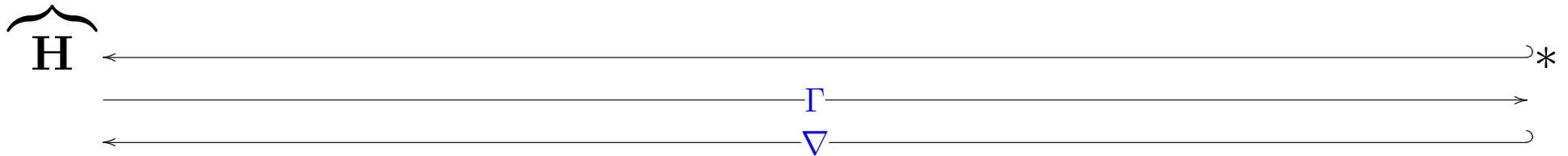
The Modalities of Super homotopy theory



pure being

The terminal functor factors into a system of dualities = adjunctions.

supergeometric
 ∞ -groupoids



The Modalities of Super homotopy theory



discrete

\mathfrak{b}
 \vee
 $\emptyset \dashv *$

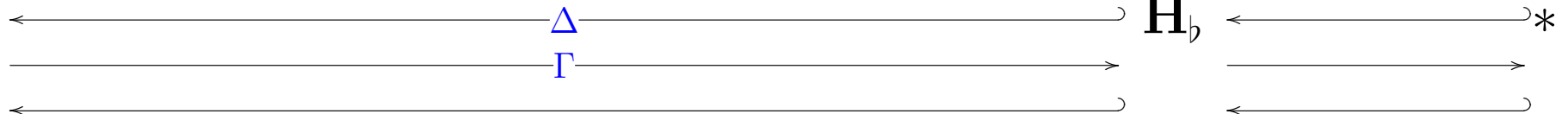
The terminal functor factors into a system of dualities = adjunctions.

supergeometric
 ∞ -groupoids

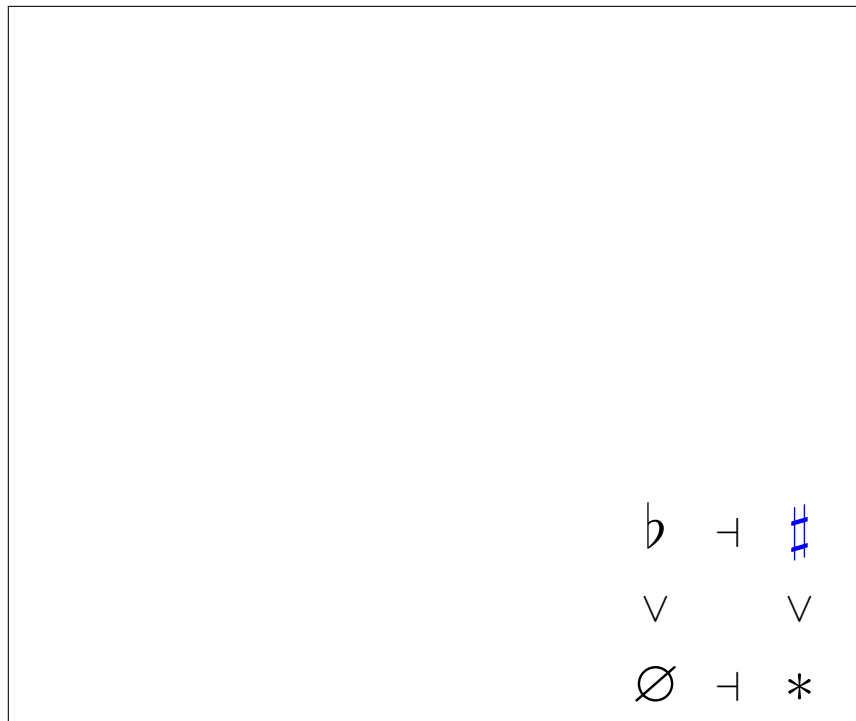
\mathfrak{H}

geometrically discrete
 ∞ -groupoids

$\mathfrak{H}_{\mathfrak{b}}$



The Modalities of Super homotopy theory



continuous

$$\begin{array}{ccc}
 \mathfrak{b} & \dashv & \# \\
 \vee & & \vee \\
 \emptyset & \dashv & *
 \end{array}$$

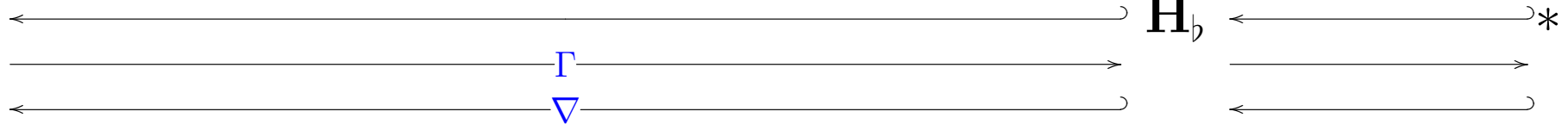
The terminal functor factors into a system of dualities = adjunctions.

supergeometric
 ∞ -groupoids

\mathfrak{H}

geometrically discrete
 ∞ -groupoids

$\mathfrak{H}_{\mathfrak{b}}$



The Modalities of Super homotopy theory



shaped

$$\begin{array}{cccc}
 \int & \dashv & \flat & \dashv & \sharp \\
 & & \vee & & \vee \\
 & & \emptyset & \dashv & *
 \end{array}$$

The terminal functor factors into a system of dualities = adjunctions.

supergeometric
 ∞ -groupoids

\mathbf{H}

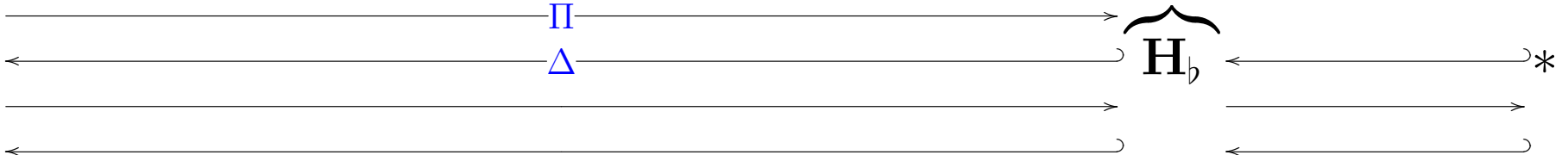
geometrically discrete
 ∞ -groupoids

\mathbf{H}_\flat

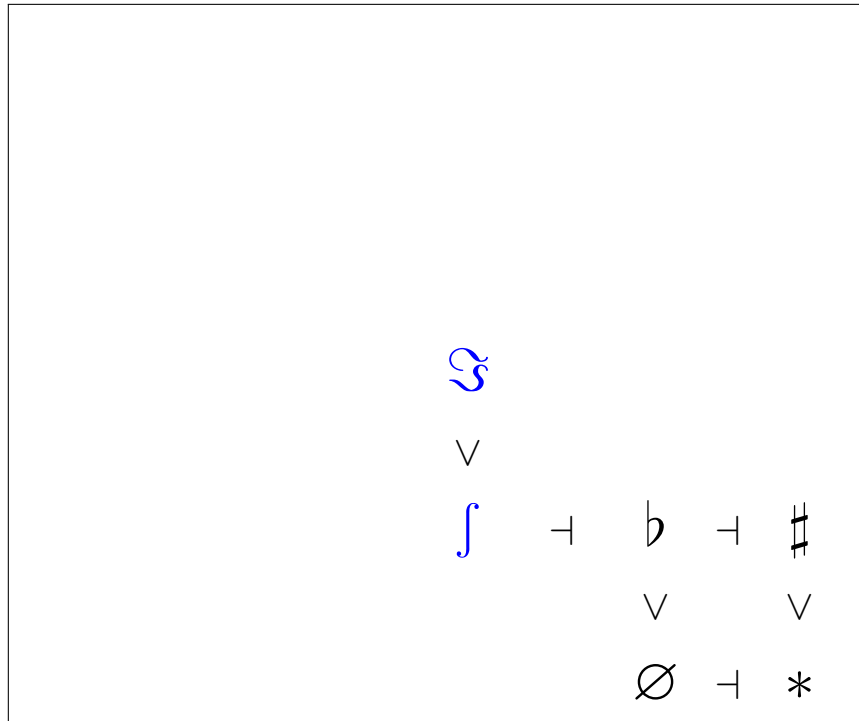
Π

Δ

*

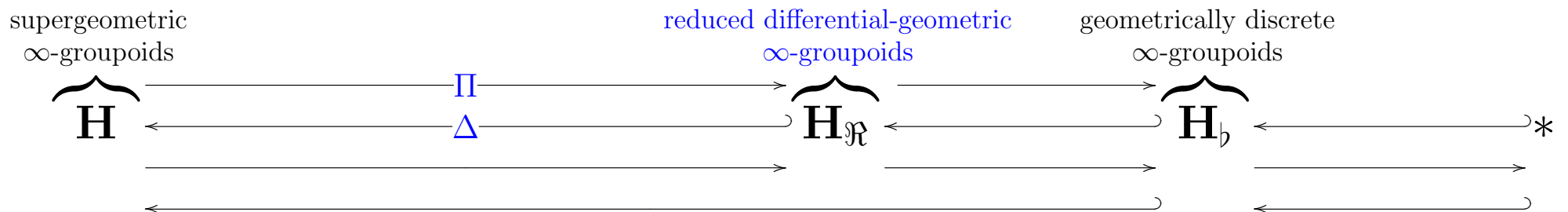


The Modalities of Super homotopy theory

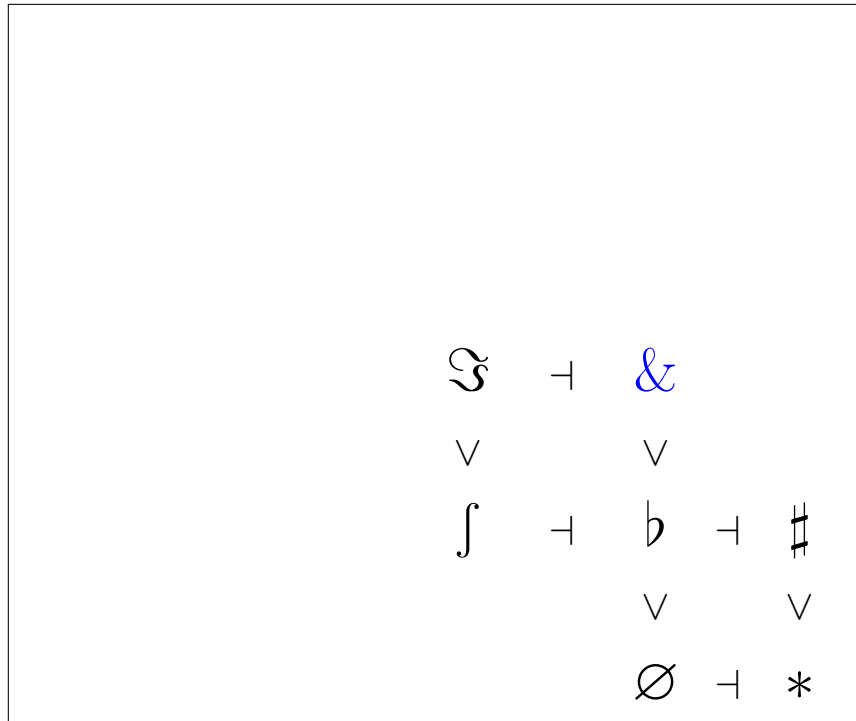


infinitesimally shaped

The terminal functor factors into a system of dualities = adjunctions.

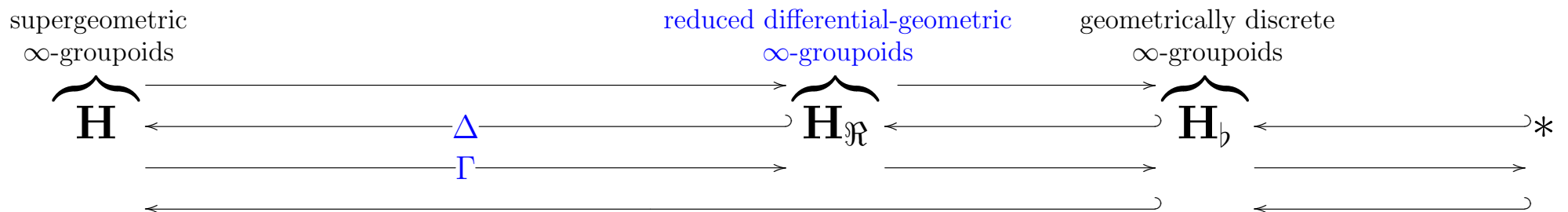


The Modalities of Super homotopy theory

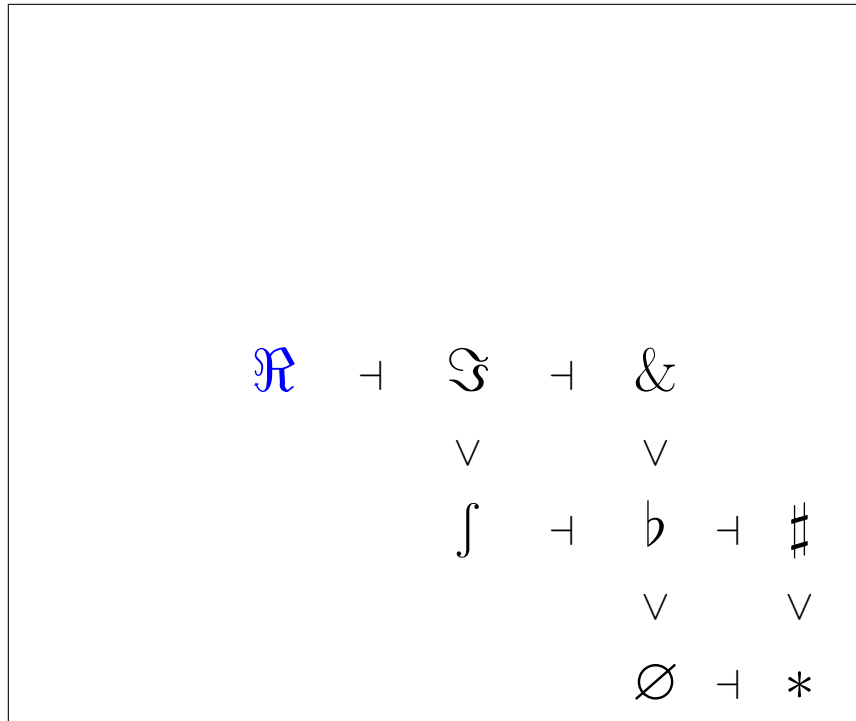


infinitesimally discrete

The terminal functor factors into a system of dualities = adjunctions.

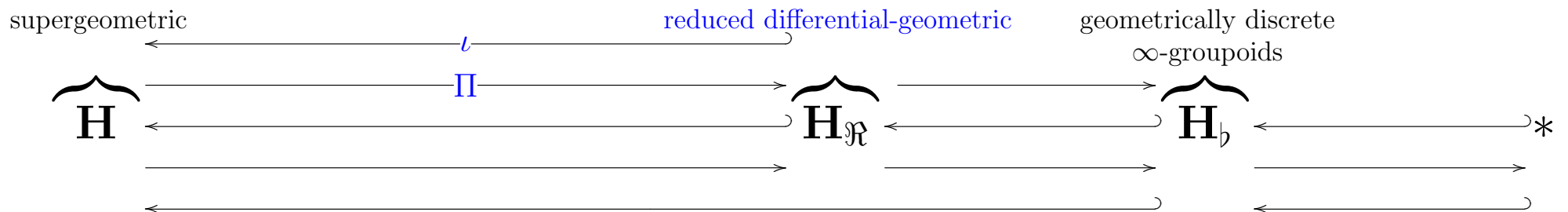


The Modalities of Super homotopy theory

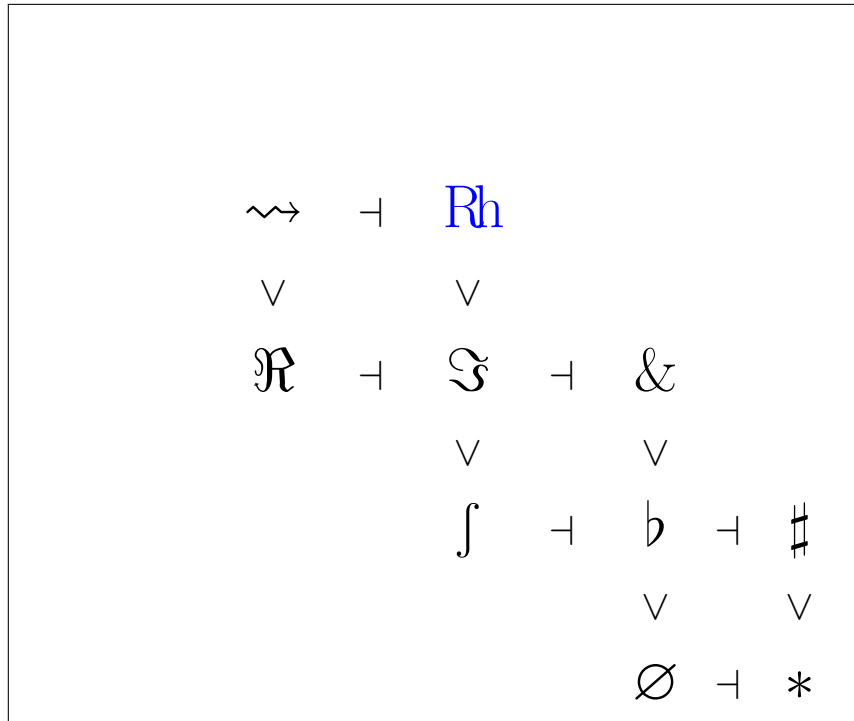


reduced

The terminal functor factors into a system of dualities = adjunctions.

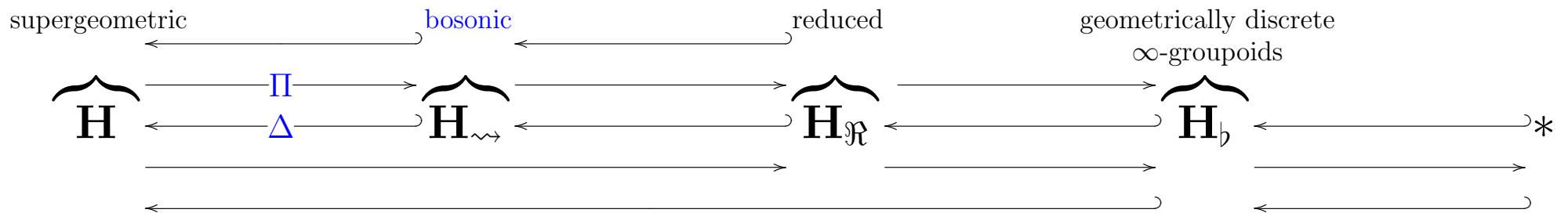


The Modalities of Super homotopy theory

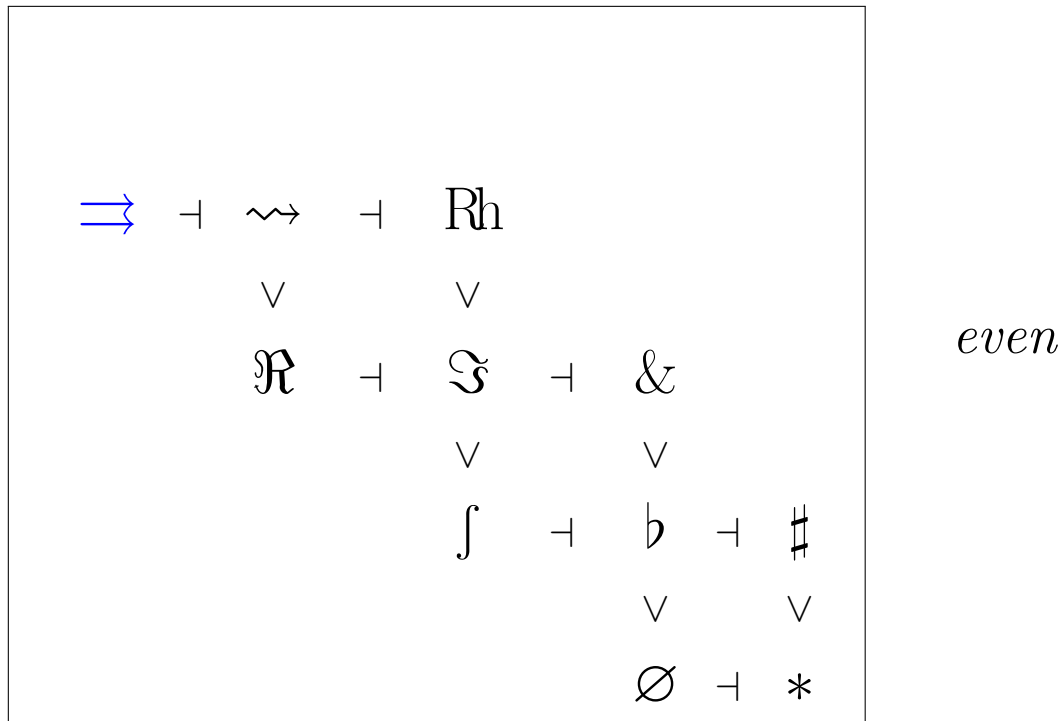


rheonomic

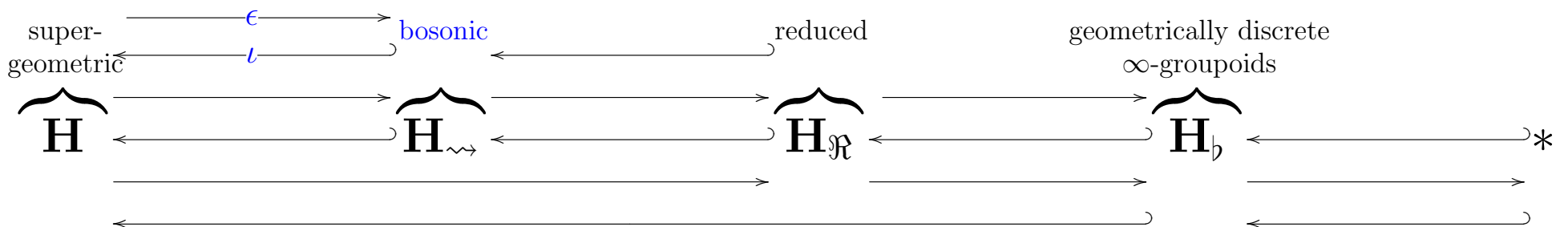
The terminal functor factors into a system of dualities = adjunctions.



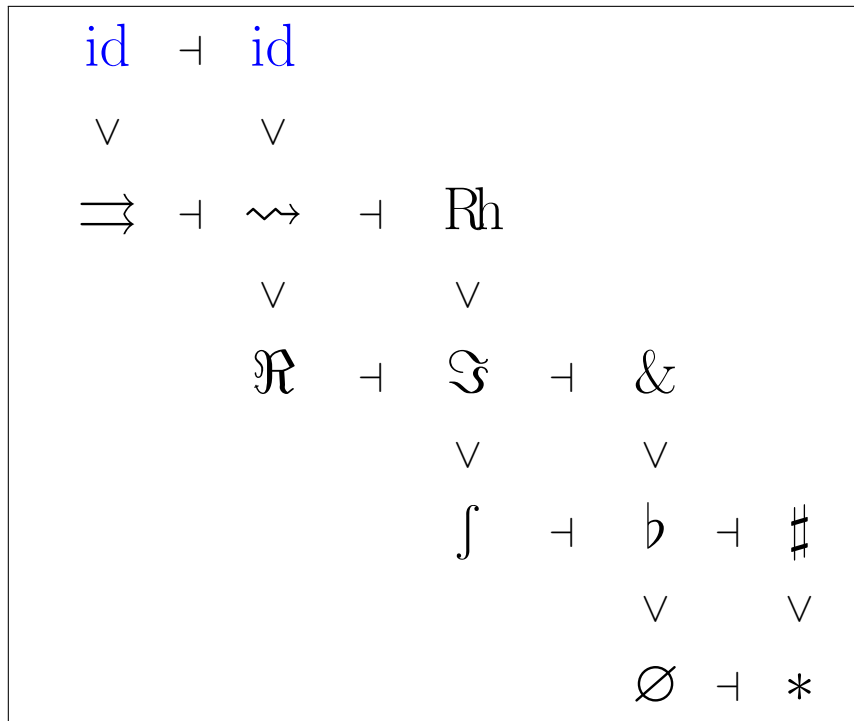
The Modalities of Super homotopy theory



The terminal functor factors into a system of dualities = adjunctions.

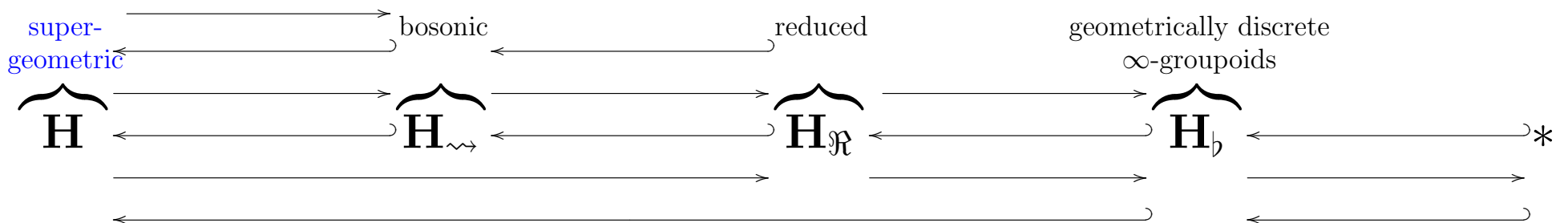


The Modalities of Super homotopy theory

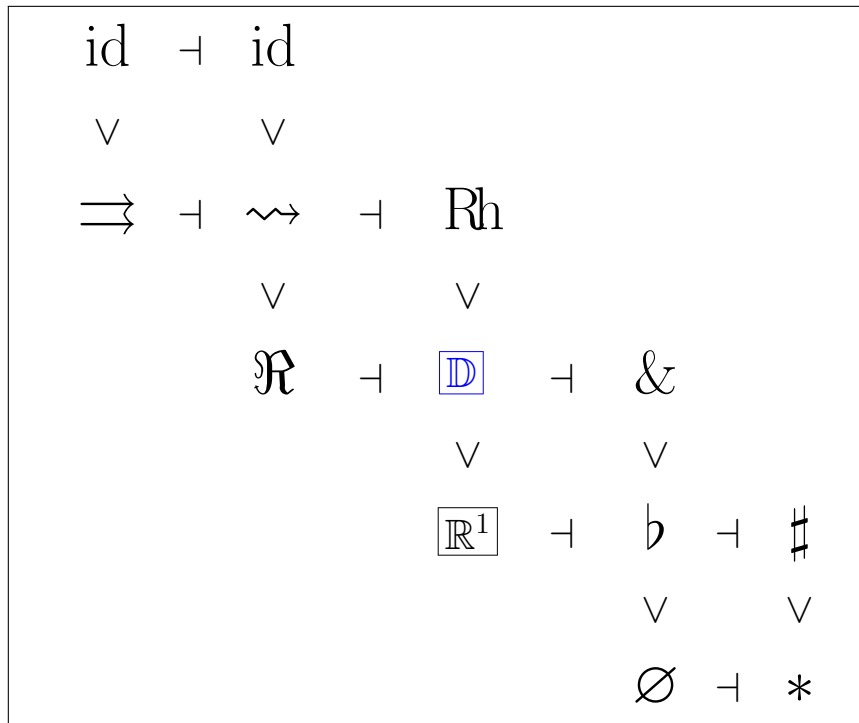


super-geometric

The terminal functor factors into a system of dualities = adjunctions.

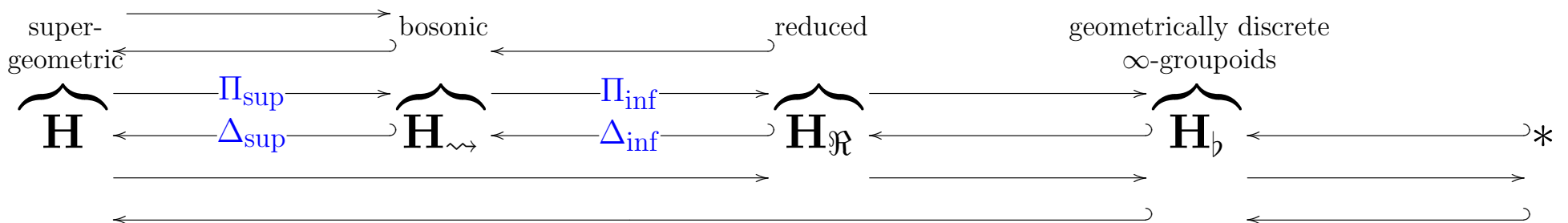


The Modalities of Super homotopy theory

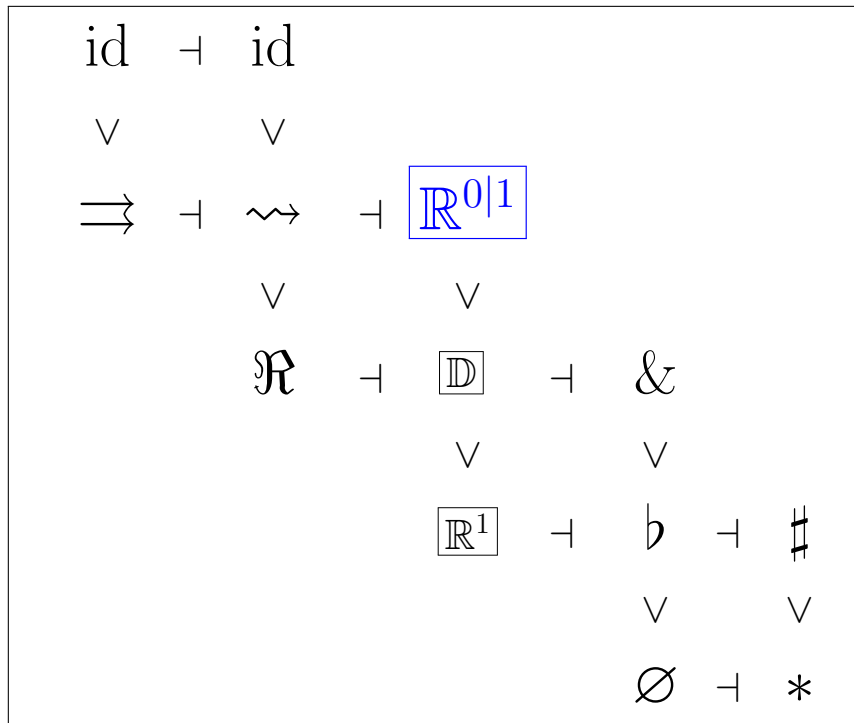


infinitum-local

The central modalities are motivic \mathbb{A}^1 -localizations.

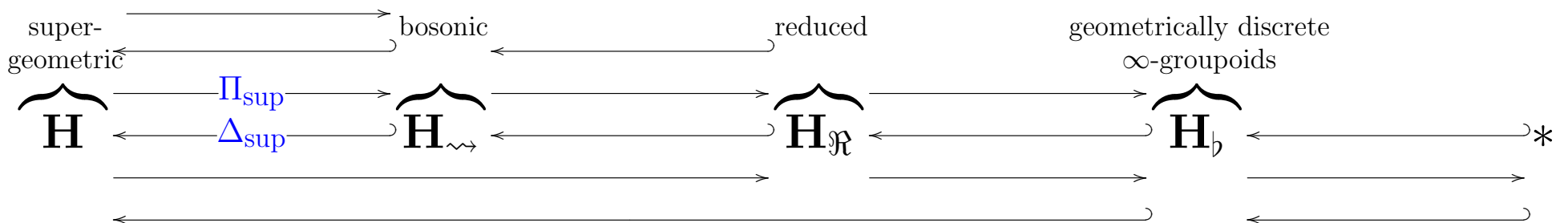


The Modalities of Super homotopy theory

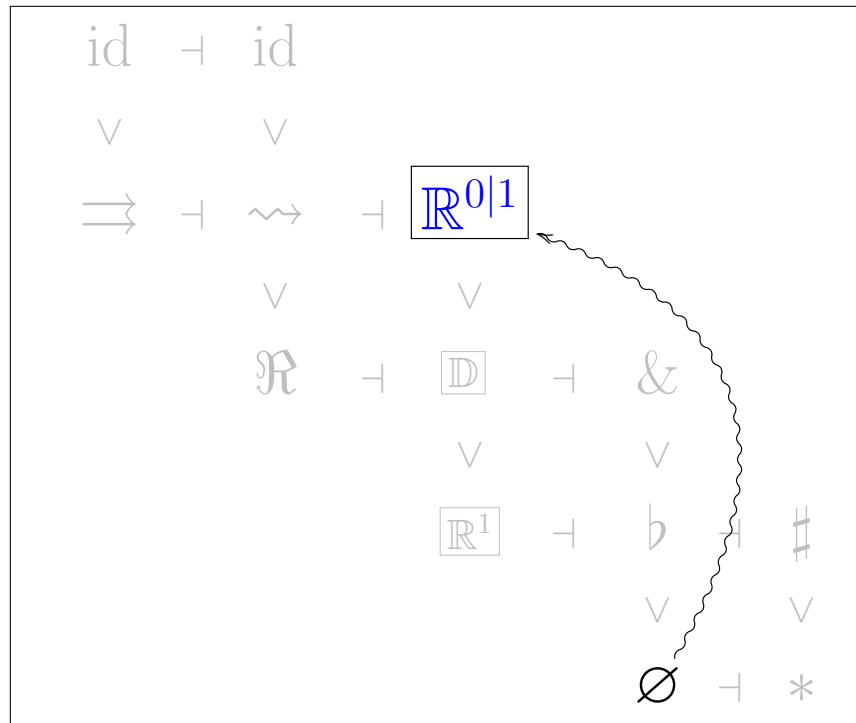


superpoint-local

The central modalities are motivic \mathbb{A}^1 -localizations.

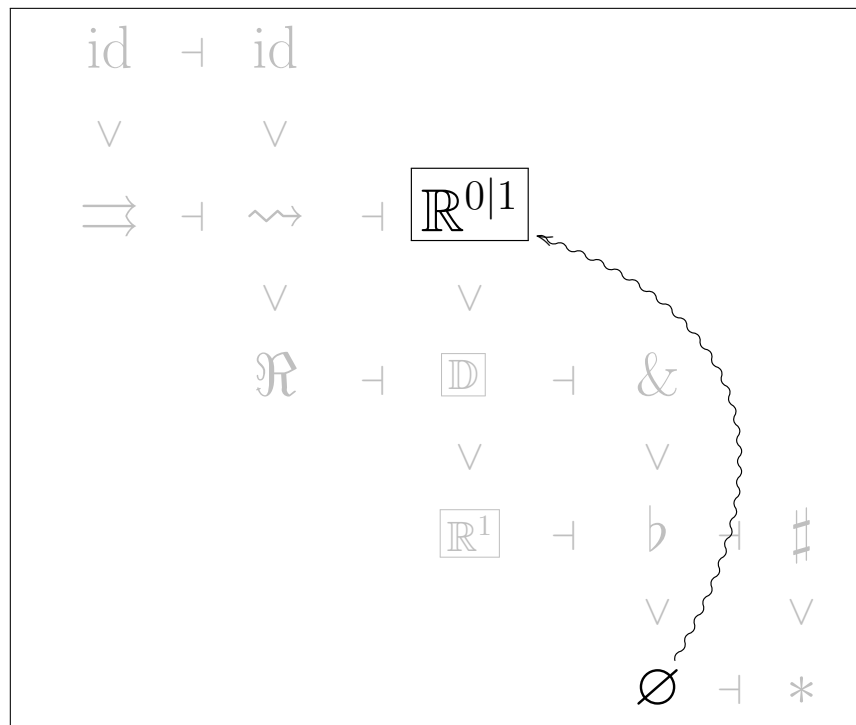


The Modalities of Super homotopy theory



\Rightarrow emergence of Atom of Superspace from \nothing

The Modalities of Super homotopy theory



\Rightarrow emergence of Atom of Superspace from \text{nothing}

now [apply the microscope of homotopy theory](#)

to discover what emerges, in turn, out of the superpoint...

Rational
Super homotopy theory
and the fundamental super p -Branes

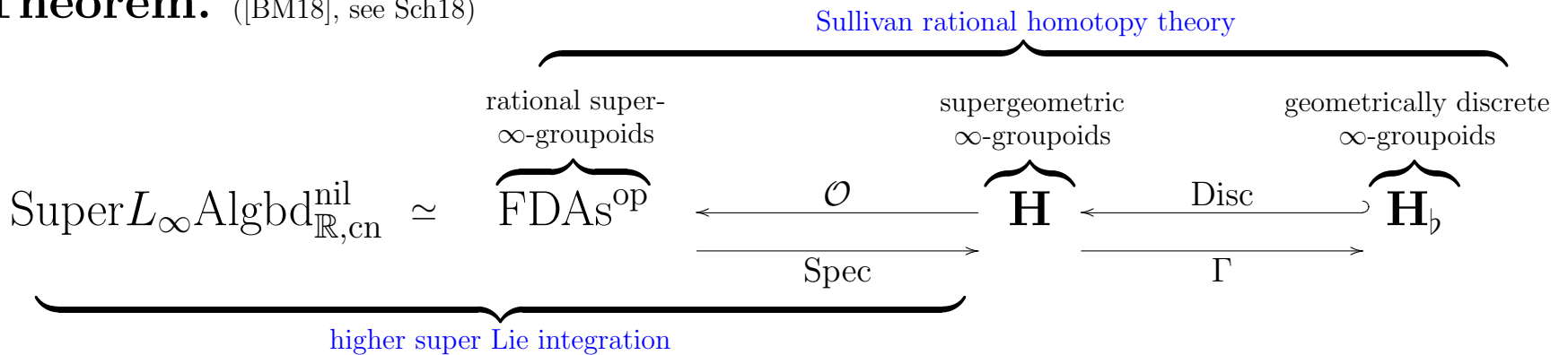
[back to Part I](#)

Higher super Lie theory and Rational homotopy

$\left. \begin{array}{l} \text{infinitesimal} \\ \text{rational} \end{array} \right\}$ approximation of super-homotopy by $\left\{ \begin{array}{l} \text{higher Lie integration} \\ \text{Sullivan construction} \end{array} \right.$

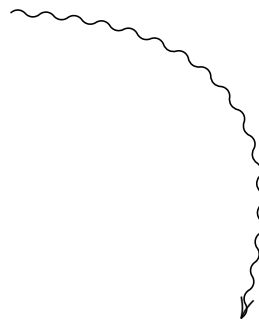
Definition. $\underbrace{\text{FDAs}}_{\substack{\text{terminology} \\ \text{common in} \\ \text{supergravity} \\ ([\text{vanNieuwenhuizen82}])}} := \underbrace{\text{dgcSuperAlg}_{\mathbb{R}, \text{cn}}}_{\substack{\infty\text{-category of} \\ \text{differential} \\ \text{graded-commutative} \\ \text{superalgebras}}} \xleftarrow[\cong]{\text{CE}} \left(\underbrace{\text{Super}L_\infty\text{Algbd}_{\mathbb{R}, \text{cn}}^{\text{nil}}}_{\substack{\infty\text{-category of} \\ \text{nilpotent} \\ \text{super } L_\infty\text{-algebroids}}} \right)^{\text{op}}$

Theorem. ([BM18], see Sch18)



Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

the Atom of Superspace



$\mathbb{R}^{0|1}$

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

$\mathbb{R}^{0|1}$

Type I

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

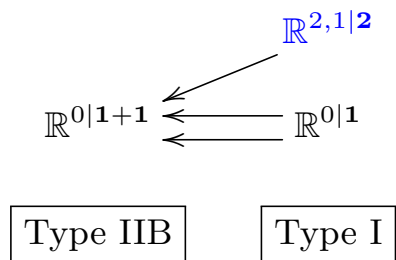
$$\mathbb{R}^{0|1+1} \leftarrow \leftarrow \mathbb{R}^{0|1}$$

Type IIB

Type I

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

[HS17]



universal central extension: $3d$ super-Minkowski spacetime

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

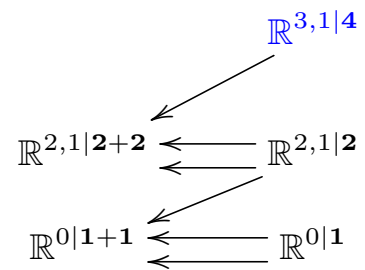
[HS17]

$$\begin{array}{ccc} \mathbb{R}^{2,1|2+2} & \longleftarrow & \mathbb{R}^{2,1|2} \\ & \longleftarrow & \\ \mathbb{R}^{0|1+1} & \longleftarrow & \mathbb{R}^{0|1} \end{array}$$

Type IIB

Type I

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet



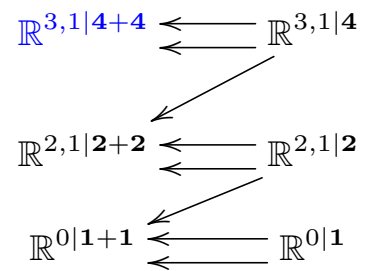
[HS17]

Type IIB

Type I

universal invariant central extension: 4d super-Minkowski spacetime

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

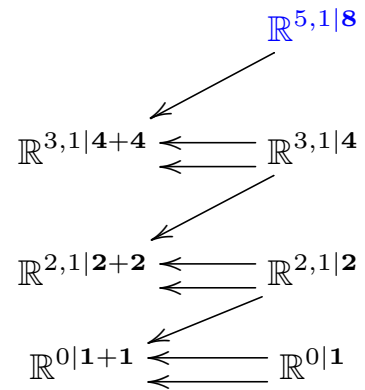


[HS17]

Type IIB

Type I

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

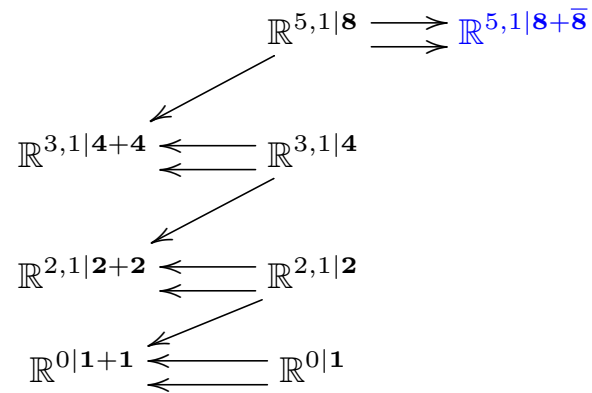


[HS17]

Type IIB

Type I

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet



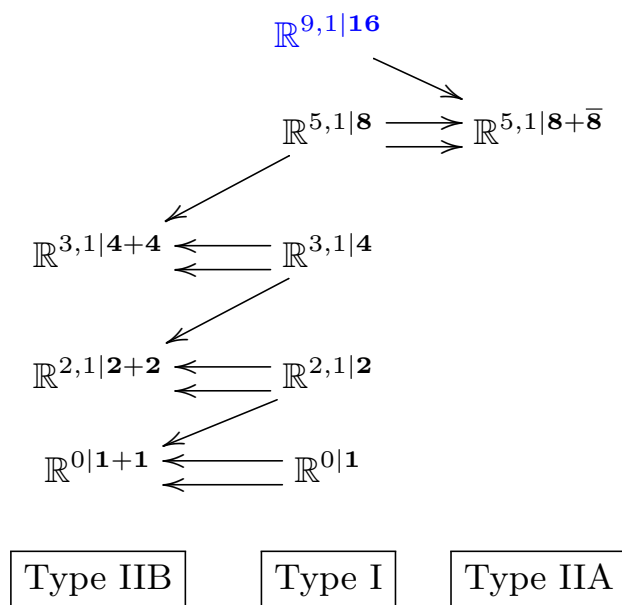
[HS17]

Type IIB

Type I

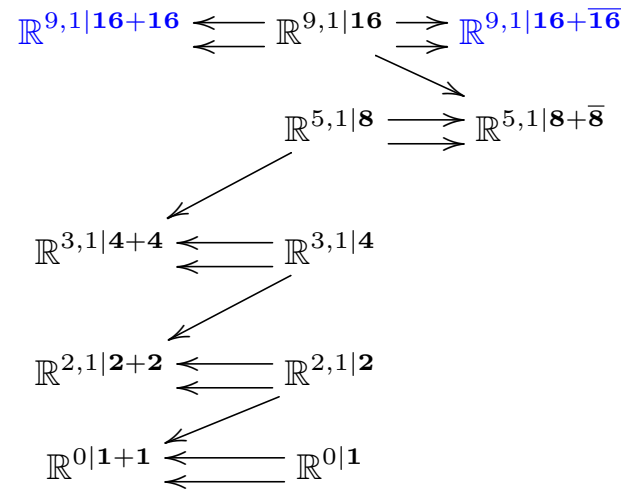
Type IIA

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet



[HS17]

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet



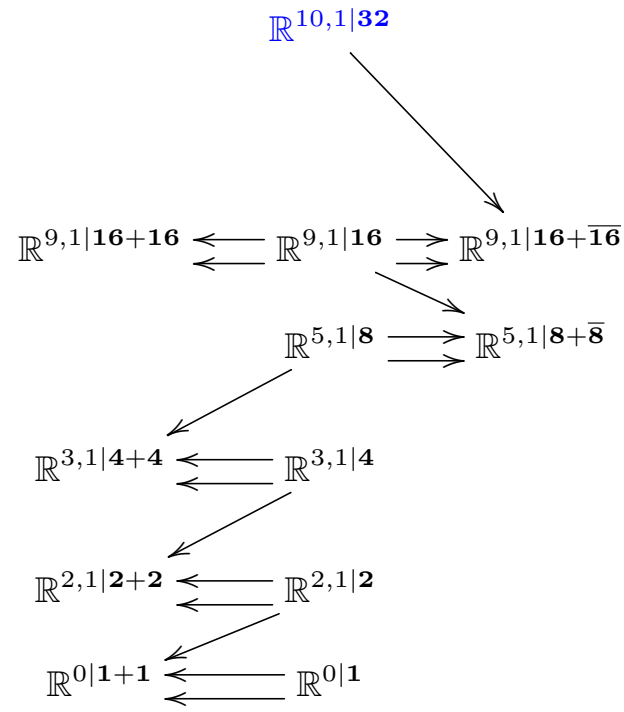
[HS17]

Type IIB

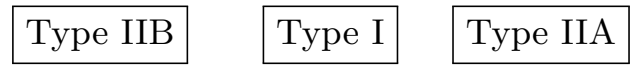
Type I

Type IIA

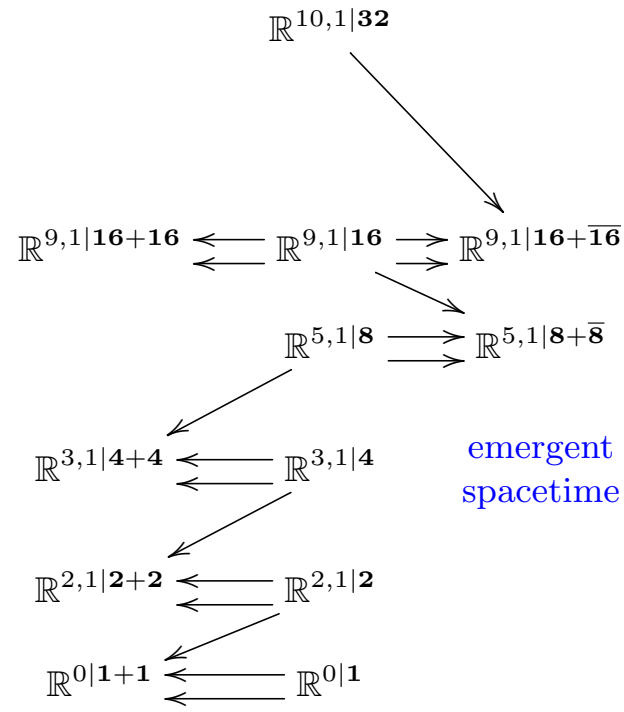
Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet



[HS17]

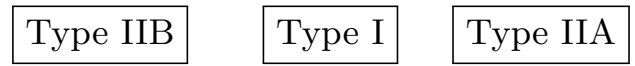


Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

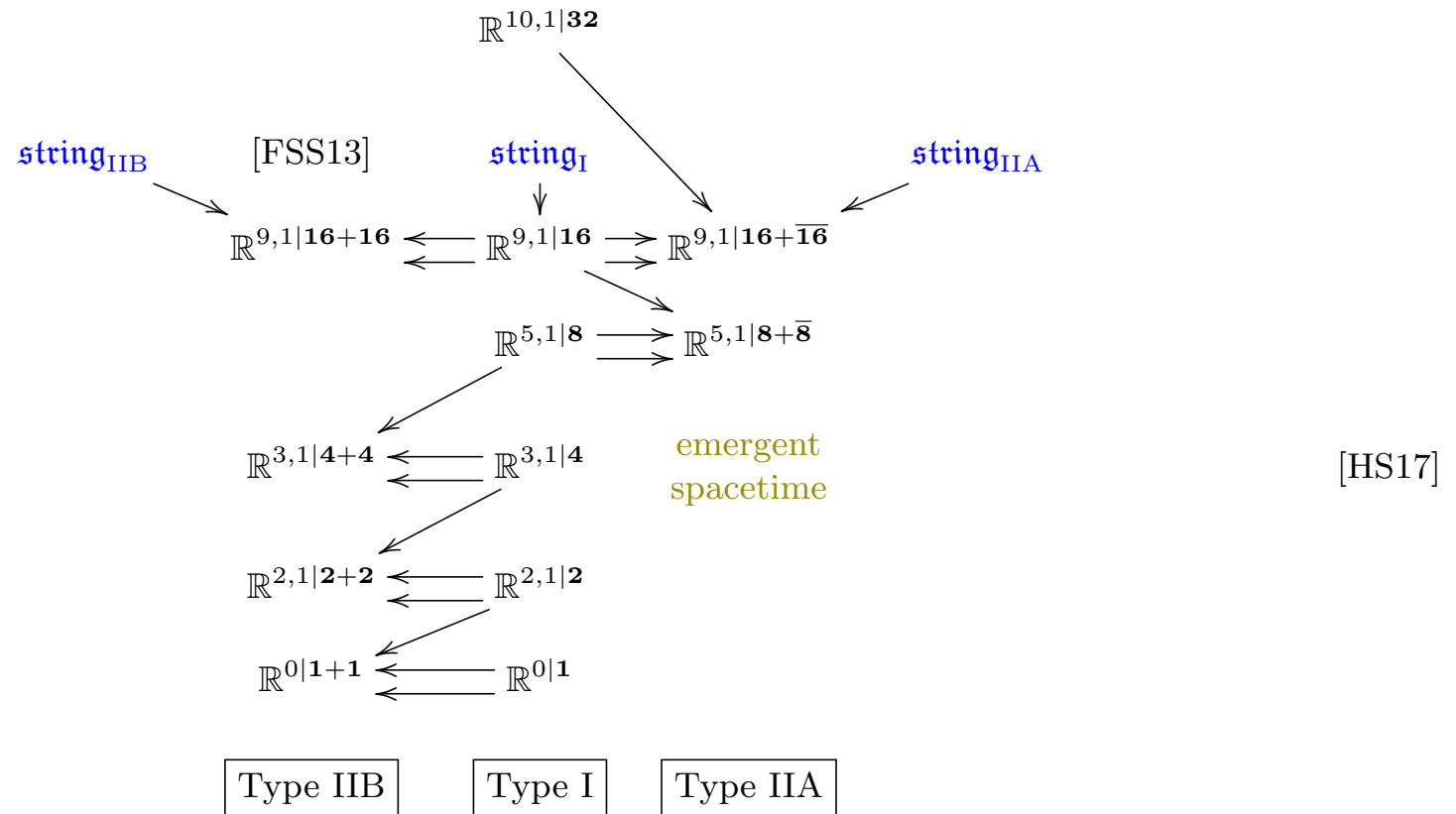


emergent
spacetime

[HS17]

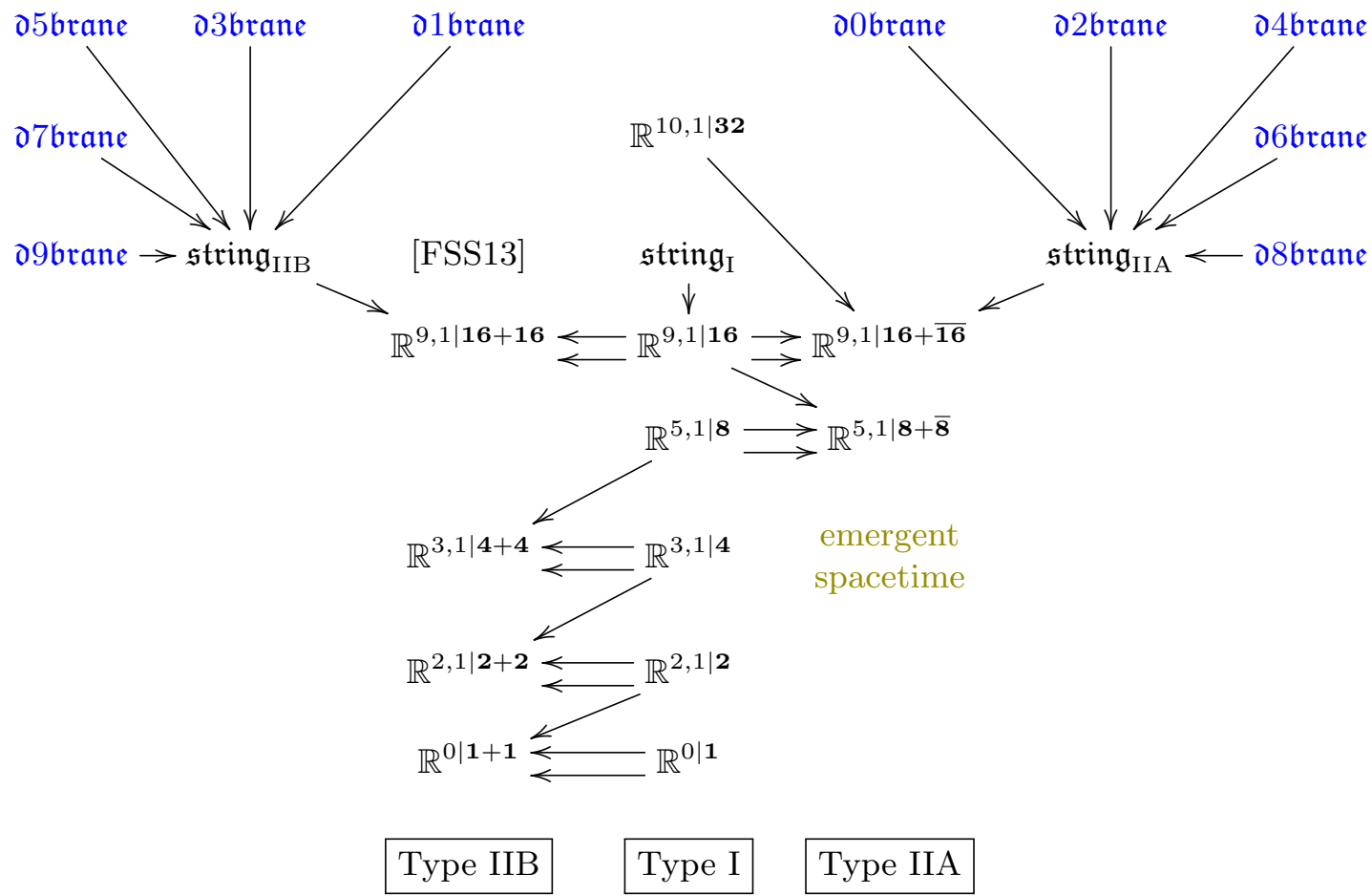


Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

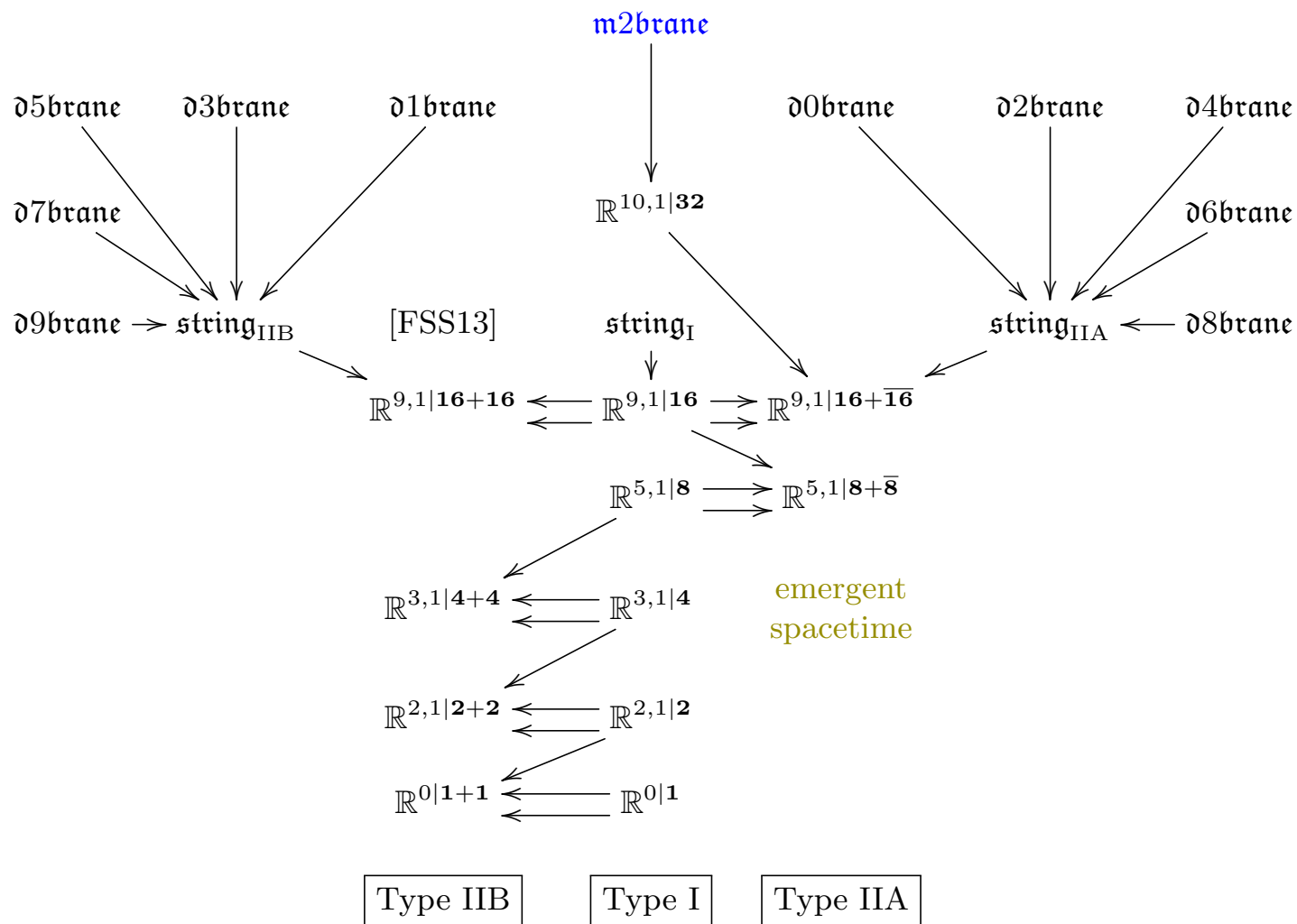


universal *higher* central invariant extension: stringy extended super-spacetimes

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

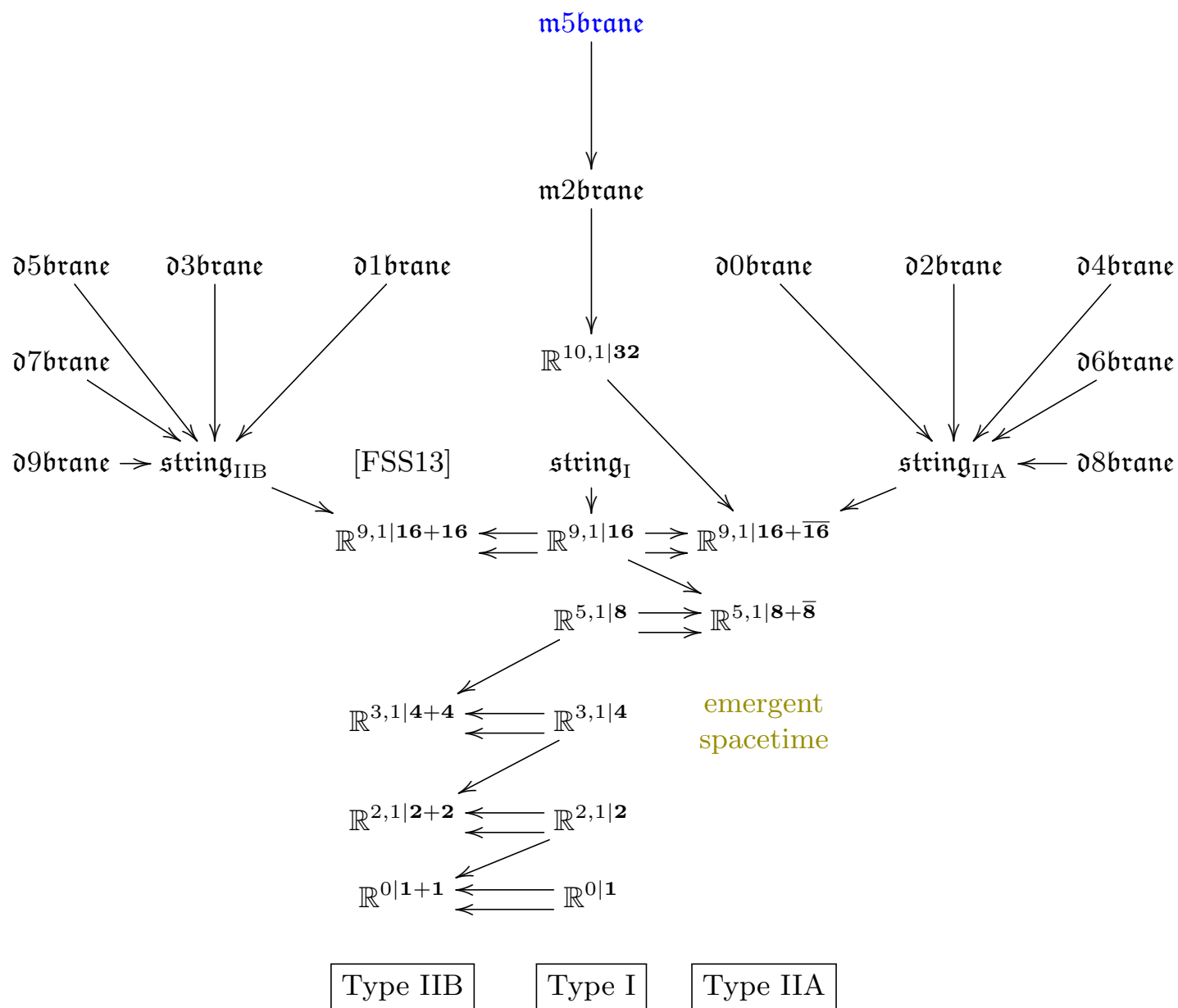


Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet



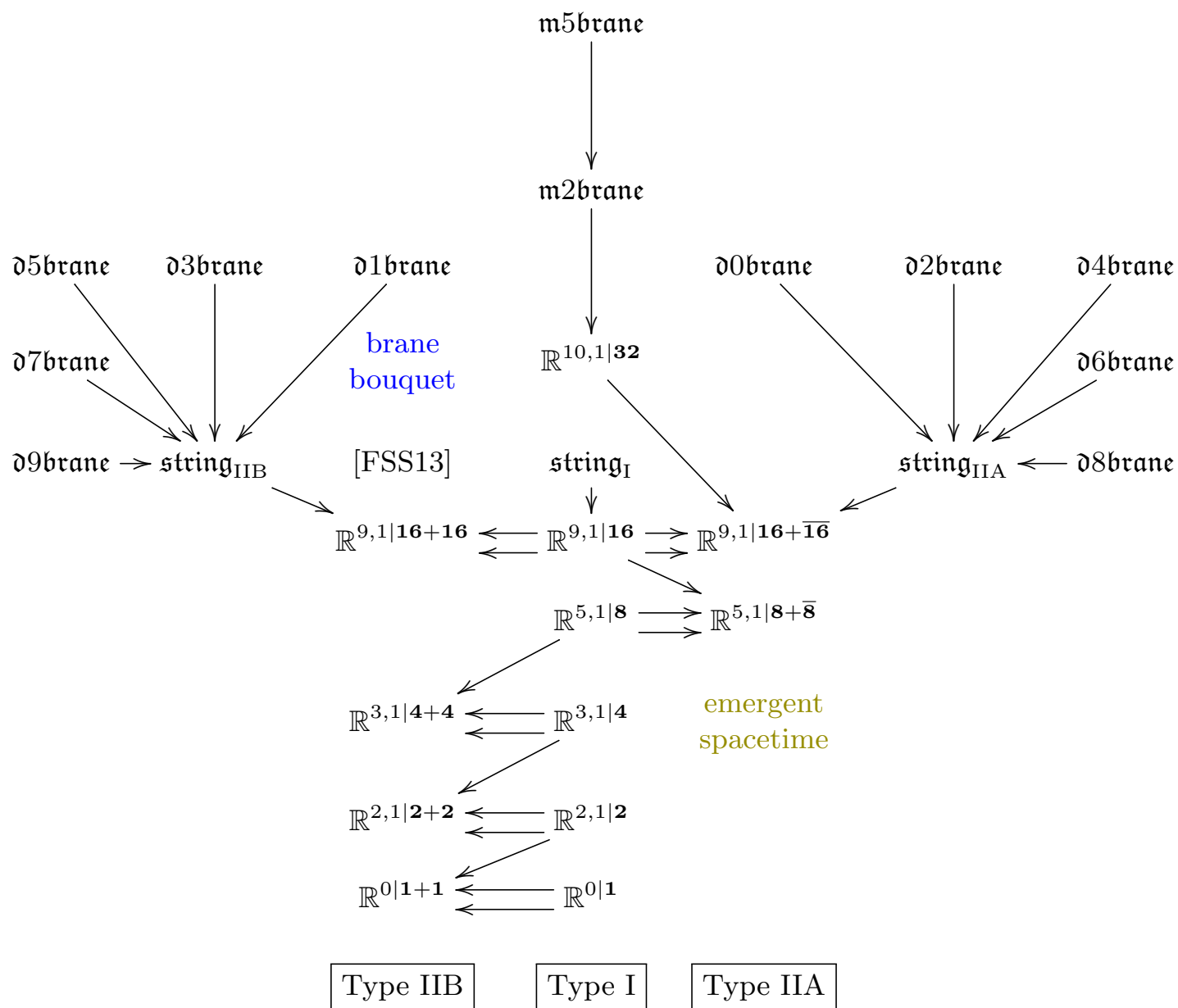
Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

[FSS15]

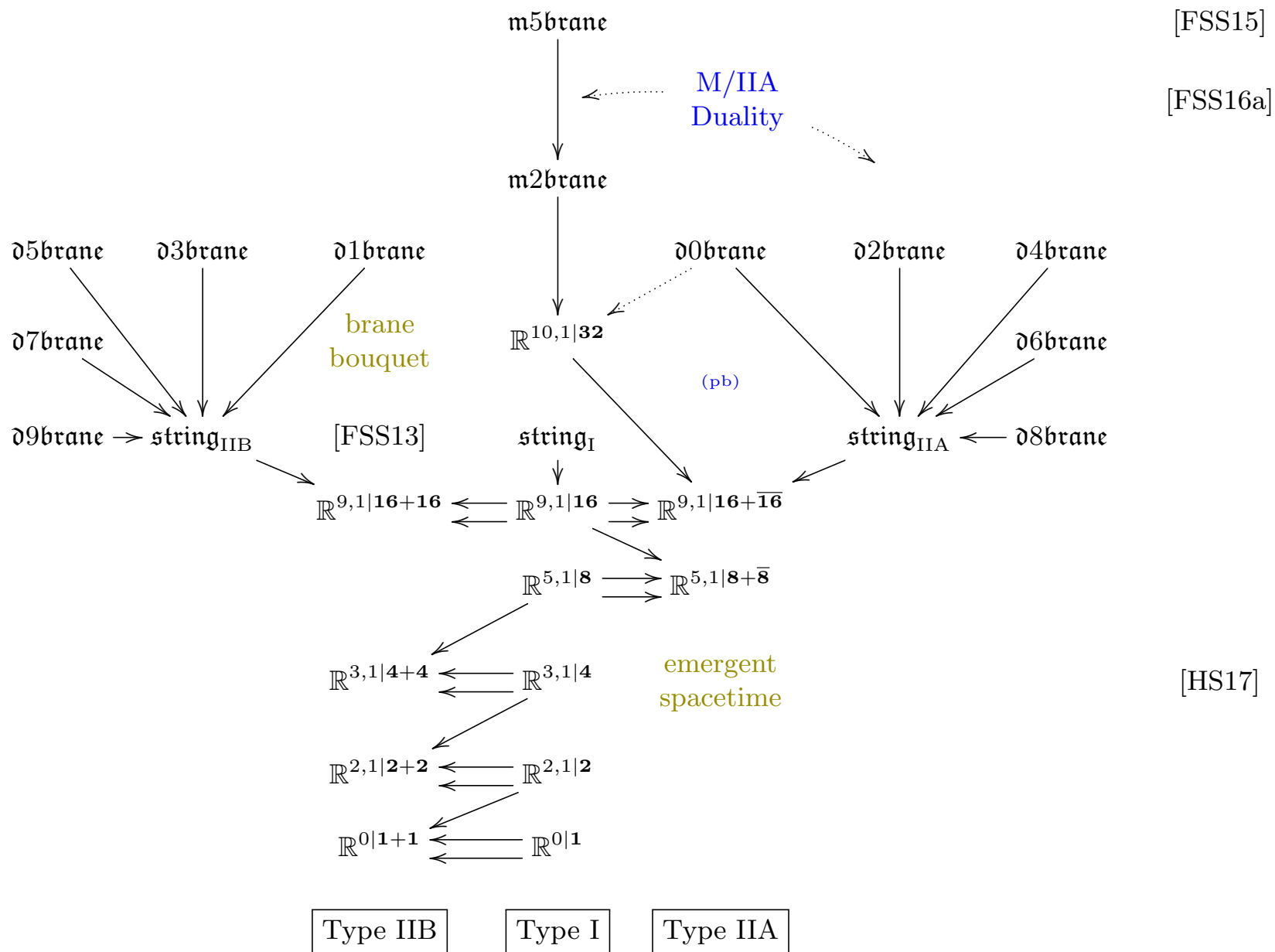


Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

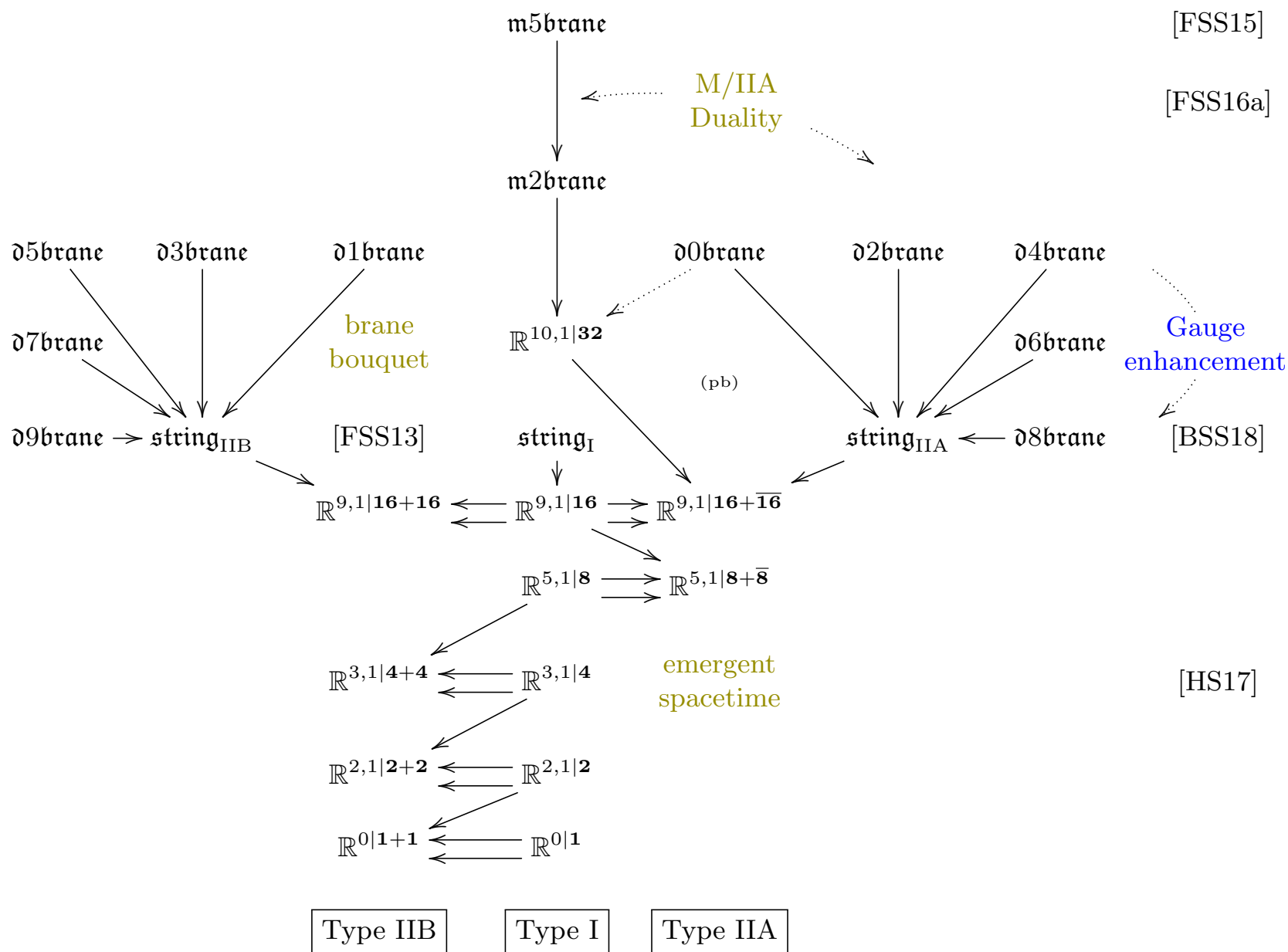
[FSS15]



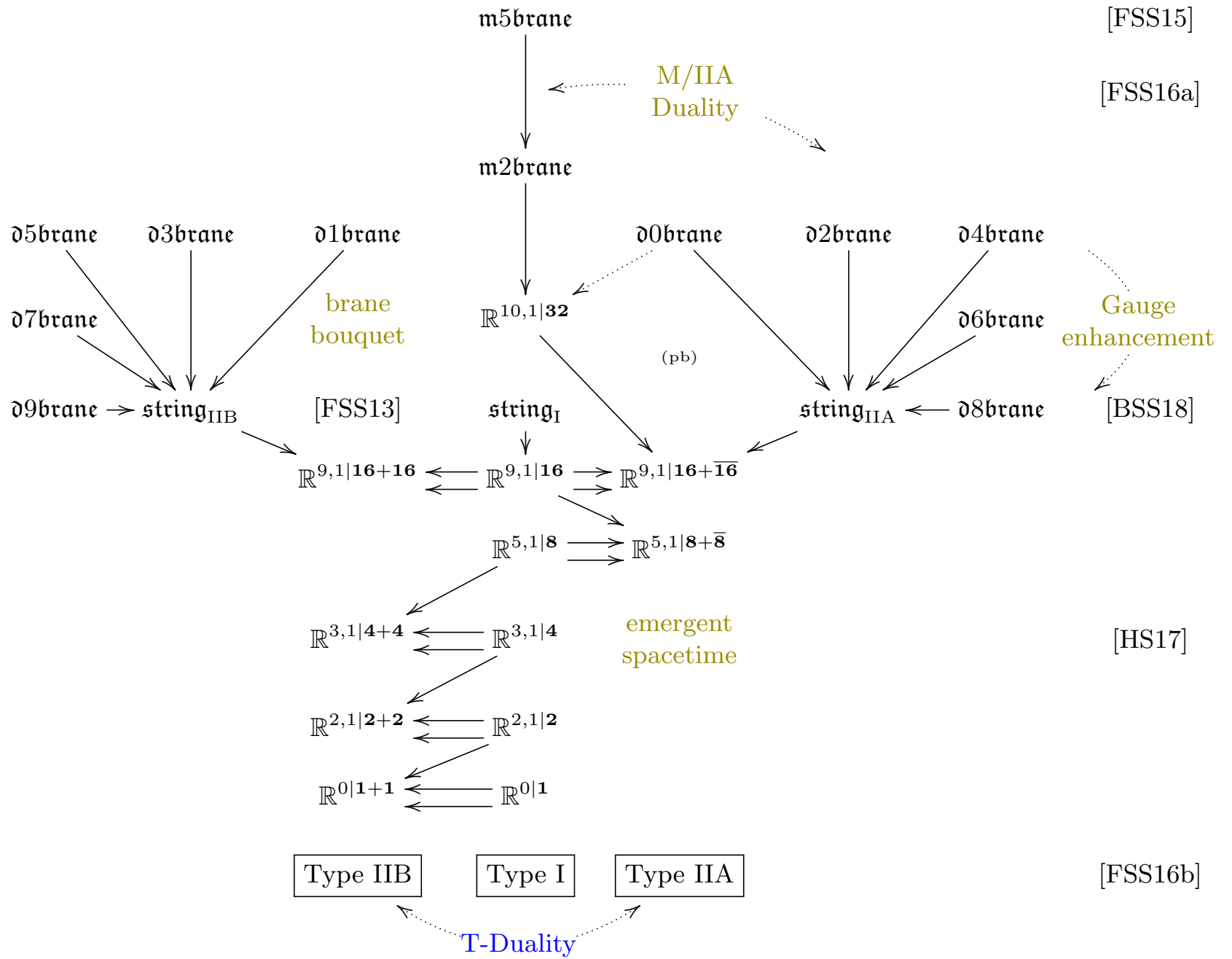
Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet



Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet



Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

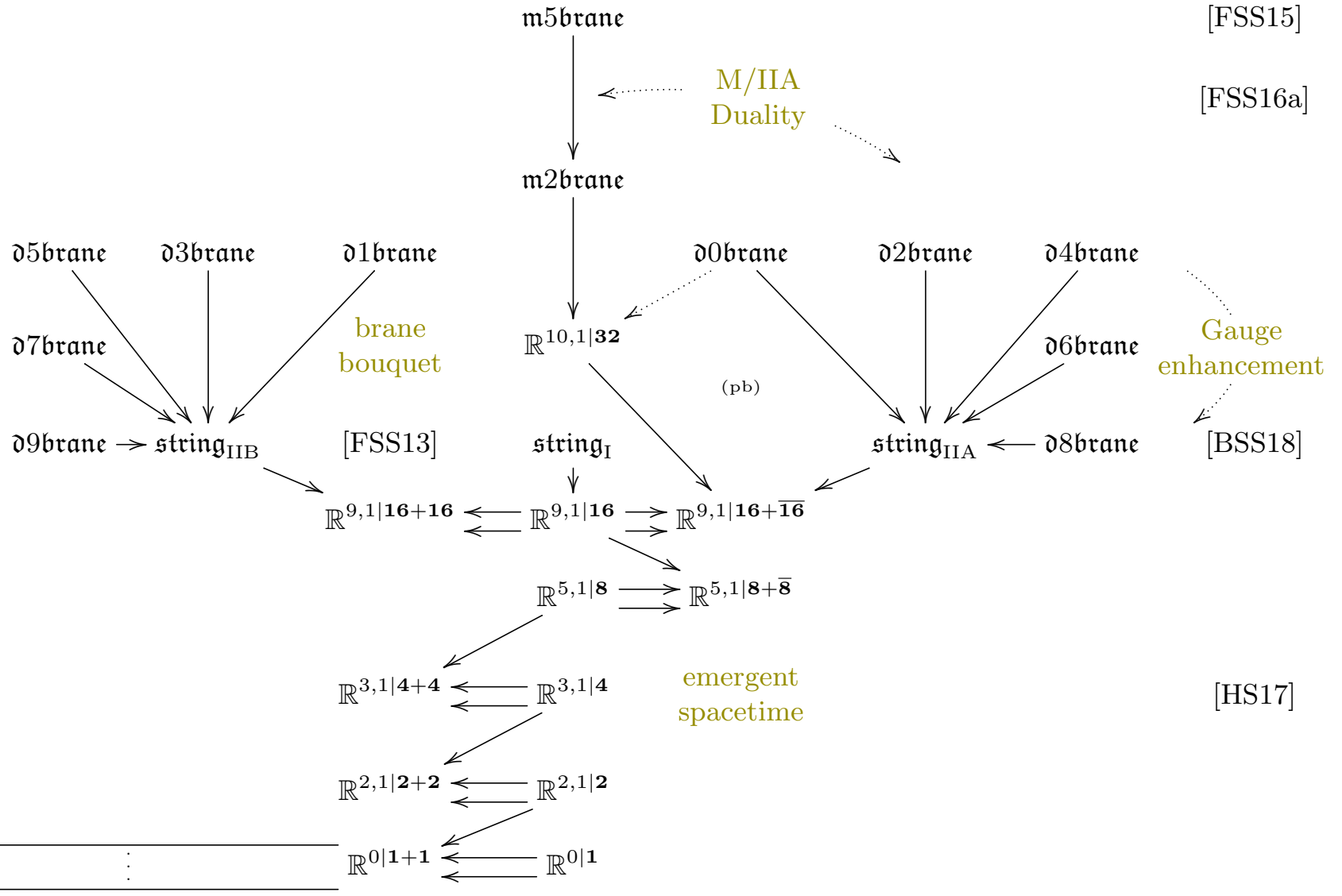


Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

[FSS18]

[FSS15]

[FSS16a]



[HS17]

Exceptional

Type IIB

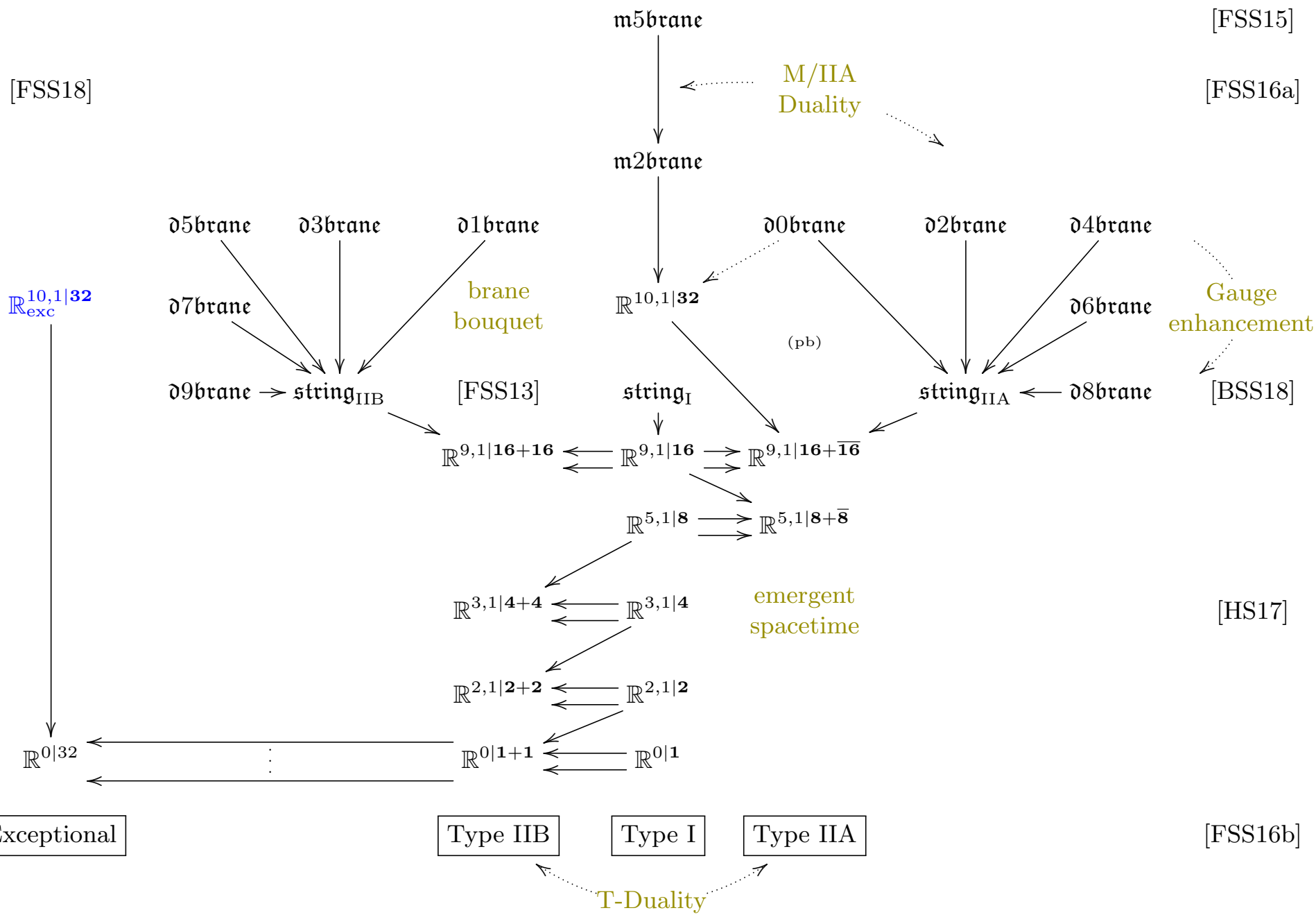
Type I

Type IIA

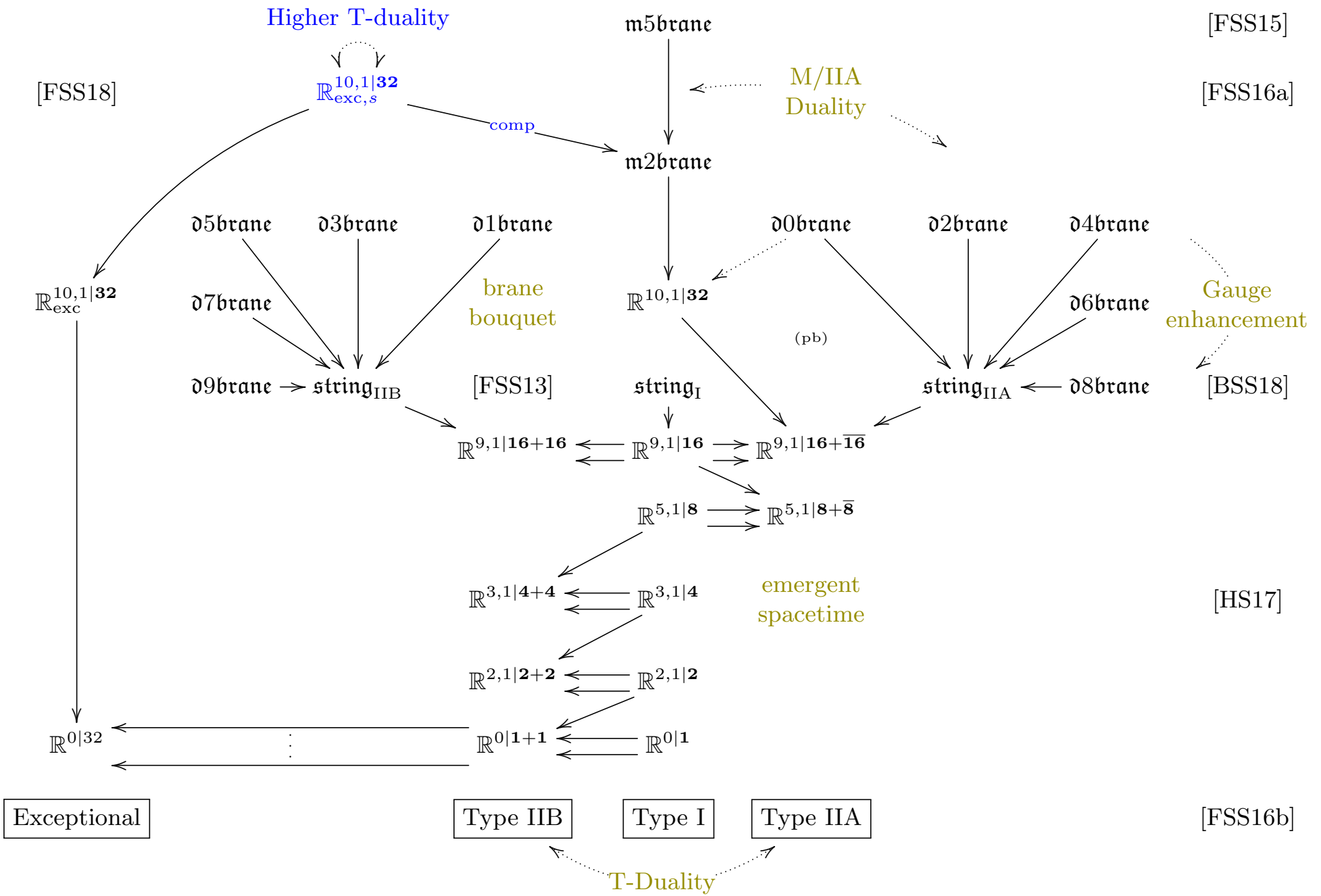
[FSS16b]

T-Duality

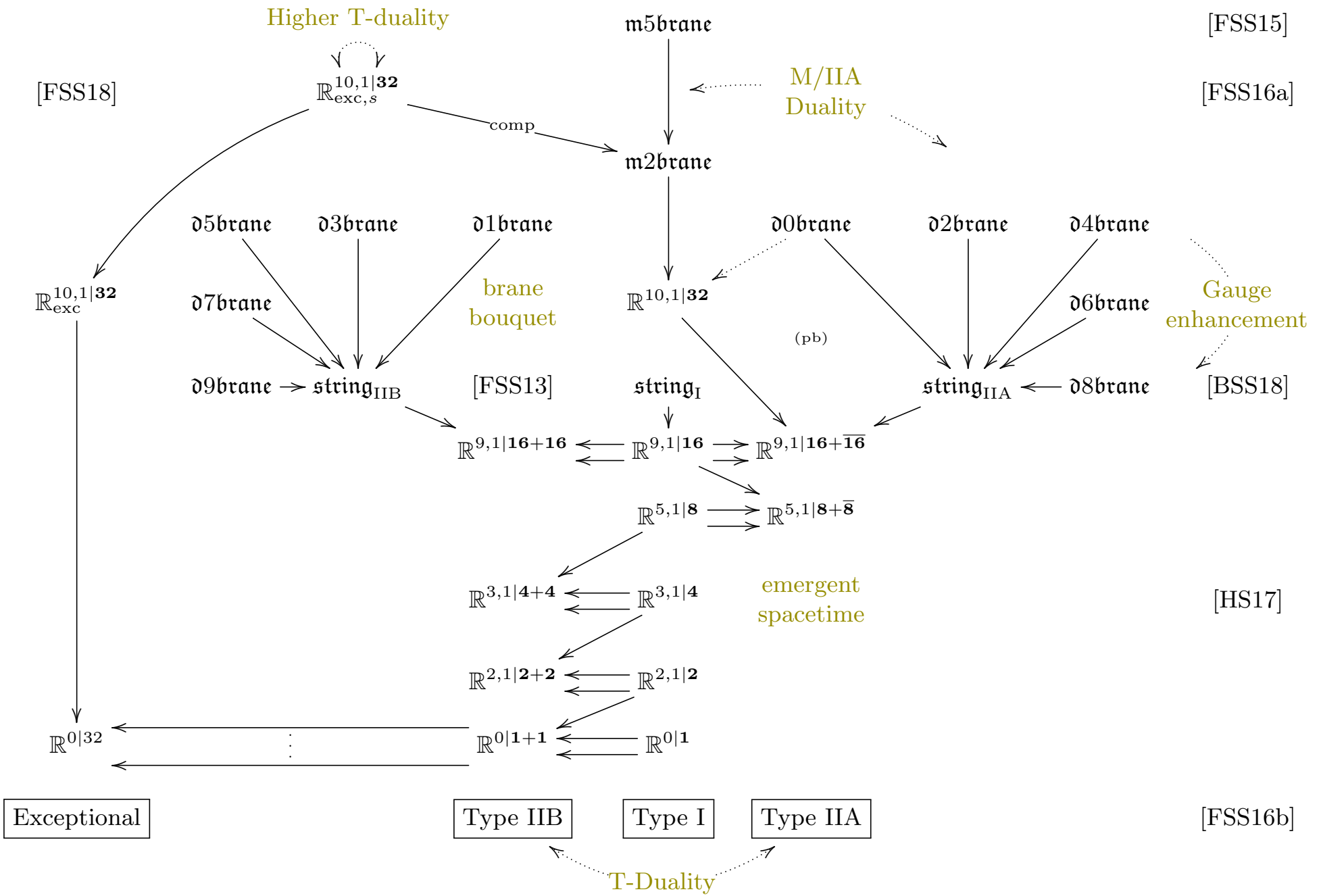
Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet



Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

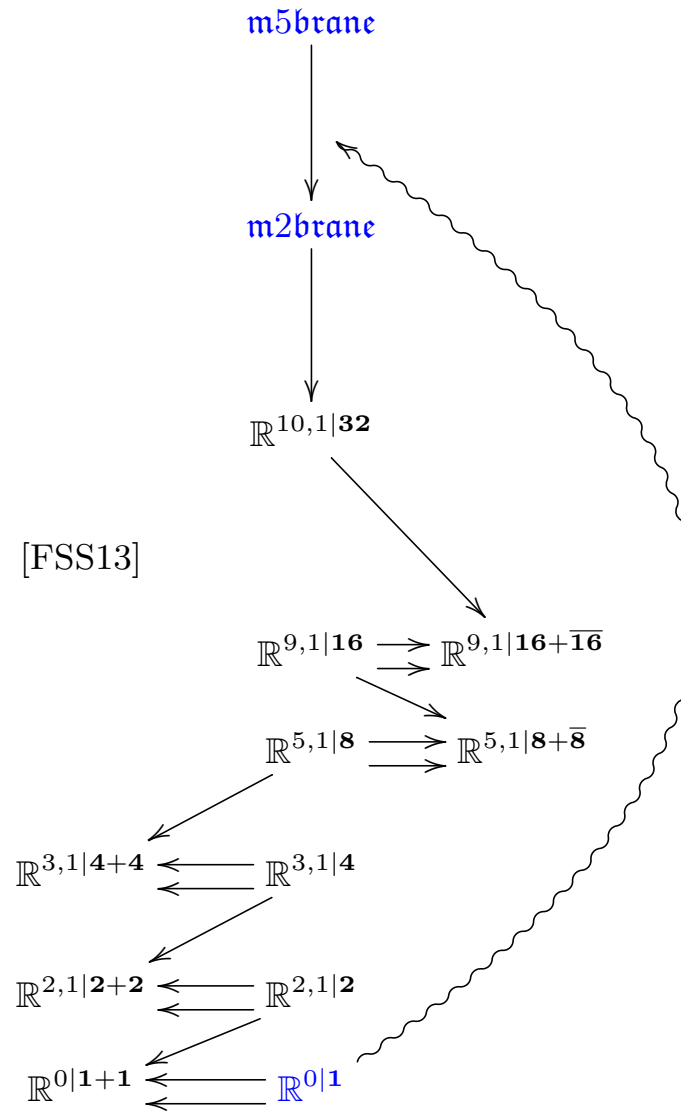


Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet



Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

[FSS15]



emergence of fundamental M-branes from the Atom of Superspace

Universal central invariant super- L_∞ extensions of $\mathbb{R}^{0|1}$: Brane bouquet

[FSS15]

m5brane



m2brane

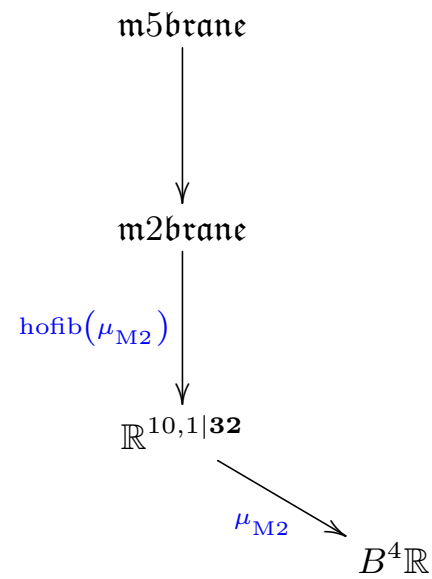


$\mathbb{R}^{10,1|32}$

zoom in on the fundamental M-brane super-extensions

The fundamental M2/M5-brane cocycle

[FSS15]

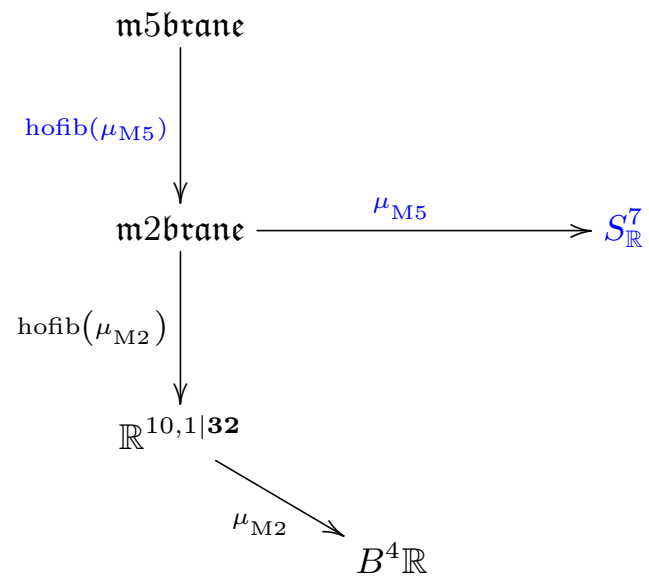


$$\mu_{M2} = dL_{M2}^{\text{WZW}} = \frac{i}{2} (\bar{\psi} \Gamma_{a_1 a_2} \psi) \wedge e^{a_1} \wedge e^{a_2}$$

the [WZW-curvature](#) of the Green-Schwarz-type sigma-model [super-membrane](#)

The fundamental M2/M5-brane cocycle

[FSS15]

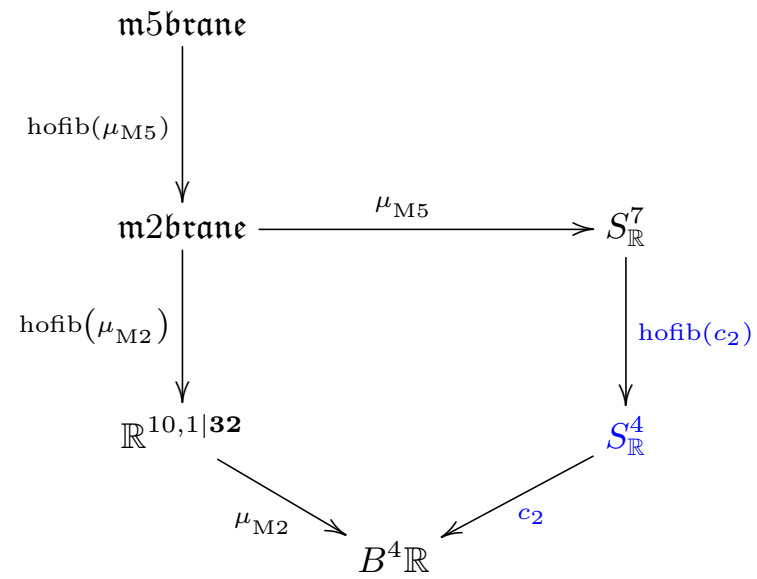


$$\mu_{M5} = dL_{M5}^{\text{WZW}} = \frac{1}{5!} (\bar{\psi} \Gamma_{a_1 \dots a_5} \psi) \wedge e^{a_1} \wedge \dots \wedge e^{a_5} + c_3 \wedge \frac{i}{2} (\bar{\psi} \Gamma_{a_1 a_2} \psi) \wedge e^{a_1} \wedge e^{a_2}$$

the **WZW-curvature** of the Green-Schwarz-type sigma-model **super-fivebrane**

The fundamental M2/M5-brane cocycle

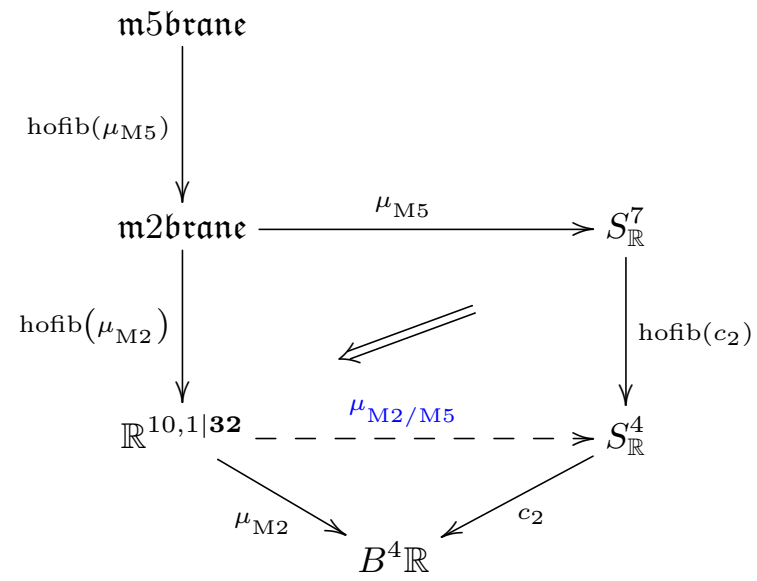
[FSS15]



the [quaternionic Hopf fibration](#) (in rational homotopy theory)

The fundamental M2/M5-brane cocycle

[FSS15]



the unified M2/M5-cocycle

The fundamental M2/M5-brane cocycle

[FSS15]

$$\mathbb{R}^{10,1|\mathbf{32}} \xrightarrow{\mu_{\text{M2/M5}}} S_{\mathbb{R}}^4$$

the unified M2/M5-cocycle is in rational Cohomotopy in degree 4

The fundamental M2/M5-brane cocycle

[FSS15]

$$\mathbb{R}^{10,1|\mathbf{32}} \xrightarrow{\mu_{\text{M2/M5}}} S_{\mathbb{R}}^4$$

$$\frac{i}{2} (\bar{\psi} \Gamma_{a_1 a_2} \psi) \wedge e^{a_1} \wedge e^{a_2} \longleftarrow G_4$$

$$\frac{1}{5!} (\bar{\psi} \Gamma_{a_1 \dots a_5} \psi) \wedge e^{a_1} \dots e^{a_5} \longleftarrow G_7$$

$$\text{Sullivan model: } \mathcal{O}(S_{\mathbb{R}}^4) \simeq \mathbb{R}[G_4, G_7] / \left(\begin{array}{l} dG_4 = 0 \\ dG_7 = -\frac{1}{2} G_4 \wedge G_4 \end{array} \right)$$

= 11d supergravity equations of motion of the C -field ([Sati13, Sect. 2.5])

The fundamental M2/M5-brane cocycle

[FSS15]

$$\mathbb{R}^{10,1|\mathbf{32}} \xrightarrow{\mu_{\text{M2/M5}}} S_{\mathbb{R}}^4$$

the unified M2/M5-cocycle

$$\begin{array}{ccc} \mathbb{R}^{10,1|\mathbf{32}} & \xrightarrow{\mu_{\text{M2/M5}}} & S_{\mathbb{R}}^4 \\ & \downarrow \text{double dimensional reduction \& gauge enhancement} & \\ \mathbb{R}^{9,1|\mathbf{16}+\overline{\mathbf{16}}} & \xrightarrow{\mu_{F1/D2p}} & \text{ku} // B^2\mathbb{R} \end{array}$$

D-brane charge in twisted K-theory, rationally
[BSS18]

The rational conclusion.

In $\left\{ \begin{array}{l} \text{infinitesimal} \\ \text{rational} \end{array} \right\}$ approximation

brane charge quantization follows from first principles

and reveals this situation:

brane species	cohomology theory of charge quantization
D-branes	twisted K-theory
M-branes	Cohomotopy in degree 4

The rational conclusion.

In $\left\{ \begin{array}{c} \text{infinitesimal} \\ \text{rational} \end{array} \right\}$ approximation

brane charge quantization follows from first principles
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brane species	cohomology theory of charge quantization
D-branes	twisted K-theory
M-branes	Cohomotopy in degree 4

Lift beyond $\left\{ \begin{array}{c} \text{infinitesimal} \\ \text{rational} \end{array} \right\}$ approximation is not unique

but one lift of rational Cohomotopy is *minimal* (in number of cells):
actual Cohomotopy represented by the actual 4-sphere

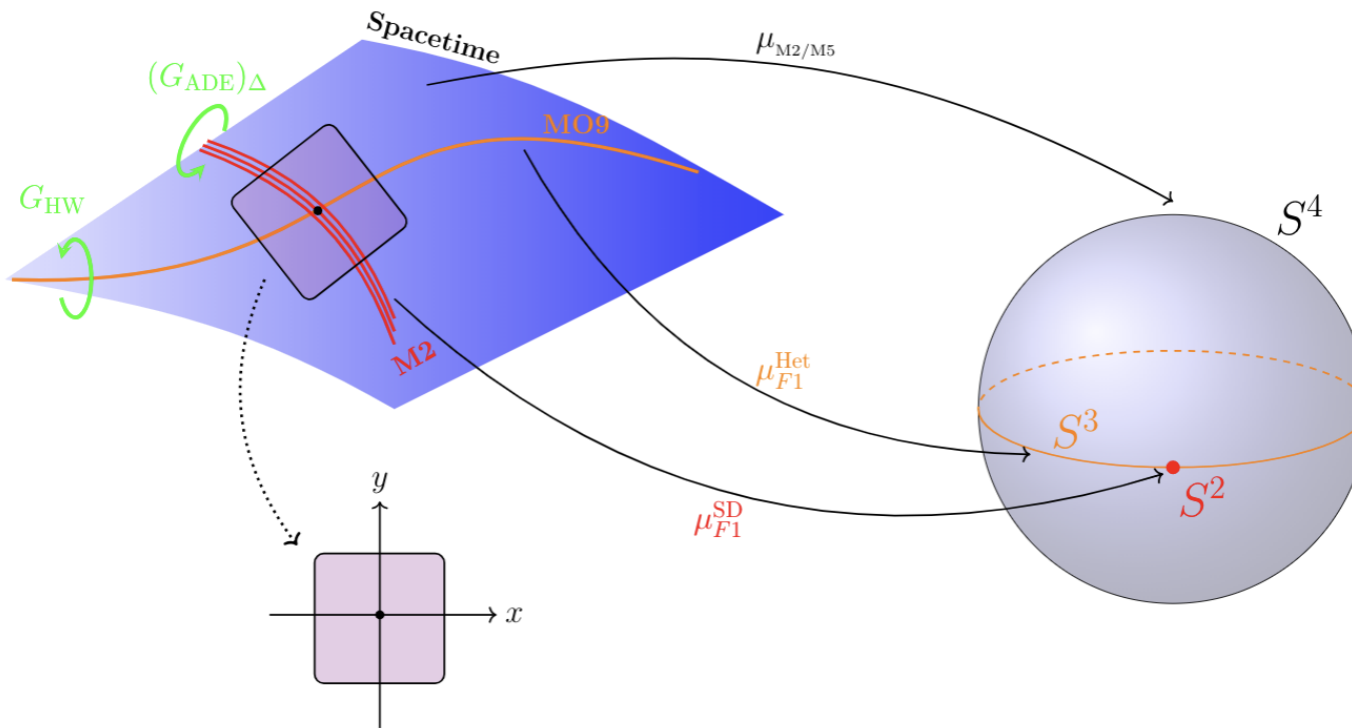
$$\begin{array}{ccc} & & S^4 \\ & \nearrow \text{cocycle in} & \downarrow \text{rationalization} \\ & \text{actual Cohomotopy} & \\ X & \xrightarrow{\text{cocycle in}} & S^4_{\mathbb{R}} \\ & \text{rational cohomotopy} & \end{array}$$

Towards microscopic M-theory

1. Construct

differential equivariant Cohomotopy \widehat{S}_γ^4
of 11d super-orbifold spacetimes \mathcal{X}

2. lifting super-tangent-space-wise the fundamental M2/M5-brane cocycle.



3. Compare the resulting observables on M-brane charge quantized supergravity field moduli with expected limiting corners of M-theory

**Global equivariant
Super homotopy theory**
and the C -field at singularities

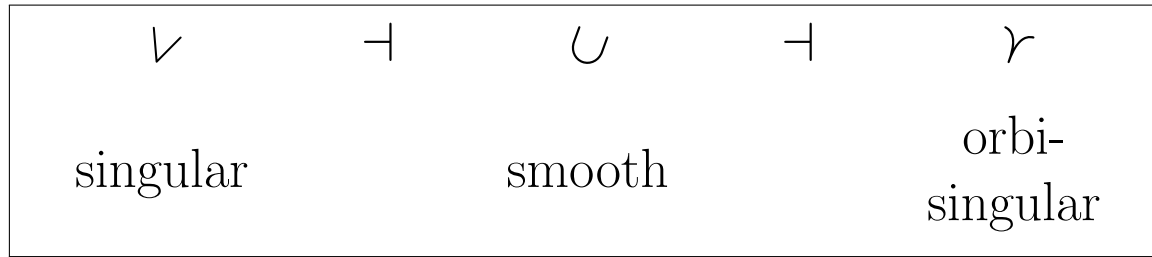
[back to Part I](#)

orbifolded

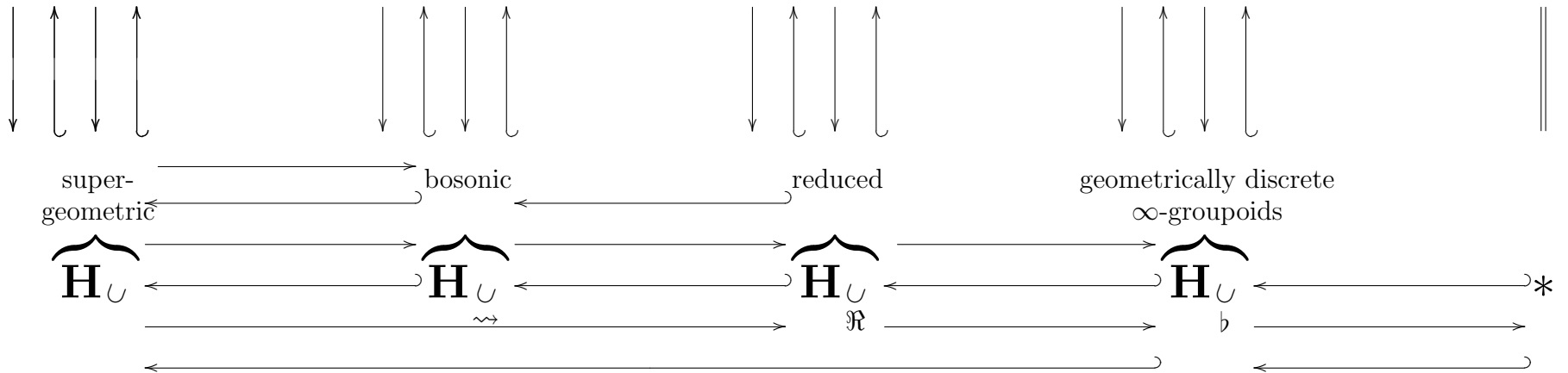
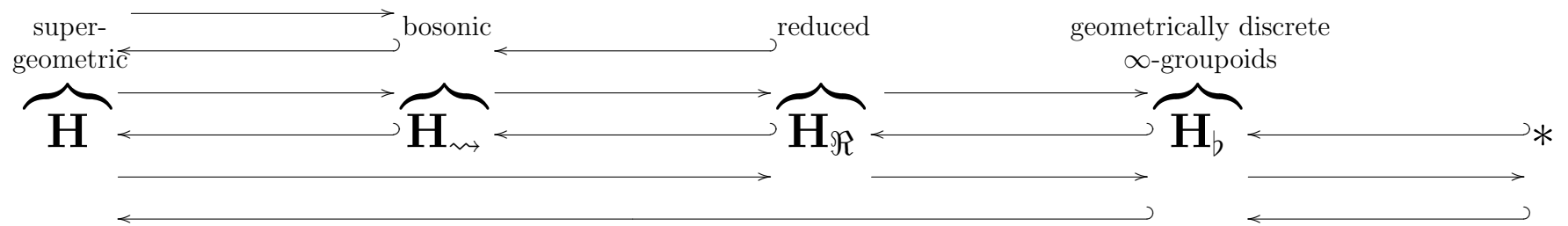


**Global equivariant
Super homotopy theory**
and the C -field at singularities

The modalities of global equivariant homotopy theory

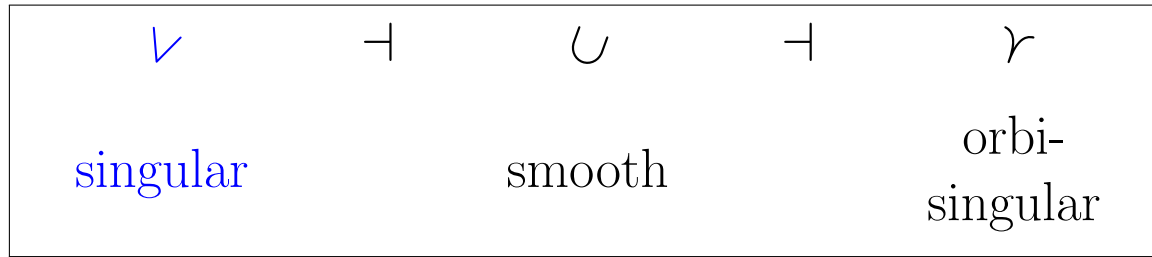


possibly singular/orbifolded

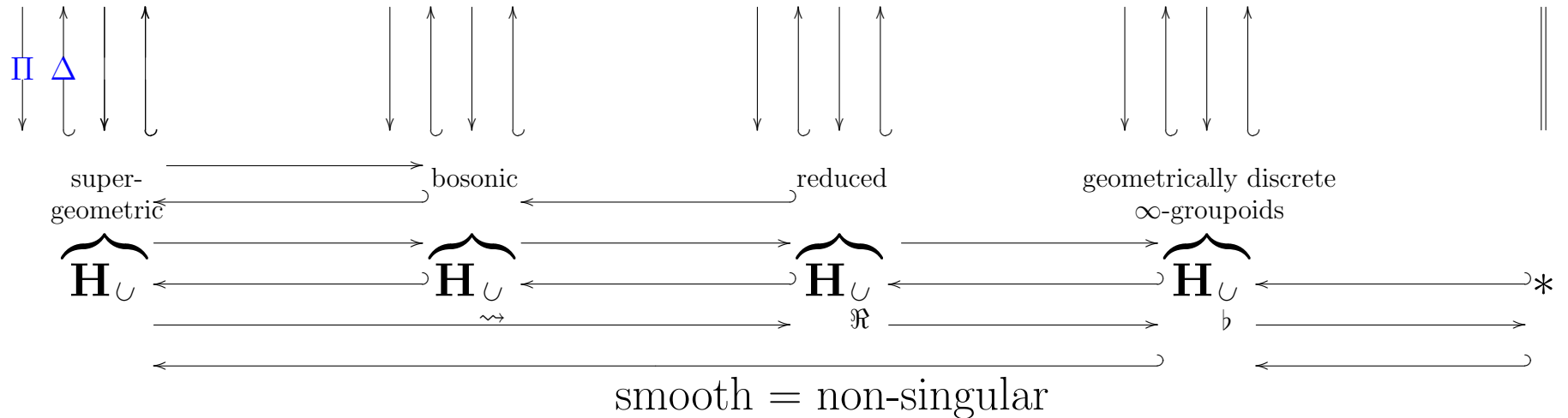
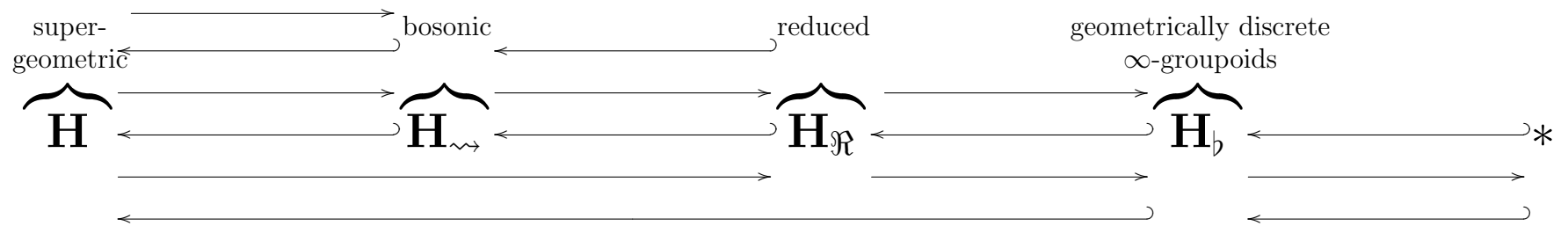


smooth = non-singular

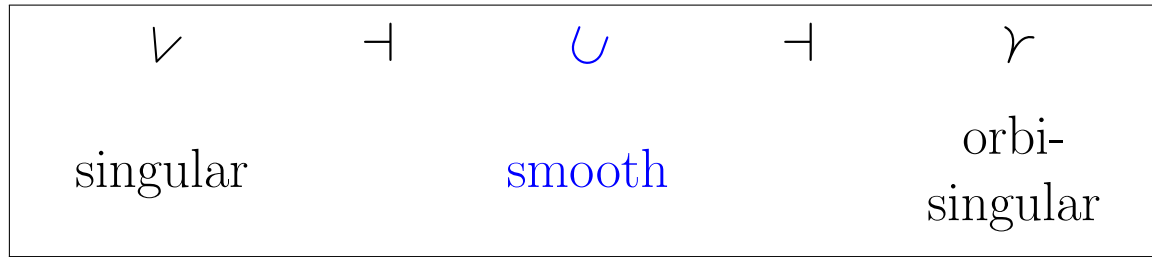
The modalities of global equivariant homotopy theory



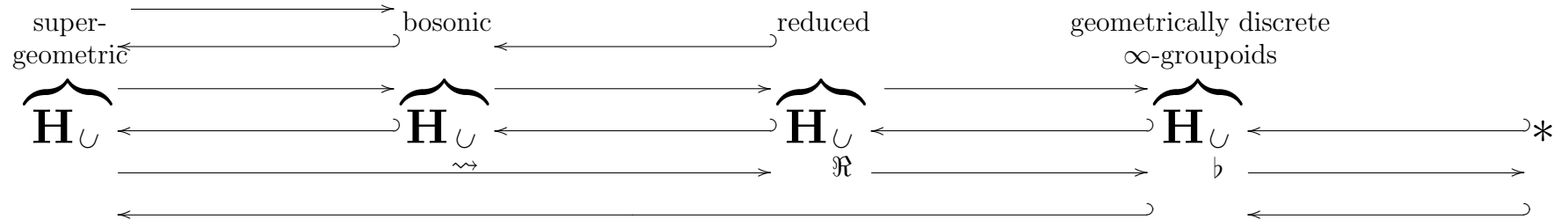
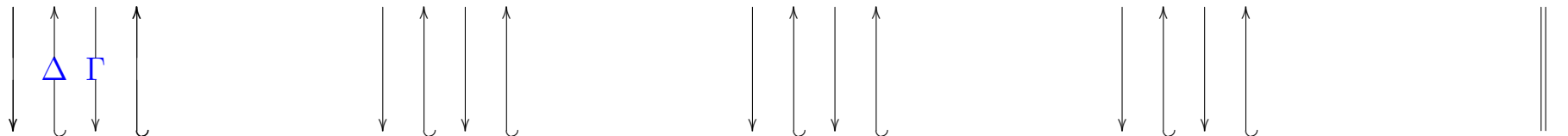
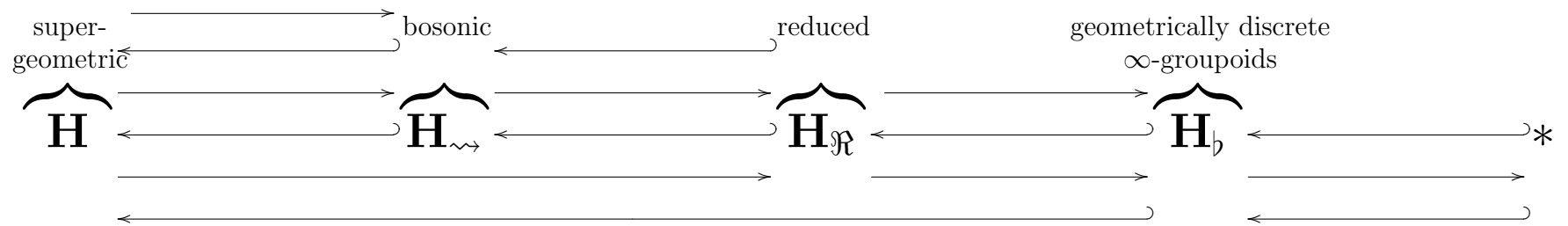
possibly singular/orbifolded



The modalities of global equivariant homotopy theory

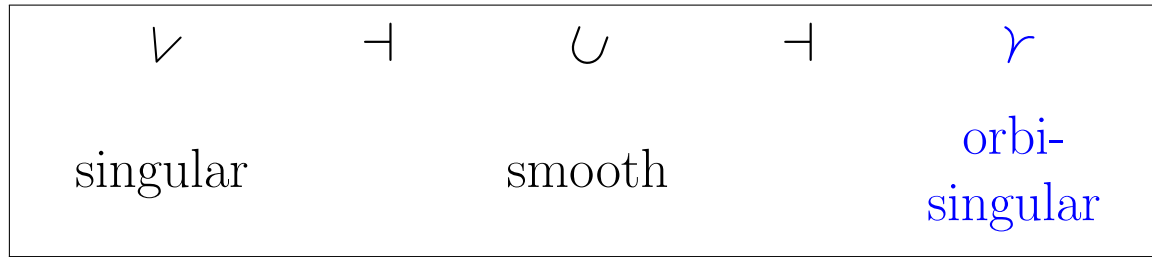


possibly singular/orbifolded

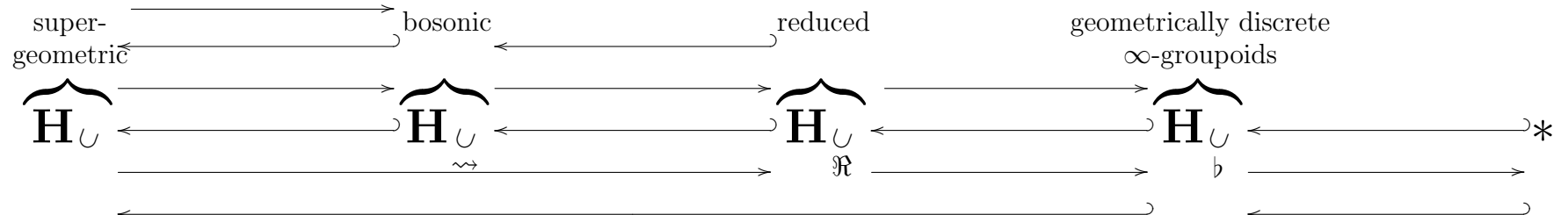
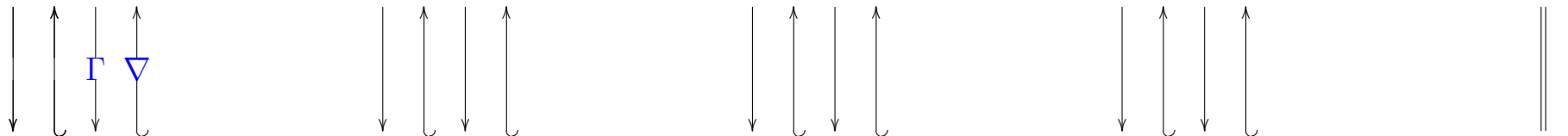
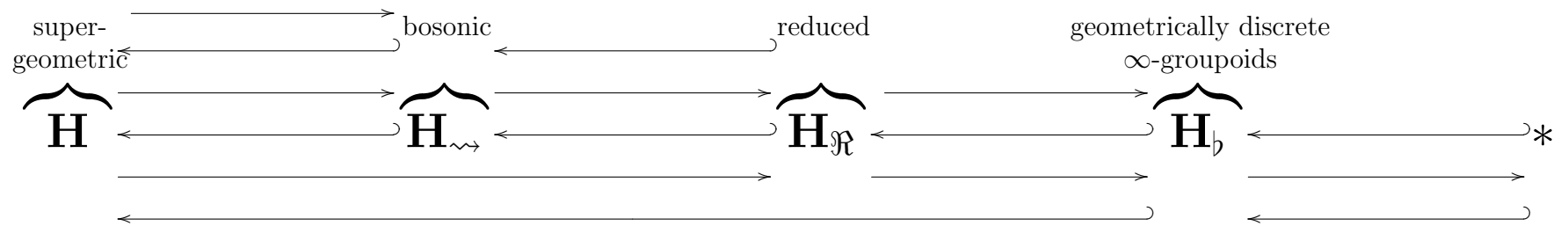


smooth = non-singular

The modalities of global equivariant homotopy theory



possibly singular/orbifolded



smooth = non-singular

Super-Orbifolds – Abstract definition

Let $\underbrace{V}_{\substack{\text{tangent} \\ \text{space} \\ \text{model}}}$, $\underbrace{G}_{\substack{\text{generic} \\ \text{singularity} \\ \text{type}}}$ $\in \text{Grp}(\mathbf{H})$ be group objects.

Definition. A G -orbi V -fold is

- an object $\mathcal{X} \in \mathbf{H}/\mathbf{B}G_\gamma$

which is

1. 0-truncated: $\tau_0(\mathcal{X}) \simeq \mathcal{X}$
2. orbi-singular: $\gamma(\mathcal{X}) \simeq \mathcal{X}$

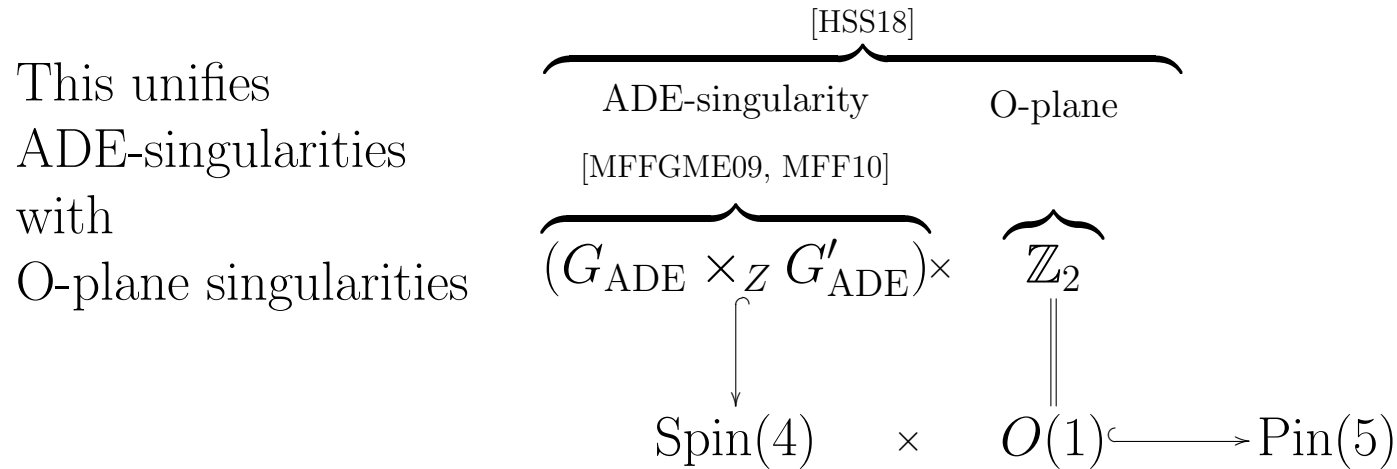
3. a V -fold: there exists a V -atlas

$$\begin{array}{ccc} & U & \\ p_V \swarrow & & \searrow p_X \\ V & & \mathcal{X}_\cup \end{array}$$

- (a) p_X is a covering: $(\tau_{-1})_{/X}(p_X) \simeq *$
- (b) p_X is a local diffeomorphism: $\mathfrak{S}_{/X}(p_X) \simeq p_X$
- (c) p_V is a local diffeomorphism: $\mathfrak{S}_{/V}(p_V) \simeq p_V$

The global equivariant 4-sphere

In the following $G := \text{Pin}(5)^b$
 the unoriented spin group in 5d, regarded as geometrically discrete.



Write $\mathbf{S}^4 \in \text{SmoothManifolds}$ for the smooth 4-sphere.

$\hookrightarrow \mathbf{H}$

with $S^4 := \int (\mathbf{S}^4) \in \infty\text{Groupoids}$ its shape.

Then

$\mathbf{S}_r^4 := \left(\mathbf{S}^4 // \flat\text{Pin}(5) \right)_r \in \mathbf{H}_{/\mathbf{B}\text{Pin}(5)^b_r}$ is a $\text{Pin}(5)^b$ -orbi \mathbb{R}^4 -fold

$S_r^4 := \int \left(\mathbf{S}^4 // \text{Pin}(5)^b \right)_r$ is its shape orbi-space

Equivariant Cohomotopy of Super-orbifolds

Let

$$\mathbb{R}^{10,1|\mathbf{32}} \in \text{Grp}(\mathbf{H}) \quad D = 1, \mathcal{N} = 1 \text{ translational supersymmetry}$$

$$\mathcal{X} \in \mathbf{H}_{/\mathbf{BPin}(5)_r^b} \quad \text{a } \text{Pin}(5)^b\text{-orbi } \mathbb{R}^{10,1|\mathbf{32}}\text{-fold.}$$

Definition.

The cocycle space of *equivariant Cohomotopy* of \mathcal{X} is

$$\mathbf{H}_{/\mathbf{BPin}(5)_r^b}(\mathcal{X}, S_r^4) = \left\{ \begin{array}{ccc} & \text{cocycle in} & \\ & \text{equivariant Cohomotopy} & \\ \mathcal{X} & \xrightarrow{\quad\quad\quad} & S_r^4 \\ & \swarrow \quad \searrow & \\ & \mathbf{BPin}(5)_r^b & \end{array} \right\}$$

and so the cohomology set is

$$H(\mathcal{X}, S_r^4) := \pi_0 \mathbf{H}_{/\mathbf{BPin}(5)_r^b}(\mathcal{X}, S_r^4)$$

Differential Equivariant Cohomotopy of Super-Orbifolds

Definition. $\Omega_{\text{flat}}(-, \mathfrak{L}S_r^4) \in \mathbf{H}/\mathbb{B}\text{bPin}(5)$

is the universal moduli space of

$$\mathbb{R}^{d|N} \times \mathbb{D} \times \left(\begin{array}{c} \mathbb{B}K \\ \downarrow \\ \mathbb{B}\text{bPin}(5) \end{array} \right) \mapsto \overbrace{\Omega_{\text{flat}}\left(\mathbb{R}^{d|N} \times \mathbb{D}, \mathfrak{L}\underbrace{(S^4)^K}_{\substack{\text{fixed point sphere} \\ \underbrace{\hspace{2cm}} \\ L_\infty\text{-algebra dual to its} \\ \text{minimal Sullivan model}}}\right)}^{\text{flat } L_\infty\text{-algebra valued} \\ \text{super-differential forms}}$$

Claim: $\int\left(\Omega_{\text{flat}}(-, \mathfrak{L}S_r^4)\right) \simeq (S_r^4)_{\mathbb{R}}$

Definition. The *differential equivariant 4-sphere* is

$$\widehat{S}_r^4 := S_r^4 \times_{(S_r^4)_{\mathbb{R}}} \Omega_{\text{flat}}(-, \mathfrak{L}S_r^4)$$

Hence *differential equivariant Cohomotopy in degree 4* is

$$H(\mathcal{X}, \widehat{S}_r^4) := \pi_0 \mathbf{H}/_{\mathbb{B}\text{Pin}(5)_r^b}(\mathcal{X}, \widehat{S}_r^4)$$

M-brane charge quantized C -field

C -field	electromagnetic field ("A-field")
<p style="text-align: center;">flux forms</p> $\mathcal{X} \xrightarrow{(G_4, G_7)} \Omega_{\text{flat}}(-, \mathfrak{L}S_r^4)$	<p style="text-align: center;">Faraday tensor</p> $X \xrightarrow{F_2} \underbrace{\Omega_{\text{flat}}(-, \mathfrak{L}BU(1))}_{=\Omega_{\text{clsd}}^2(-)}$
<p style="text-align: center;"><i>M-brane charge quantized C-field</i></p> $\begin{array}{ccc} & \widehat{S}_r^4 & \\ & \nearrow \hat{C} & \\ \mathcal{X} & \xrightarrow{(G_4, G_7)} & \Omega_{\text{flat}}(-, \mathfrak{L}S_r^4) \\ & & \downarrow \end{array}$	<p style="text-align: center;">Dirac charge quantized electromagnetic field</p> $\begin{array}{ccc} & \mathbf{BU}(1)_{\text{conn}} & \\ & \nearrow \hat{A} & \\ X & \xrightarrow{F_2} & \Omega_{\text{clsd}}^2(-) \\ & & \downarrow \end{array}$

Super Cartan geometry

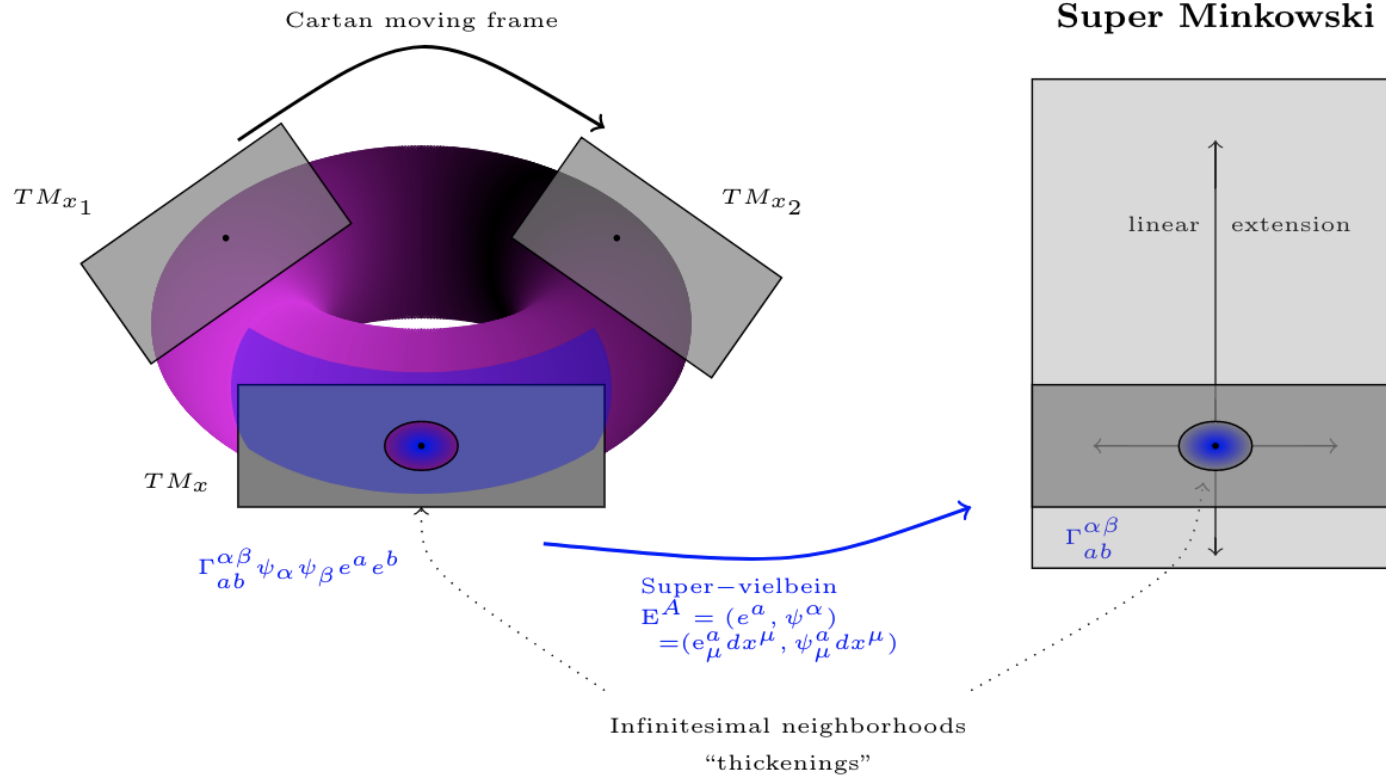
and 11d orbifold supergravity

[back to Part I](#)

Cartan geometry formalizes Einstein principle of equivalence

Spacetime is locally equivalent to Minkowski spacetime, namely in the infinitesimal neighbourhood of every point

We now generalize this from manifolds to super-orbifolds...



G -Structures on orbi V -folds ([Wellen17, Sch13])

Def.: infinitesimal disk around origin: $\mathbb{D}^V := V \times_{\mathfrak{S}(V)} \{e\} \hookrightarrow V$

Prop.: every orbi V -fold \mathcal{X} carries its canonical V -frame bundle $\mathcal{X}_\cup \xrightarrow{\text{frame}} \mathbf{BAut}(\mathbb{D}^V)$

Def.: for $G \xrightarrow{\text{homom.}} \mathbf{Aut}(\mathbb{D}^V)$ a G -structure is a lift (E is the *vielbein*)

$$\begin{array}{ccc}
 \mathcal{X}_\cup & \overset{\text{frame}}{\dashrightarrow} & \mathbf{BG} \\
 & \swarrow \text{frame} \quad \nwarrow E & \\
 & \mathbf{BAut}(\mathbb{D}^V) &
 \end{array}$$

Prop.: V itself carries canonical G -structure given by left translation

$$\begin{array}{ccc}
 V & \overset{\text{frame}}{\dashrightarrow} & \mathbf{BG} \\
 & \swarrow \text{frame} \quad \nwarrow E_{\text{li}} & \\
 & \mathbf{BAut}(\mathbb{D}^V) &
 \end{array}$$

Def.: a G -structure is *torsion-free and flat* if it coincides with this canonical one on each infinitesimal disk $E|_{\mathbb{D}_x^V} \simeq (E_{\text{li}})|_{\mathbb{D}_e^V}$

11d Supergravity from Super homotopy theory

Consider now $V = \mathbb{R}^{10,1|32}$ and \mathcal{X} an orbi $\mathbb{R}^{10,1|32}$ -fold.

Claim:

$$G := \mathbf{Aut}_{\text{Grp}}^{\rightsquigarrow}(\mathbb{R}^{10,1|32}) \simeq \text{Spin}(10, 1)$$

$$\begin{aligned} G\text{-structure on } \mathcal{X} &\simeq \text{super-vielbein on } \mathcal{X} \\ &\simeq \text{metric/field of gravity} \end{aligned}$$

$$\begin{aligned} G\text{-structure is torsion-free:} &\Leftrightarrow \text{super-torsion on } \mathcal{X} \text{ vanishes} \\ &\Leftrightarrow \mathcal{X} \text{ is solution to 11d supergravity} \\ &\quad \begin{matrix} \text{[CaLe93]} \\ \text{[How97]} \end{matrix} \text{ with vanishing bosonic flux} \end{aligned}$$

$$G\text{-structure is flat:} \quad \Leftrightarrow \mathcal{X} \text{ is a "flat" super-orbifold solution to 11d supergravity}$$

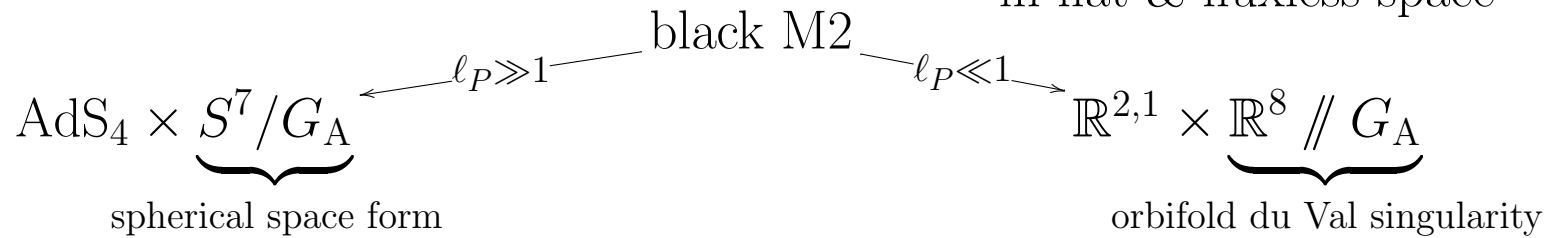
\Rightarrow all $\left\{ \begin{array}{l} \text{curvature} \\ \& G_4\text{-flux} \end{array} \right\}$ hence all $\left\{ \begin{array}{l} \text{higher curvature corrections} \\ \& \text{flux quantization} \end{array} \right\}$
 crammed into orbifold singularities
 and thus *taken care of by the equivariance*
 of charge quantization in differential equivariant Cohomotopy

Flat & fluxless except at curvature- & flux- singularities

Plausibility check:

Black M2/M5-brane solutions to SuGra interpolate ([AFFHS98]) between:

Planck-scale curved throat $\xleftrightarrow{\text{near/far horizon geometry}}$ orbifold-singularity
in flat & fluxless space



inconsistent:

Planck-scale throat ($\ell_P \gg 1$)

spurious in SuGra ($\ell_P \ll 1$)

(evaded only by

macroscopic $N \gg 1$)

consistent:

all Planck-scale geometry

crammed into orbi-singularity

(necessary for

microscopic $N = 1$)

Hence, indeed, a consistent & complete picture:

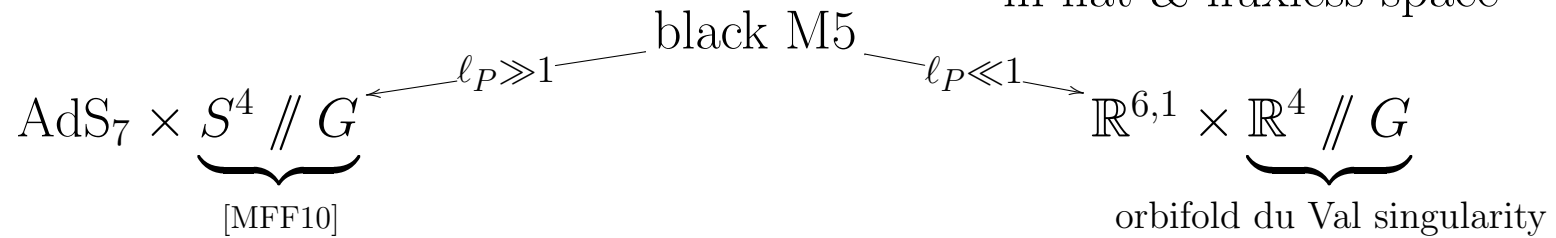
1. is flat & fluxless away from singularities,
2. has hidden degrees of freedom inside the singularities.

Flat & fluxless except at curvature- & flux- singularities

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Hence, indeed, a consistent & complete picture:

1. is flat & fluxless away from singularities,
2. has hidden degrees of freedom inside the singularities.

In conclusion

the following enhancement of 11d supergravity
naturally emerges out of super homotopy theory

The full covariant phase space is...

CovariantPhaseSpace :=

$$\bigsqcup_{\substack{[\mathcal{X}] \in \\ \text{Pin}(5)^b \text{Orbi-} \\ \mathbb{R}^{10,1|32} \text{Folds}(\mathbf{H}) \sim}} \left(\begin{array}{ccc} & \text{CovariantPhaseSpace}_{\mathcal{X}} & \\ & \swarrow & \searrow \\ \underline{\mathbf{H}}\left(\mathcal{X}_{\cup}, \mathbf{BSpin}(10, 1)\right)^{\tau=0} & & \underline{\mathbf{H}}\left(\mathcal{X}_{\cap}, \widehat{S^4}_{\cap}\right) / \mathbf{BPin}(5)^b_{\cap} \\ & \swarrow & \searrow \\ & \underline{\mathbf{H}}\left(\mathcal{X}_{\cup}, \mathbf{\Omega}(-, \mathbb{S}^4)\right) & \end{array} \right)$$

$$\begin{aligned} \text{Observables} &:= \widehat{H}\mathbb{R}_{\cap} \left(\text{CovariantPhaseSpace} \right) \\ &= \widehat{H}\mathbb{R}_{\cap} \left(\sum_{\text{Pin}(5)^b}^{\infty} \text{CovariantPhaseSpace} \right) \end{aligned}$$

In conclusion

the following enhancement of 11d supergravity naturally emerges out of super homotopy theory

... for each class of super-orbifolds \mathcal{X} ...

CovariantPhaseSpace :=

$$\bigsqcup_{\substack{[\mathcal{X}] \in \\ \text{Pin}(5)^b \text{Orbi-} \\ \mathbb{R}^{10,1|32} \text{Folds}(\mathbf{H}) \sim}} \left(\begin{array}{ccc} & \text{CovariantPhaseSpace}_{\mathcal{X}} & \\ & \swarrow & \searrow \\ \underline{\mathbf{H}}\left(\mathcal{X}_{\cup}, \mathbf{BSpin}(10, 1)\right)^{\tau=0} & & \underline{\mathbf{H}}\left(\mathcal{X}_{\cap}, \widehat{S^4}_{\cap}\right) / \mathbf{BPin}(5)^b_{\cap} \\ & \searrow & \swarrow \\ & \underline{\mathbf{H}}\left(\mathcal{X}_{\cup}, \mathbf{\Omega}(-, \mathbb{S}^4)\right) & \end{array} \right)$$

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In conclusion

the following enhancement of 11d supergravity
naturally emerges out of super homotopy theory

... a super-torsion-free Spin-structure
encoding the fields of supergravity...

CovariantPhaseSpace :=

$$\bigsqcup_{\substack{[\mathcal{X}] \in \\ \text{Pin}(5)^b \text{Orbi-} \\ \mathbb{R}^{10,1|32} \text{Folds}(\mathbf{H}) \sim}} \left(\begin{array}{ccc} & \text{CovariantPhaseSpace}_{\mathcal{X}} & \\ & \swarrow & \searrow \\ \underline{\mathbf{H}}\left(\mathcal{X}_{\cup}, \mathbf{BSpin}(10, 1)\right)^{\tau=0} & & \underline{\mathbf{H}}\left(\mathcal{X}_{\cap}, \widehat{S^4}_{\cap}\right) / \mathbf{BPin}(5)^b_{\cap} \\ & \searrow & \swarrow \\ & \underline{\mathbf{H}}\left(\mathcal{X}_{\cup}, \mathbf{\Omega}(-, \mathbb{S}^4)\right) & \end{array} \right)$$

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In conclusion

the following enhancement of 11d supergravity naturally emerges out of super homotopy theory

... equipped with a compatible lift of the flux forms to a cocycle in differential equivariant Cohomotopy (charge quantization).

CovariantPhaseSpace :=

$$\bigsqcup_{\substack{[\mathcal{X}] \in \\ \text{Pin}(5)^b \text{Orbi-} \\ \mathbb{R}^{10,1|32} \text{Folds}(\mathbf{H}) \sim}} \left(\begin{array}{ccc} & \text{CovariantPhaseSpace}_{\mathcal{X}} & \\ & \swarrow & \searrow \\ \underline{\mathbf{H}}\left(\mathcal{X}_{\cup}, \mathbf{BSpin}(10, 1)\right)^{\tau=0} & & \underline{\mathbf{H}}\left(\mathcal{X}_r, \widehat{S}_r^4\right) / \mathbf{BPin}(5)^b_r \\ & \swarrow \text{(pb)} & \searrow \\ & \underline{\mathbf{H}}\left(\mathcal{X}_{\cup}, \Omega(-, \mathbb{S}^4)\right) & \end{array} \right)$$

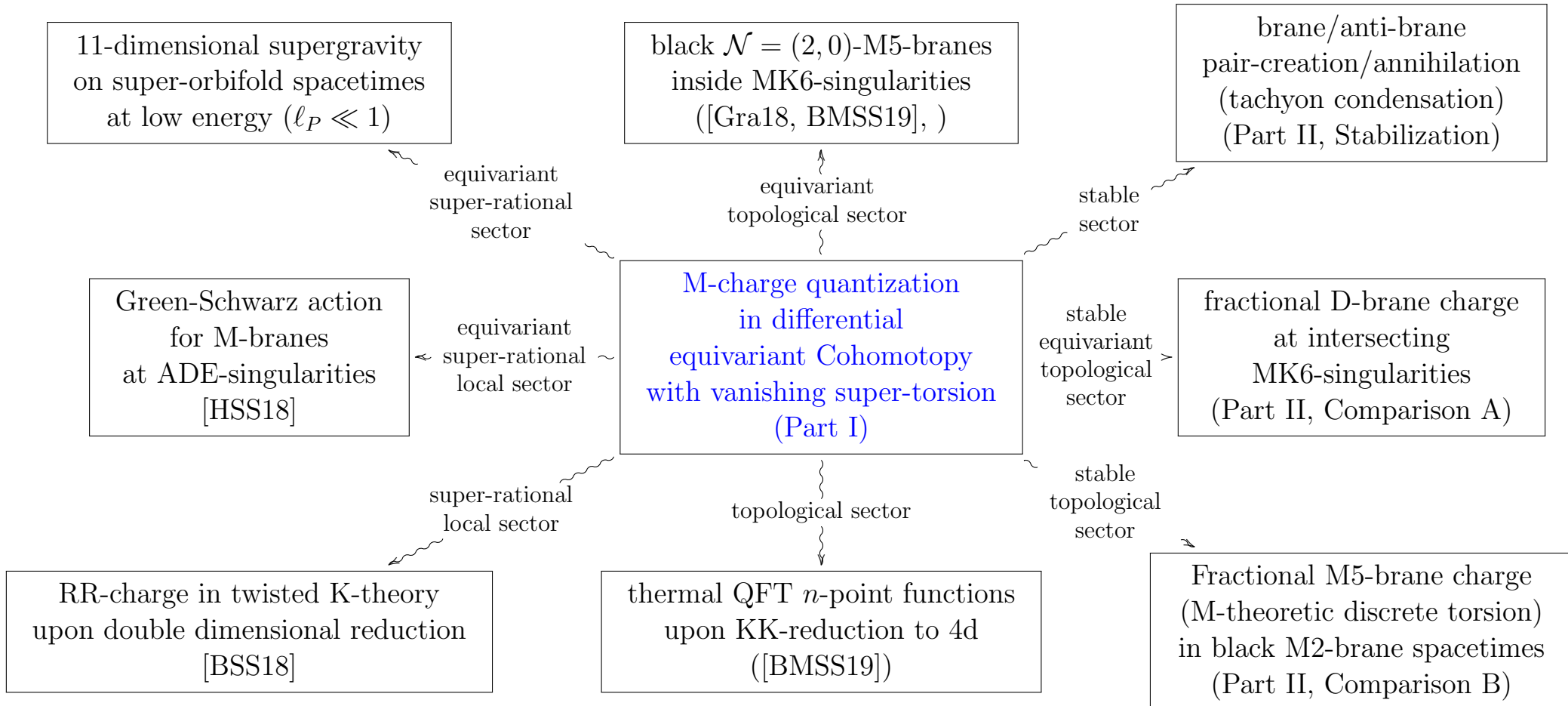
$$\begin{aligned} \text{Observables} &:= \widehat{H}\mathbb{R}_r\left(\text{CovariantPhaseSpace}\right) \\ &= \widehat{H}\mathbb{R}_r\left(\sum_{\text{Pin}(5)^b}^{\infty} \text{CovariantPhaseSpace}\right) \end{aligned}$$

Part II.

Some corners of M-theory

Part II.

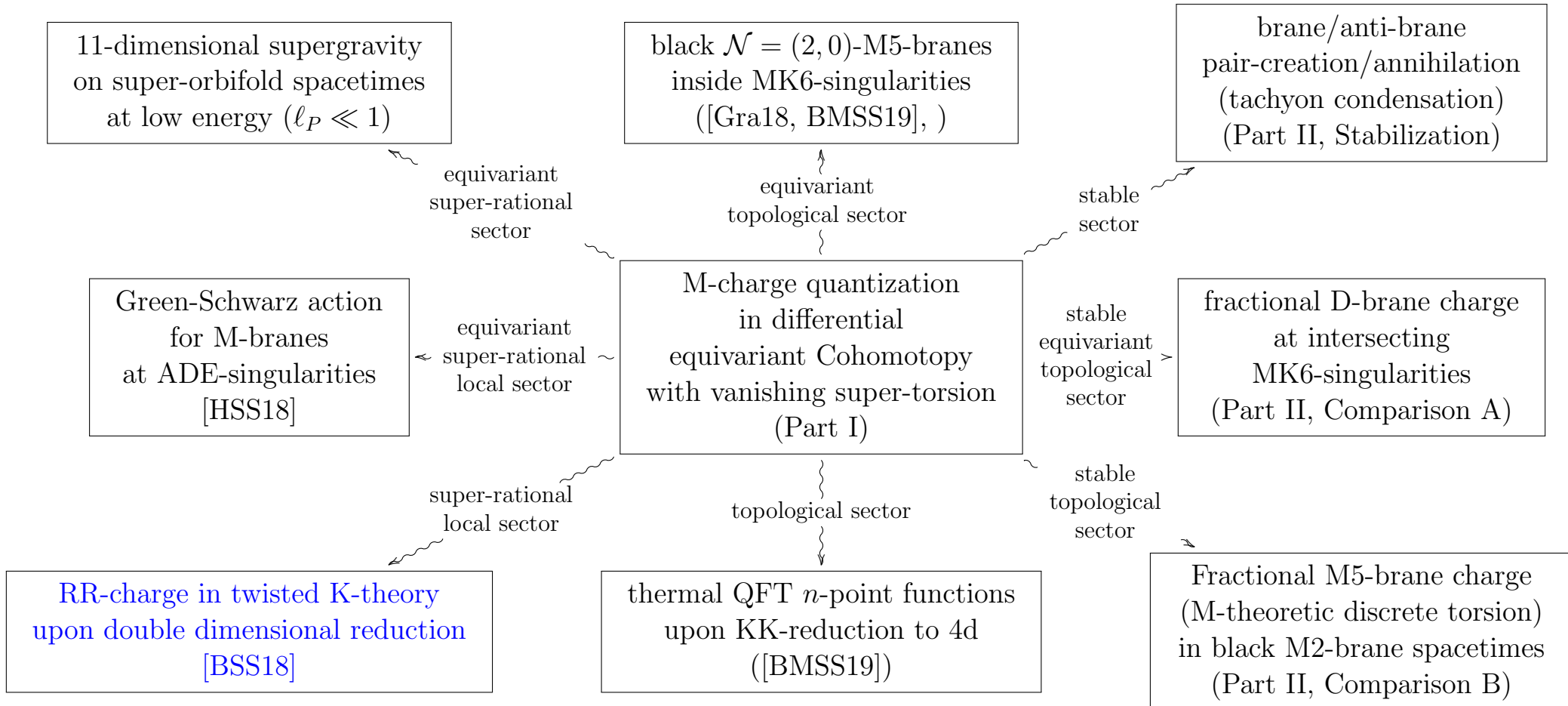
Some corners of M-theory



[back to Contents](#)

Part II.

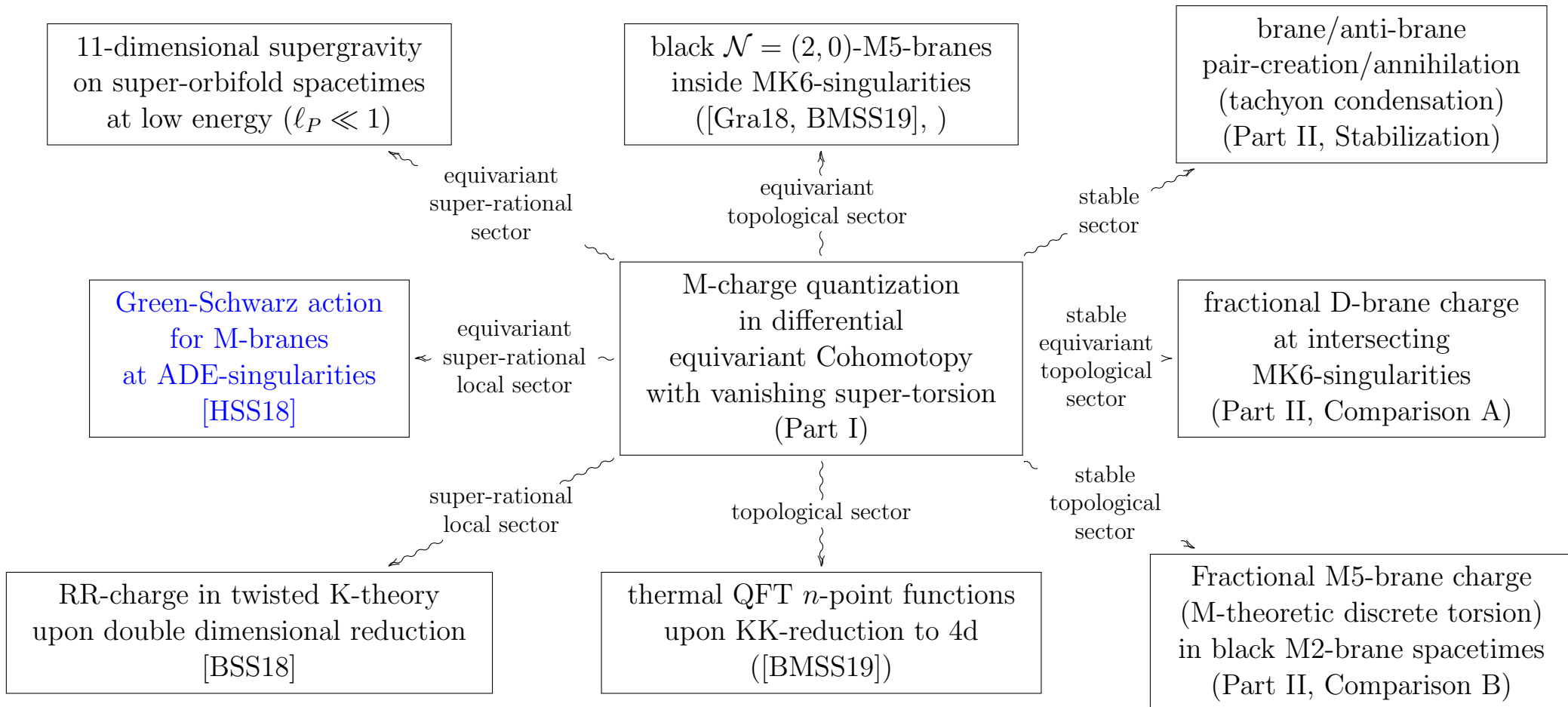
Some corners of M-theory



[back to Contents](#)

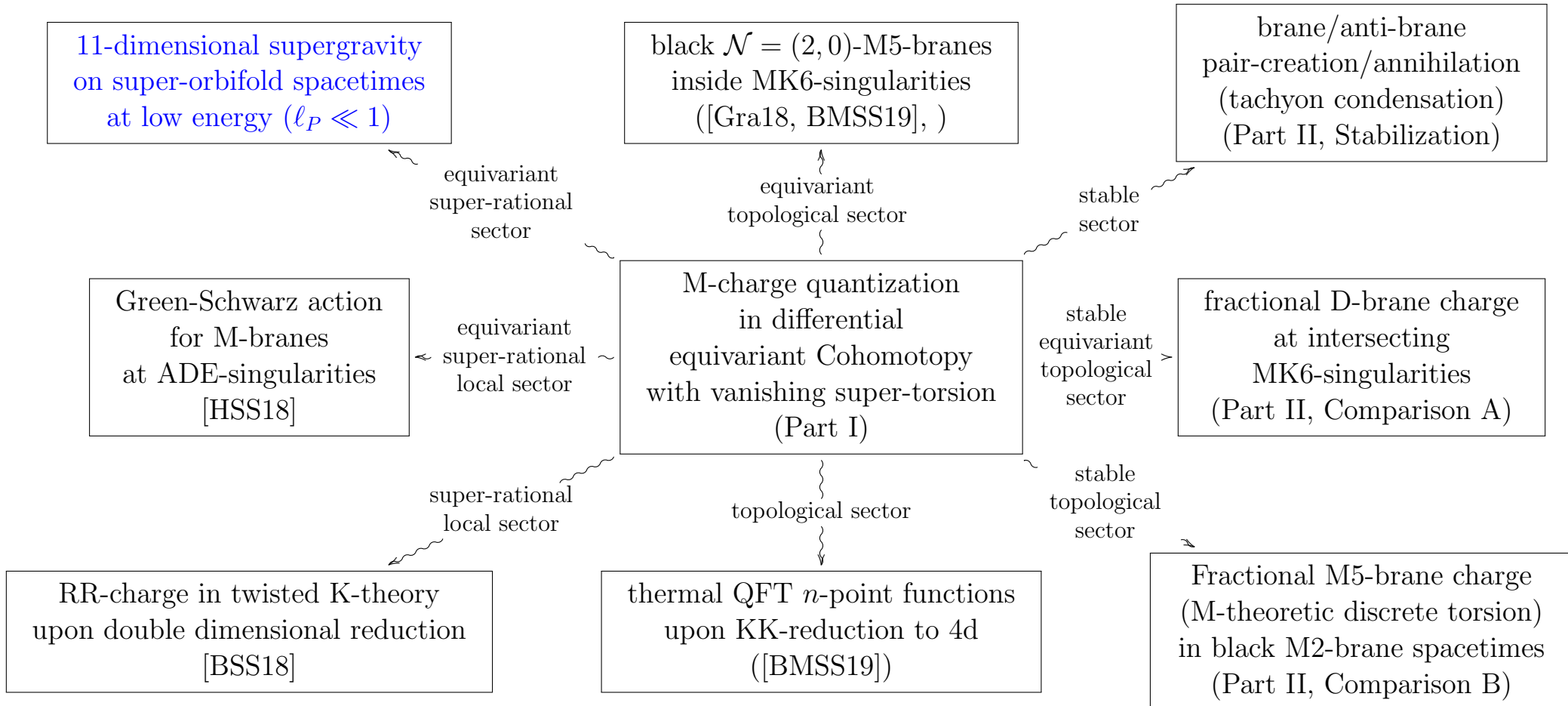
Part II.

Some corners of M-theory



Part II.

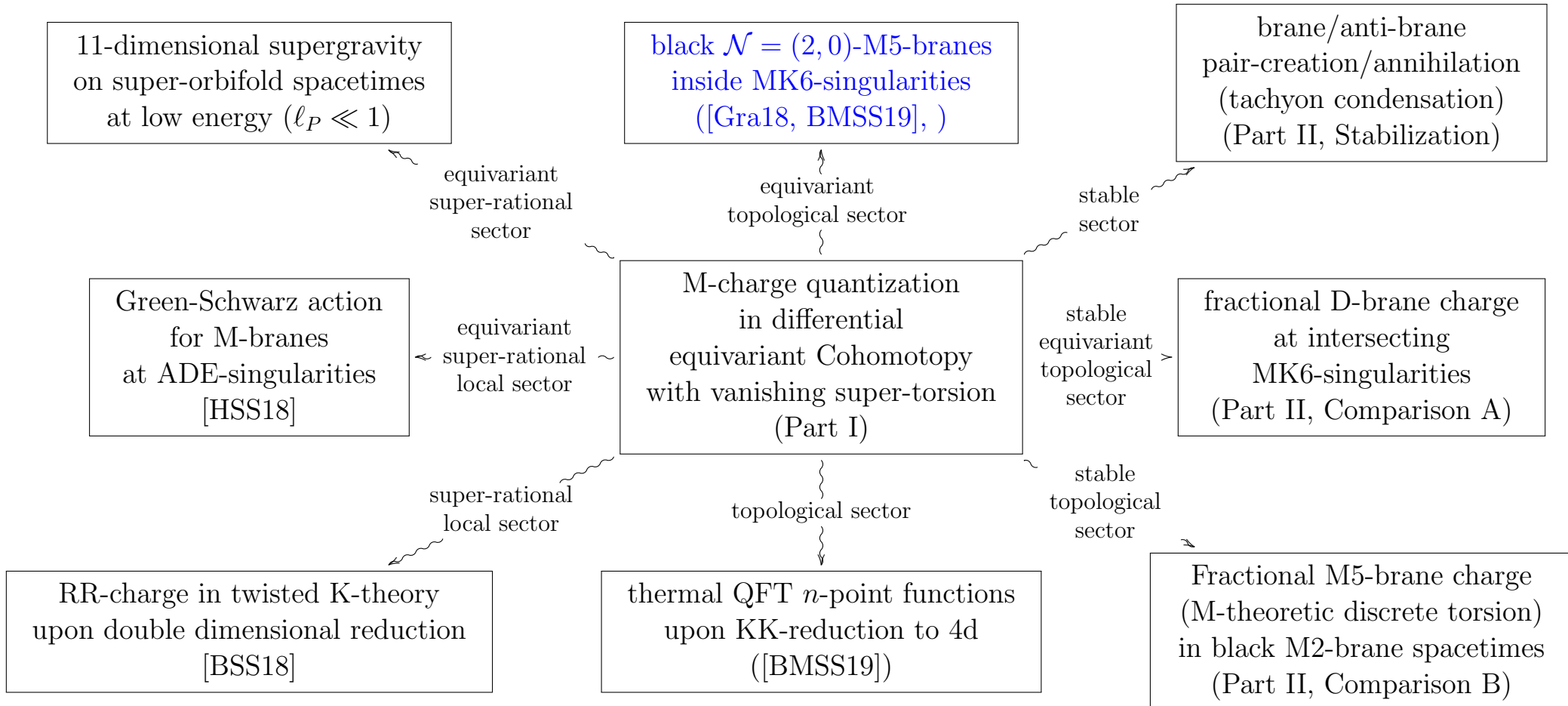
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[back to Contents](#)

Part II.

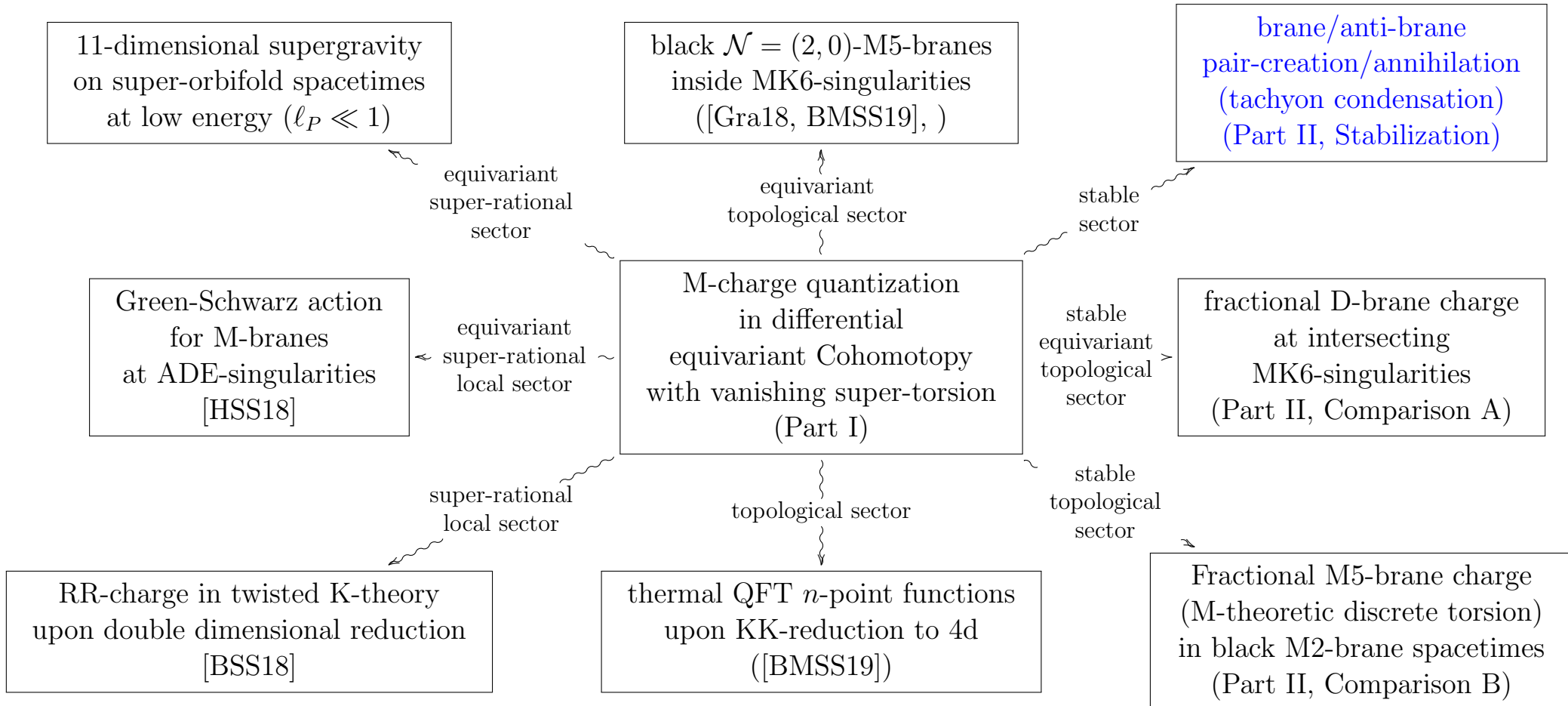
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[back to Contents](#)

Part II.

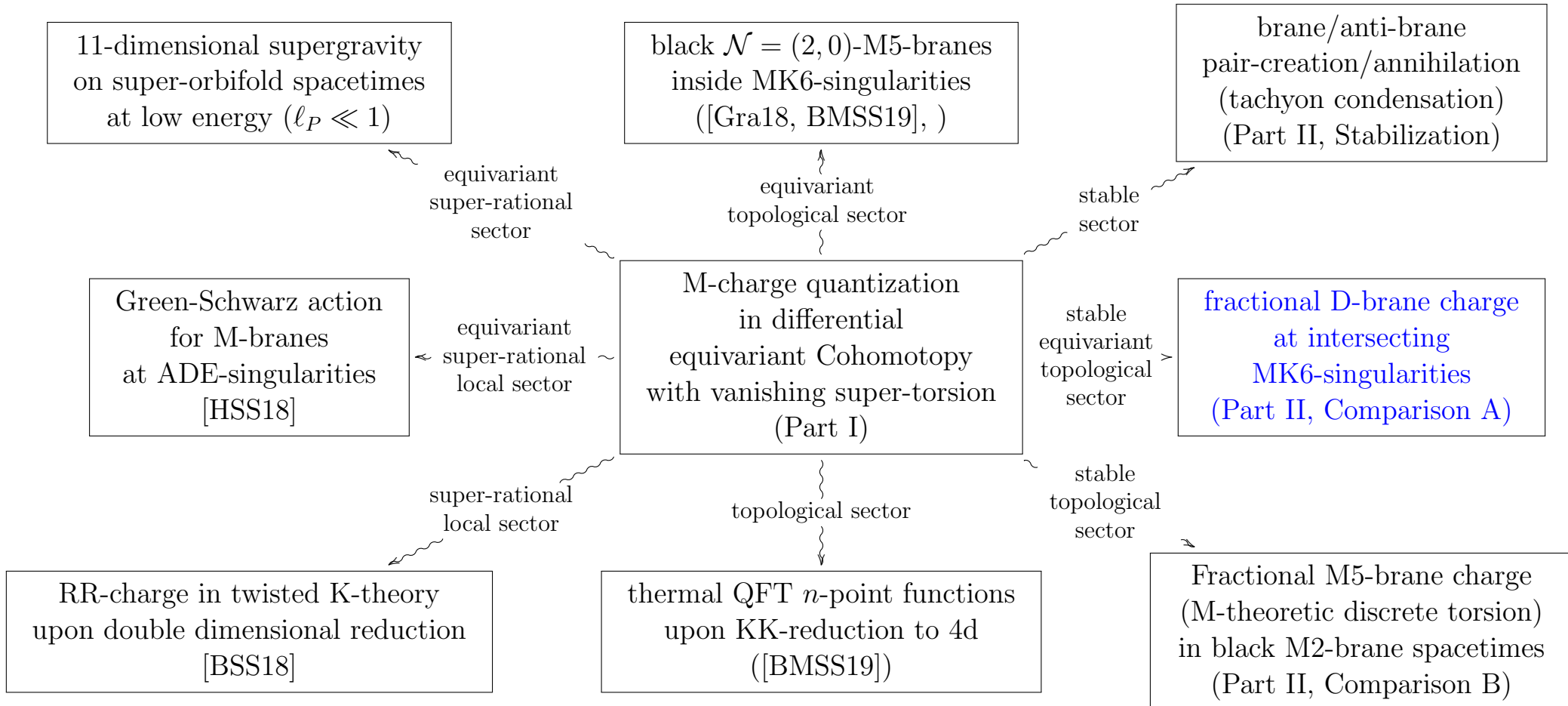
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[back to Contents](#)

Part II.

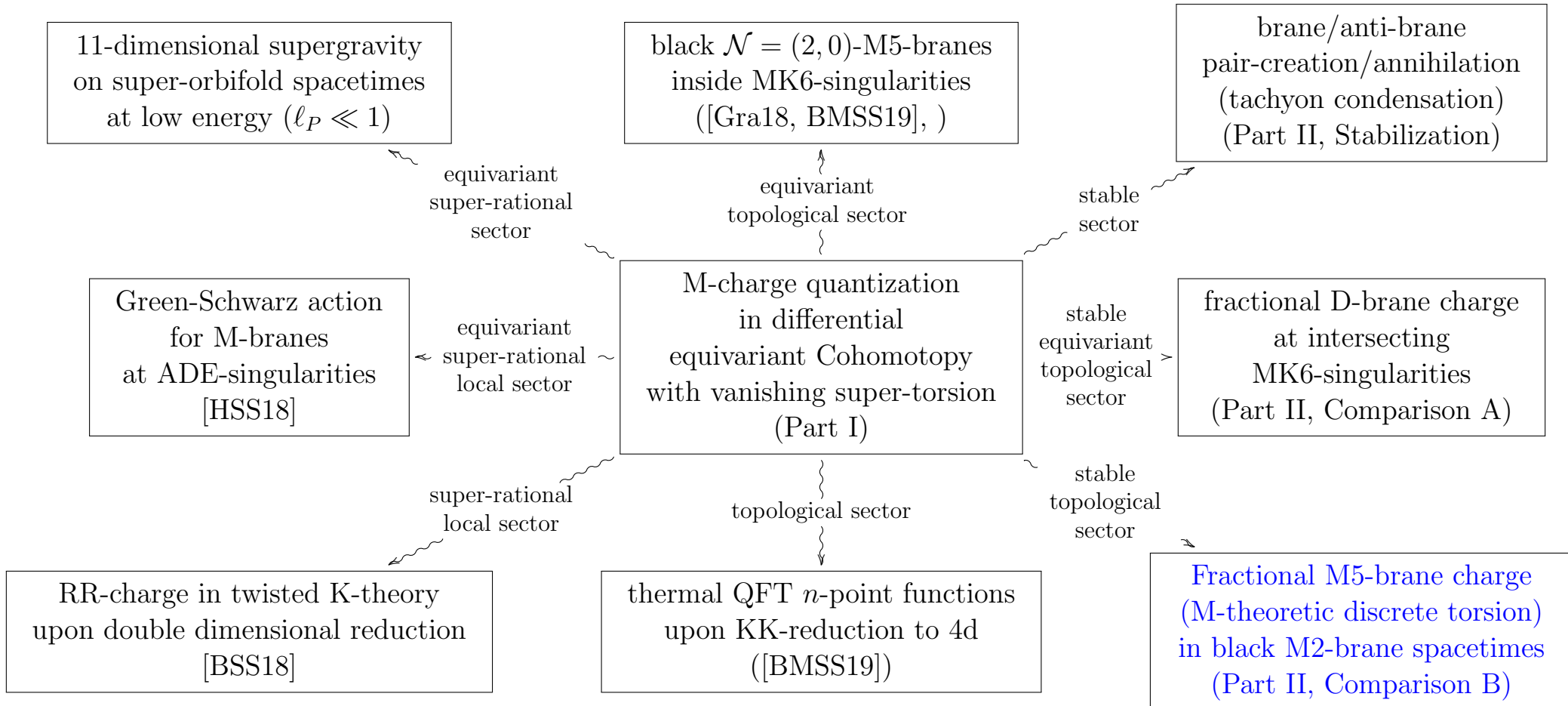
Some corners of M-theory



[back to Contents](#)

Part II.

Some corners of M-theory



[back to Contents](#)

Stable homotopy theory

and anti-branes

[back to Part II](#)

Stable Cohomotopy from observing Cohomotopy

Any space of observables is *linear*:

Observables may be added and subtracted.

Linear + Homotopy theory = *Stable homotopy theory*

$$\underbrace{\widehat{H}\mathbb{R}_G(\text{CovariantPhaseSpace})}_{\substack{\text{homotopy-linear} \\ \text{space of observables}}} = \widehat{H}\mathbb{R}_G\left(\underbrace{\Sigma_G^\infty \text{CovariantPhaseSpace}}_{\substack{\text{free homotopy-linearized} \\ \text{covariant phase space}}}\right)$$

There is canonical comparison map

$$\underbrace{\Sigma^\infty \mathbf{H}/\mathbf{B}\text{Pin}(5)_r^b(\mathcal{X}, S_r^4)}_{\text{Cohomotopy}} \xrightarrow{\text{homotopy-linear approximation}} \underbrace{\text{Stab}(\mathbf{H})_{\text{Pin}_r^b}\left(\Sigma_G^\infty \mathcal{X}, \overset{\substack{\text{4-shifted} \\ \text{equivariant} \\ \text{sphere spectrum}}}{S_r^4}\right)}_{\text{stable Cohomotopy}}$$

to fields in *stable equivariant Cohomotopy*

We now discuss what this means...

Brane charge – 1st order approximation

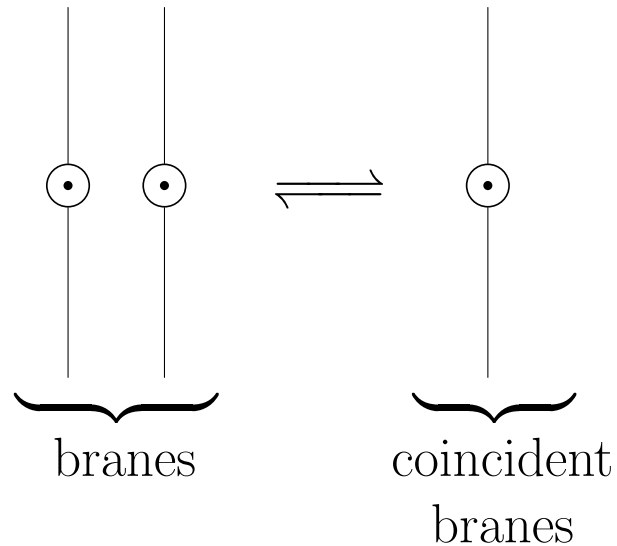
Let $(C, +)$ be an abelian semigroup
a “commutative monoid” of charges.

charge $\in (C, +)$

c_1 c_2

$c_1 + c_2$

singular locus



Fundamental example: the natural numbers

$$(C, +) = (\mathbb{N}, +)$$

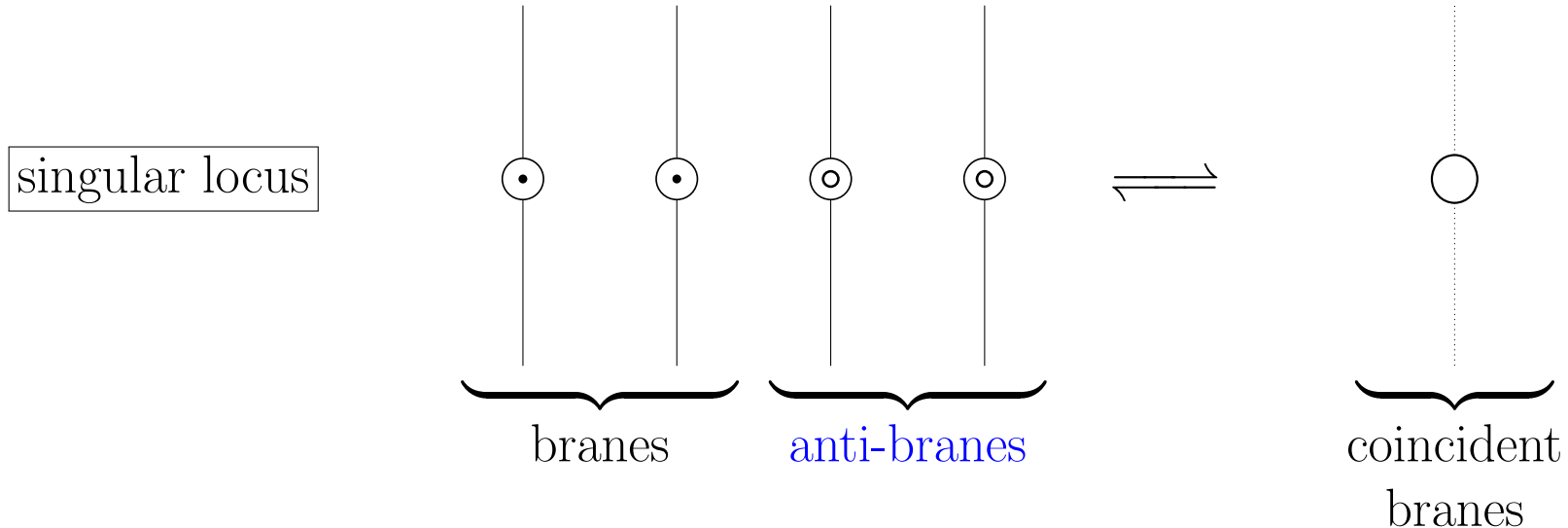
charge c = number of coincident branes \Leftrightarrow each brane carries unit charge

Brane charge – 2nd order approximation

Including **anti-brane charges**, hence negative brane charges, means to pass to the **abelian group completion** of the charge monoid:

$$K(C, +) := \left\{ (c^+, c^-) \mid c^\pm \in C \right\} / \left((c, c) \sim 0 \right)$$

charge $\in K(C, +)$	$(c_1, 0)$	$(c_2, 0)$	$(0, c_2)$	$(0, c_1)$	$(c_1 + c_2, c_1 + c_2) \sim 0$
----------------------	------------	------------	------------	------------	---------------------------------



Fundamental example: the integers:

$$K(\mathbb{N}, +) = (\mathbb{Z}, +)$$

$(c^+, c^-) =$ number of coincident branes
minus number of anti-branes

\Leftrightarrow

each brane carries unit charge
each anti-brane carries
negative unit charge

Brane charge – 3rd order approximation

The categorification of *commutative monoid* is *symmetric monoidal category*

$$(C, +) \qquad (C, \oplus)$$

$$(C, +) = \pi_0(C, \oplus)$$

- **Fundamental non-linear example**

Finite **pointed sets** with disjoint union $(C, \oplus) = (\text{Set}_{\text{fin}}^*/\sqcup)$

this categorifies the previous example: $\pi_0(\text{Set}_{\text{fin}}^*/\sqcup) \simeq (\mathbb{N}, +)$

- **Fundamental linear example** for \mathbb{F} a field

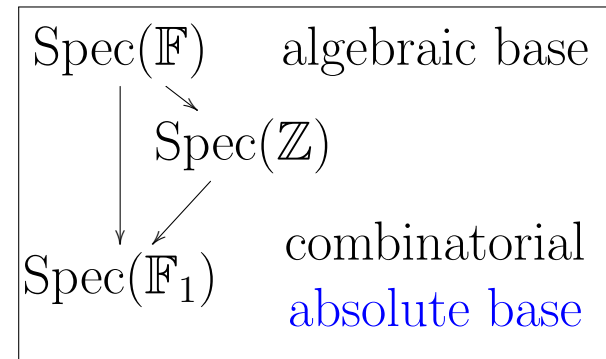
finite-dim **vector spaces** with direct sum $(C, \oplus) = (\mathbb{F}\text{Vect}_{\text{fin}}, \oplus)$

this *also* categorifies the previous example: $\pi_0(\mathbb{F}\text{Vect}_{\text{fin}}, \oplus) \simeq (\mathbb{N}, +)$

Unified on a deeper level:

pointed sets may be regarded as the vector spaces over the “absolute ground-field with one element” \mathbb{F}_1

$$(\text{Set}_{\text{fin}}^*/\sqcup) \simeq (\mathbb{F}_1\text{Vect}_{\text{fin}}, \oplus)$$



Brane charge in generalized cohomology

Brane/anti-brane annihilation may be **varying over spacetime** X

\rightsquigarrow enhance discrete abelian group of charges to a *space* of charges

The **homotopification** of *abelian group* is **∞ -loop space / spectrum**

$$(A, +) \qquad \mathcal{A}$$

$$(A, +) = \pi_0(\mathcal{A})$$

brane charge

locally constant	locally varying
$X \longrightarrow \underbrace{(A, +)}_{\substack{\text{discrete} \\ \text{abelian group}}}$	$X \longrightarrow \underbrace{\mathcal{A}}_{\substack{\infty\text{-loop space} \\ \text{or spectrum}}}$

Hence **brane charge group** on spacetime X is **generalized cohomology group**:

$$\mathcal{A}(X) := \pi_0 \text{Maps}(X, \mathcal{A})$$

Example: D-brane/anti-D-brane bound states

open string **tachyon condensation profile**:

$$X \longrightarrow \text{KU} \qquad (\text{conjecturally, or similar})$$

\simeq K-theory spectrum

Algebraic K-Theory – locally varying brane/anti-brane annihilation

in conclusion:

brane/anti-brane annihilation



abelian group completion
of charge monoid

locally varying brane charge



spectral enhancement
of charge group

combine:

The *categorification / homotopification* of *abelian group completion* is *algebraic K-theory spectrum*

$$K\left(\pi_0(\mathcal{C}, \oplus)\right)$$

$$\mathbb{K}(\mathcal{C}, \oplus) := \Omega B_{\oplus} \mathcal{C}$$

$$K\left(\pi_0(\mathcal{C}, \oplus)\right) = \pi_0\left(\mathbb{K}(\mathcal{C}, \oplus)\right)$$

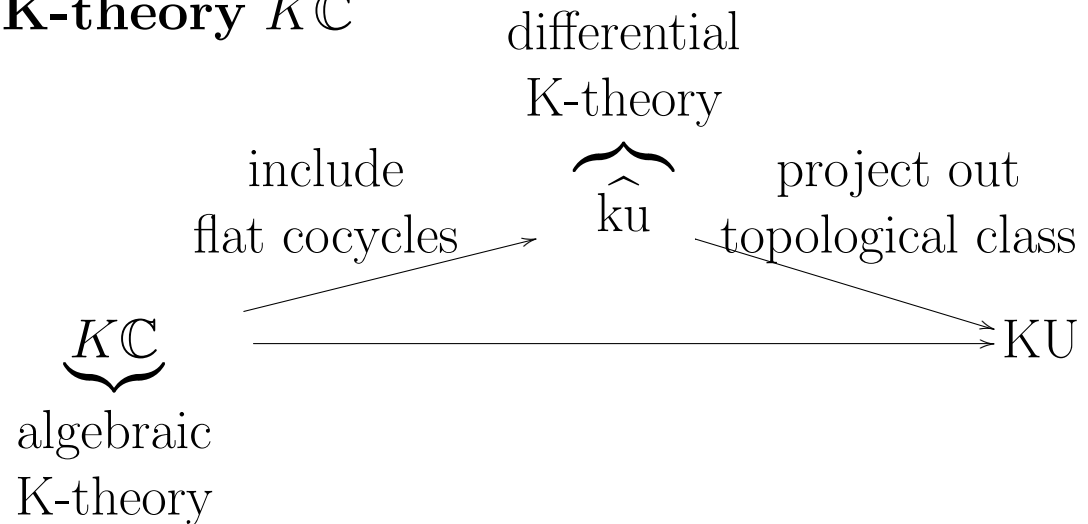
Algebraic K-Theory – Examples

- algebraic K-theory spectrum of a field \mathbb{F}

$$K\mathbb{F} \simeq \mathbb{K}(\mathbb{F}\text{Vect}_{\text{fin}})$$

- complex algebraic K-theory $K\mathbb{C}$

is *flat* K-theory:



- absolute algebraic K-theory $K\mathbb{F}_1 := \mathbb{K}(\mathbb{F}_1\text{Vect}_{\text{fin}}) \simeq \mathbb{K}(\text{Set}_{\text{fin}}^*)$
is *stable Cohomotopy theory* (Barrat-Priddy-Quillen theorem):

$$\boxed{
 \begin{array}{ccc}
 \underbrace{K\mathbb{F}_1} & \simeq & \underbrace{\mathbb{S}} \\
 \text{absolute algebraic} & & \text{sphere} \\
 \text{K-theory spectrum} & & \text{spectrum}
 \end{array}
 }$$

Brane charge on Orbifolds – Equivariant generalized cohomology

A *representation sphere* $\overset{G}{S^V}$:= one-point compactification of linear representation $\overset{G}{V}$

A *G-equivariant spectrum* \mathcal{A} is

a spectrum of G -spaces indexed by representation spheres, hence

1. a system of pointed G -spaces $\left\{ \overset{G}{\mathcal{A}_V} \mid \overset{G}{V} \text{ a linear } G\text{-representation} \right\}$

2. with equivariant suspension morphisms $S^V \wedge \mathcal{A}_W \xrightarrow{\sigma_{V,W}} \mathcal{A}_{V \oplus W}$

Examples

• The *equivariant suspension spectrum* of a G -space $\overset{G}{S^V}$ is

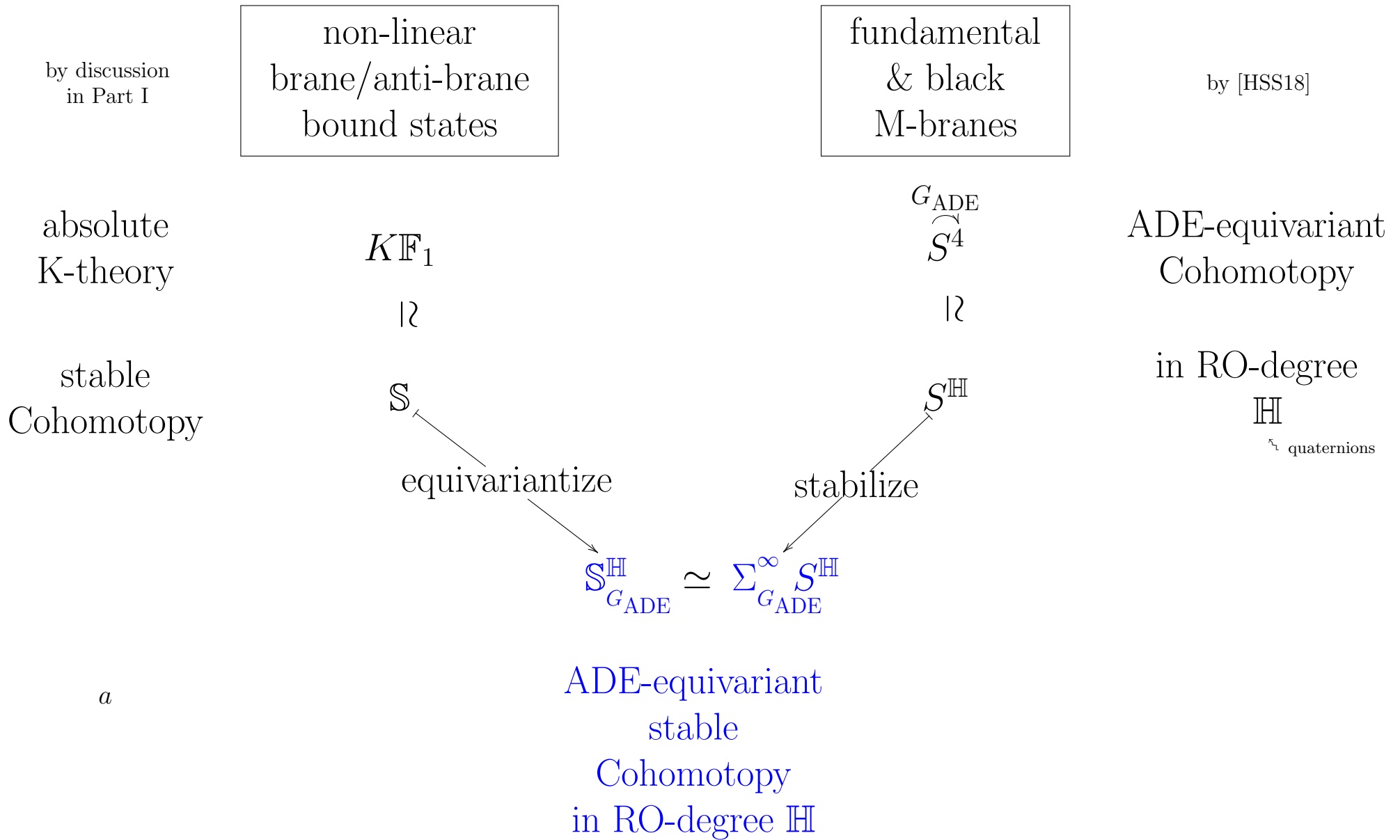
$$\Sigma_G^\infty X : V \mapsto S^V \wedge X.$$

• The *equivariant K-theory* of a contractible space is the *representation ring*

$$\mathrm{KU}_G(\overset{G}{\mathbb{R}^{d,1}}) \simeq \mathrm{KU}_G(\overset{G}{*}) \simeq R_{\mathbb{C}}(G) \simeq \mathbb{Z}[\overbrace{\rho_1, \dots, \rho_n}^{\text{irreps}}] \quad \text{“fractional D-branes”}$$

In conclusion, from Part I:

A compelling candidate for M-brane charge cohomology theory is...



Hypothesis H:

The
generalized cohomology theory
for
M-brane charge

is

ADE-equivariant
stable
Cohomotopy
in RO-degree \mathbb{H}

Hypothesis **H** predicts M-brane charge groups:

$$\mathbb{S}_{G_{\text{ADE}}}^{\mathbb{H}} \left(\underbrace{X}_{\text{11d spacetime orbifold}} \right)^{G_{\text{ADE}}}$$

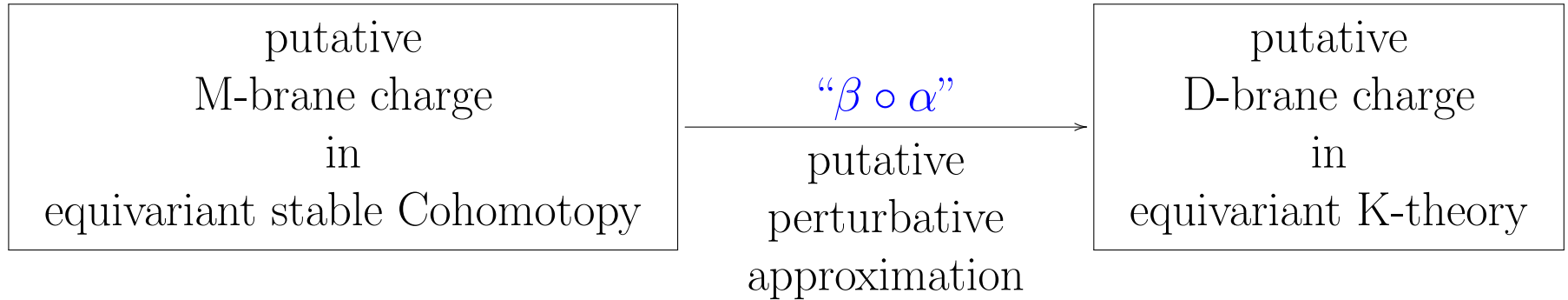
How does this compare to / clarify folklore of perturbative string theory:

- intersecting M2-branes \rightsquigarrow fractional D-branes ?
- M-theoretic “discrete torsion” of fractional M2-branes ?
- GUT at E-type singularities ?
- ...

This we discuss now \longrightarrow

Strategy for testing Hypothesis **H**

1. **Identify** suitable comparison homomorphism



2. **Compute:**

the co-kernel of $\beta \circ \alpha$; reflects	D-brane configurations that do not lift to M-theory
the kernel of $\beta \circ \alpha$; reflects	M-brane degrees of freedom invisible in perturbative string theory

Hypothesis H finds support if the **cokernel of $\beta \circ \alpha$** is

1. **small** \Leftrightarrow putative M-brane charge mostly reproduces string theory folklore,
2. **plausible** \Leftrightarrow the putative D-brane states in the co-kernel are dubious.

If so, Hypothesis **H** predicts the **kernel of $\beta \circ \alpha$** as hidden M-theoretic DOFs.

Outline

Since the sphere spectrum \mathbb{S}
is the *initial* commutative ring spectrum,
there is a unique multiplicative comparison morphism
from stable cohomotopy
to *every* other multiplicative cohomology theory \mathcal{A} ,
called the
equivariant generalized Boardman homomorphism

$$\mathbb{S}_G^\alpha(X) \xrightarrow{G} \mathcal{A}_G^\alpha(X)$$

Here we present two cases:

1. **Comparison map **A**** to
K-theory and RR-charge of fractional D-branes
2. **Comparison map **B**** to
ordinary cohomology and “discrete torsion” of fractional M5-branes

Comparison \mathbf{A} to
K-theory
and
fractional RR-charge of D-branes

[back to Part II](#)

Finite subgroups $G_{ADE} \subset SU(2)$ – Classification

Dynkin Label	Finite subgroup of $SU(2)$	Name of group
$A_{n \geq 1}$	Z_{n+1}	Cyclic
$D_{n \geq 4}$	$2D_{2(n-2)}$	Binary dihedral
E_6	$2T$	Binary tetrahedral
E_7	$2O$	Binary octahedral
E_8	$2I$	Binary icosahedral

Assumption: In the following, consider finite groups

$$G = G_{DE} \subset \mathbb{E} \subset SU(2)$$

in the D- or E-series

and

in the exceptional subgroup lattice.

next slide →

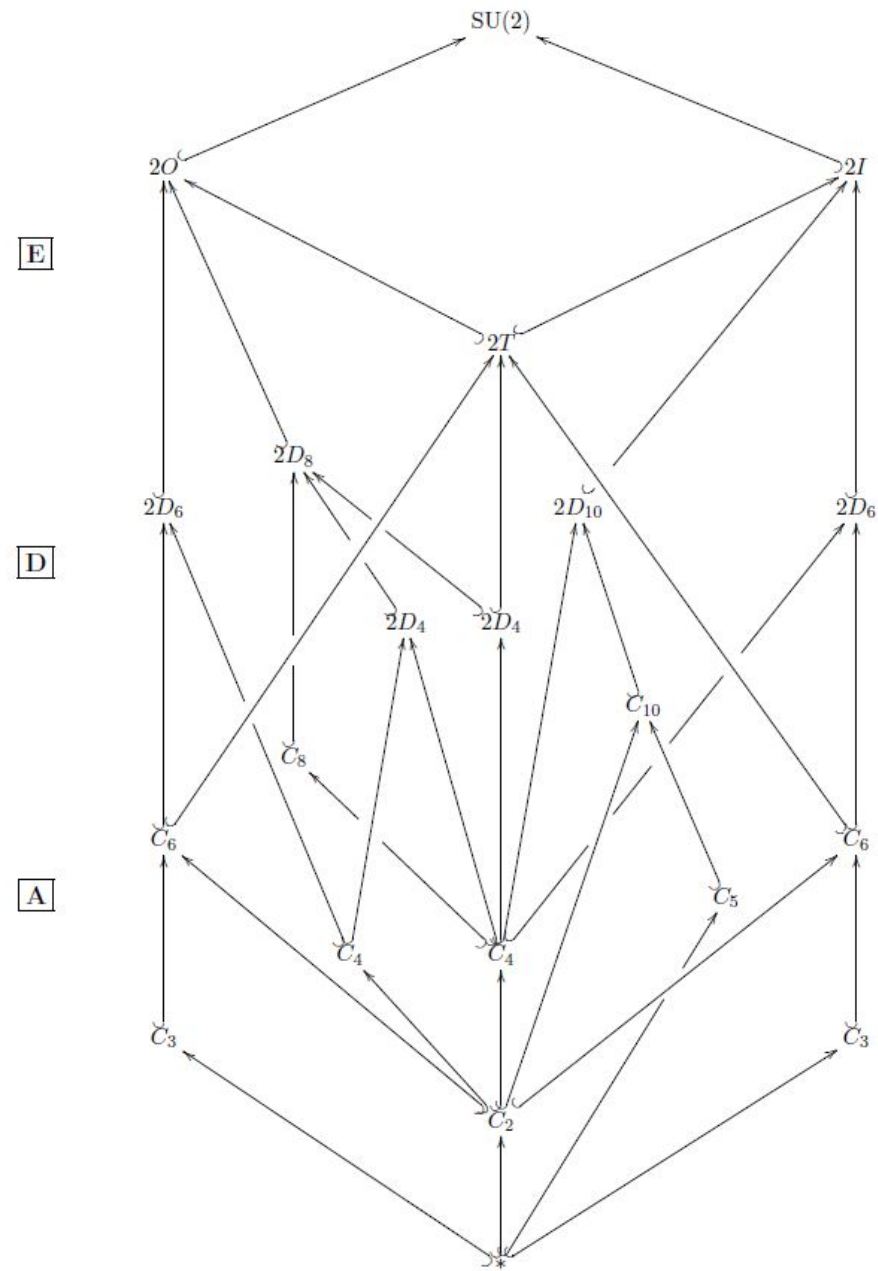
Implies in particular:

G -orbi-folds are G -orienti-folds,

the relevant K-theory for fractional D-brane charge

at G -fixed points is KO-theory

Finite subgroups $G_{ADE} \subset SU(2)$ – Exceptional subgroup lattice



The comparison homomorphism A

charge group
of M-branes
at two intersecting
MK6 ADE-singularities

visible in
M-theory
(?)

charge lattice
of fractional M-branes
at an MK6-singularity

visible in
M-theory
(?)

charge lattice
of fractional D-branes
at orientifold singularity

expected in
perturbative
string theory

equivariant stable cohomotopy in $RO(G \times G')$ -degree \mathbb{H}

equivariant KO-theory
in degree 0

$$\begin{array}{ccccc}
 \mathbb{S}_{(G \times G')}^{\mathbb{H}}(\mathbb{R}^{10,1}) & \xrightarrow{\text{zoom in onto one MK6}} & \mathbb{S}_G^{(\mathbb{H}^{G'})}(\mathbb{R}^{6,1}) & \xrightarrow{\text{perturbative approximation}} & \text{KO}_G^0(\mathbb{R}^{6,1}) \\
 \wr & & \wr & & \wr \\
 \mathbb{S}_{(G \times G')}^{\mathbb{H}}(*) & \xrightarrow[\text{form geometric } G'\text{-fixed spectrum}]{\alpha} & \mathbb{S}_G^0(*) & \xrightarrow{\text{Boardman homomorphism}} & \text{KO}_G^0(*) \\
 & & \wr & & \wr \\
 & & A(G) & \xrightarrow[\text{linearize } G\text{-actions}]{\beta} & \text{RO}(G)
 \end{array}$$

module over
 $(G \times G')$ -Burnside ring

Thm. 1

G -Burnside ring

Thm. 2

G -Representation ring

Theorem 1

charge group
of M-branes
at two intersecting
MK6 ADE-singularities

visible in
M-theory
(?)

charge lattice
of fractional M-branes
at an MK6-singularity

visible in
M-theory
(?)

Proof.
Use Prop. II 9.13 in
[LewisMaySteinberger86].
□

equivariant stable cohomotopy in $RO(G \times G')$ -degree \mathbb{H}

$$\begin{array}{ccc}
 \mathbb{S}_{(G \times G')}^{\mathbb{H}}(\mathbb{R}^{10,1}) & \xrightarrow{\text{zoom in onto one MK6}} & \mathbb{S}_G^{(\mathbb{H}^{G'})}(\mathbb{R}^{6,1}) \\
 \wr & & \wr \\
 \mathbb{S}_{(G \times G')}^{\mathbb{H}}(*) & \xrightarrow[\text{form geometric } G'\text{-fixed spectrum}]{\alpha} & \mathbb{S}_G^0(*) \\
 & & \wr \\
 & & A(G)
 \end{array}$$

module over
 $(G \times G')$ -Burnside ring

Is surjective.

G -Burnside ring

hence: $\text{coker}(\beta \circ \alpha) \simeq \text{coker}(\beta)$

The comparison homomorphism A

charge group
of M-branes
at two intersecting
MK6 ADE-singularities

visible in
M-theory
(?)

charge lattice
of fractional M-branes
at an MK6-singularity

visible in
M-theory
(?)

charge lattice
of fractional D-branes
at orientifold singularity

expected in
perturbative
string theory

equivariant stable cohomotopy in $RO(G \times G')$ -degree \mathbb{H}

equivariant KO-theory
in degree 0

$$\begin{array}{ccccc}
 \mathbb{S}_{(G \times G')}^{\mathbb{H}}(\mathbb{R}^{10,1}) & \xrightarrow{\text{zoom in onto one MK6}} & \mathbb{S}_G^{(\mathbb{H}^{G'})}(\mathbb{R}^{6,1}) & \xrightarrow{\text{perturbative approximation}} & \text{KO}_G^0(\mathbb{R}^{6,1}) \\
 \wr & & \wr & & \wr \\
 \mathbb{S}_{(G \times G')}^{\mathbb{H}}(*) & \xrightarrow[\text{form geometric } G'\text{-fixed spectrum}]{\alpha} & \mathbb{S}_G^0(*) & \xrightarrow{\text{Bordman homomorphism}} & \text{KO}_G^0(*) \\
 & & \wr & & \wr \\
 & & A(G) & \xrightarrow[\text{linearize } G\text{-actions}]{\beta} & RO(G)
 \end{array}$$

module over
 $(G \times G')$ -Burnside ring

Thm. 1

G -Burnside ring

Thm. 2

G -Representation ring

Theorem 2

$$\underbrace{\text{irrational characters}}_{\text{RO}^{\text{irrational}}(G)} \cong \text{RO}(G) / \underbrace{\text{integral characters}}_{\text{RO}^{\text{int}}(G)}$$

charge lattice of fractional M-branes at an MK6-singularity

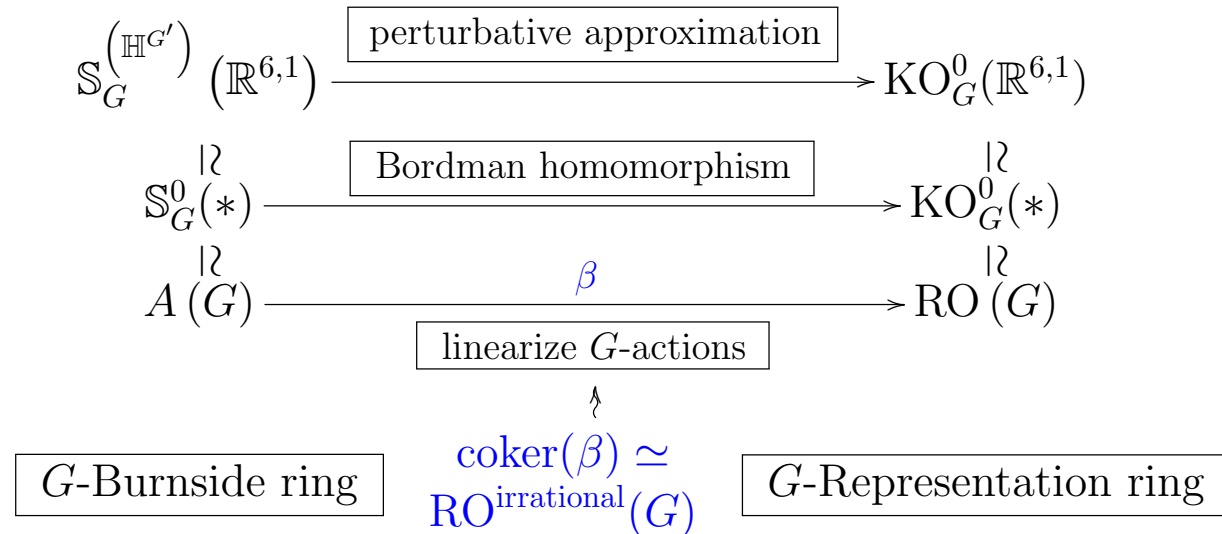
visible in M-theory (?)

charge lattice of fractional D-branes at orientifold singularity

expected in perturbative string theory

equivariant stable cohomotopy in degree 0

equivariant KO-theory in degree 0



Theorem 2 – Ingredients

G -actions
on finite sets
(not-linear)

G -representations
on \mathbb{F} -vector spaces
(linear)

Burnside ring of
virtual G -sets

Representation ring
of virtual representations

G -set

$$\begin{array}{ccc}
 (G\text{Set}_{\text{fin}}, \sqcup) & \xrightarrow{\text{linearize}_{\mathbb{F}[-]}} & (G\text{Rep}_{\mathbb{F}}, \oplus) \\
 \downarrow K & & \downarrow K \\
 A(G) & \xrightarrow{\beta} & R_{\mathbb{F}}(G) \\
 \\
 \begin{array}{c} \overline{G} \\ \underbrace{S} \end{array} & \xrightarrow{\quad} & \begin{array}{c} \overline{G} \\ \underbrace{\mathbb{F}[S]} \end{array}
 \end{array}$$

“permutation
representation”

the **cokernel** of β reflects linear algebra
invisible to pure combinatorics

the **kernel** of β reflects pure combinatorics
invisible to linear algebras

Theorem 2 – Proof

Compute:

1. set of [conjugacy classes](#) $\{[H_i]\}$ of subgroups $H \subset G$
2. the [Burnside product](#) $[G/H_i] \times [G/H_j] = \bigsqcup_{\ell} \underbrace{n_{ij}^{\ell}}_{\substack{\text{structure} \\ \text{constants}}} \cdot [G/H_{\ell}]$
3. its matrix of [total multiplicities](#) $\text{mult}_{ij} := \sum_{\ell} n_{ij}^{\ell}$
4. its integral [row reduction](#) $\underbrace{H}_{\substack{\text{upper} \\ \text{triangular}}} := \underbrace{U}_{\in \text{GL}(N, \mathbb{Z})} \cdot \text{mult}$

Lemma. *The rows of H span $\text{im}(\beta) \subset R_{\mathbb{F}}(G)$.*

This yields an [effective algorithm computing \$\text{coker}\(\beta\) = R_{\mathbb{F}}\(G\)/\text{im}\(\beta\)\$](#)

Simon Burton has implemented this algorithm in **Python**.

\Rightarrow **Proof of Theorem 2:** By brute force automatized computation. \square

Theorem 2 – Proof

	coker	$A(G) \xrightarrow{\beta_{\mathbb{F}}} R_{\mathbb{F}}(G)$			$A(G) \xrightarrow{\beta_{\mathbb{F}}^{\text{int}}} R_{\mathbb{F}}^{\text{int}}(G)$		
		ground field \mathbb{F}					
		\mathbb{Q}	\mathbb{R}	\mathbb{C}	\mathbb{Q}	\mathbb{R}	\mathbb{C}
finite group G	$2D_4$	0	0	$\frac{\mathbb{Z}[\rho_5]}{\mathbb{Z}[2\rho_5]}$	0	0	$\frac{\mathbb{Z}[\rho_5]}{\mathbb{Z}[2\rho_5]}$
	$2D_6$	0	0	$\frac{\mathbb{Z}[\rho_3, \rho_4, \rho_6]}{\mathbb{Z}[\rho_3 + \rho_4, 2\rho_6]}$	0	0	$\frac{\mathbb{Z}[\rho_6]}{\mathbb{Z}[2\rho_6]}$
	$2D_8$	0	$\frac{\mathbb{Z}[2\rho_6, 2\rho_7]}{\mathbb{Z}[2\rho_6 + 2\rho_7]}$	$\frac{\mathbb{Z}[\rho_6, \rho_7]}{\mathbb{Z}[2\rho_6 + 2\rho_7]}$	0	0	$\frac{\mathbb{Z}[\rho_6 + \rho_7]}{\mathbb{Z}[2\rho_6 + 2\rho_7]}$
	$2D_{10}$	0	$\frac{\mathbb{Z}[\rho_3, \rho_4, \rho_5, \rho_6, \rho_7, \rho_8]}{\mathbb{Z}[\rho_3 + \rho_4, \rho_5 + \rho_6, 2\rho_7 + 2\rho_8]}$	$\frac{\mathbb{Z}[\rho_3, \rho_4, \rho_5, \rho_6, 2\rho_7, 2\rho_8]}{\mathbb{Z}[\rho_3 + \rho_4, \rho_5 + \rho_6, 2\rho_7 + 2\rho_8]}$	0	0	$\frac{\mathbb{Z}[\rho_7 + \rho_8]}{\mathbb{Z}[2\rho_7 + 2\rho_8]}$
	$2D_{12}$	0	$\frac{\mathbb{Z}[\rho_7, \rho_8, \rho_9]}{\mathbb{Z}[2\rho_7, 2\rho_8 + 2\rho_9]}$	$\frac{\mathbb{Z}[2\rho_8, 2\rho_9]}{\mathbb{Z}[2\rho_8 + 2\rho_9]}$	0	0	$\frac{\mathbb{Z}[\rho_7]}{\mathbb{Z}[2\rho_7]}$
	$2T$	0	0	$\frac{\mathbb{Z}[\rho_2, \rho_2^*, \rho_4, \rho_4^*, \rho_5]}{\mathbb{Z}[\rho_2 + \rho_2^*, \rho_4 + \rho_4^*, 2\rho_5]}$	0	0	$\frac{\mathbb{Z}[\rho_5]}{\mathbb{Z}[2\rho_5]}$
	$2O$	0	$\frac{\mathbb{Z}[2\rho_6, 2\rho_7]}{\mathbb{Z}[2\rho_6 + 2\rho_7]}$	$\frac{\mathbb{Z}[\rho_6, \rho_7, \rho_8]}{\mathbb{Z}[2\rho_6 + 2\rho_7, 2\rho_8]}$	0	0	$\frac{\mathbb{Z}[\rho_8]}{\mathbb{Z}[2\rho_8]}$
	$2I$	0	$\frac{\mathbb{Z}[2\rho_2, 2\rho_3, \rho_4, \rho_5]}{\mathbb{Z}[2\rho_2 + 2\rho_3, \rho_4 + \rho_5]}$	$\frac{\mathbb{Z}[\rho_2, \rho_3, \rho_4, \rho_5, \rho_7, \rho_9]}{\mathbb{Z}[2\rho_2 + 2\rho_3, \rho_4 + \rho_5, 2\rho_7, 2\rho_9]}$	0	0	$\frac{\mathbb{Z}[\rho_2 + \rho_3, \rho_4 + \rho_5, \rho_7, \rho_9]}{\mathbb{Z}[2\rho_2 + 2\rho_3, \rho_4 + \rho_5, 2\rho_7, 2\rho_9]}$

Theorem 2 – Physics interpretation

Let $V \in K_G(*) \simeq R(G)$ a fractional D-brane \leftrightarrow G-representation,

RR-charge in the g -twisted closed string sector

is the [value of its character](#) at g :

$$Q_V^{\text{RR}}(g) = \frac{1}{|G|} \chi_V(g)$$

([DouglasGreeneMorrisson97, (3.8)], [DiGo00, (2.4)], [BCR00, (4.65) with (4.41)], [EGJ05, (4.5)], [ReSc13, 4.102])

Theorem 2 – Physics reformulation:

*Hypothesis **H** implies*

that fractional D-branes with irrational RR-charge are spurious.

Physically plausible?

Some $V \in K_G()$ must be spurious [BDHKMMS02, 4.5.2].*

Irrational RR-charge called a *paradox* in [BachasDouglasSchweigert00, (2.8)], also [Taylor00, Zho01, Rajan02], apparently unresolved.

If this is indeed a *paradox*,

then hypothesis **H** exactly resolves it.

Theorem 2 – Physics interpretation

Regard $\text{coker}(\beta)$ under **McKay correspondence** :

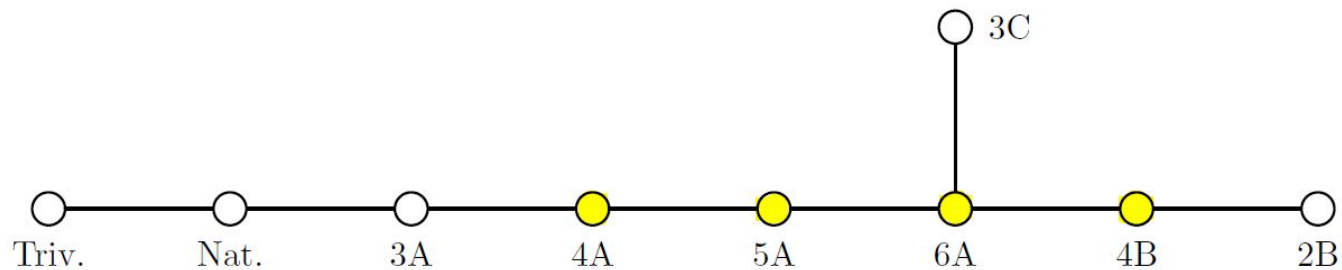
$$\{ \text{irreps } \rho \in R_{\mathbb{C}}(G_{\text{ADE}}) \} \simeq \left\{ \begin{array}{l} \text{vertices of corresponding} \\ \text{ADE-type Dynkin diagram} \end{array} \right\}$$

Most exceptional Example: $G = 2I$:

	e	a	a^2	a^3	a^2b	a^4	a^3b	a^5	a^4b
Triv.	1	1	1	1	1	1	1	1	1
Nat.	2	$\frac{1}{2}(1 + \sqrt{5})$	$-\frac{1}{2}(1 - \sqrt{5})$	$\frac{1}{2}(1 - \sqrt{5})$	1	$-\frac{1}{2}(1 + \sqrt{5})$	0	-2	-1
3A	3	$\frac{1}{2}(1 + \sqrt{5})$	$\frac{1}{2}(1 - \sqrt{5})$	$\frac{1}{2}(1 - \sqrt{5})$	0	$\frac{1}{2}(1 + \sqrt{5})$	-1	3	0
4A	4	1	-1	1	-1	-1	0	-4	1
5A	5	0	0	0	-1	0	1	5	-1
6A	6	-1	1	-1	0	1	0	-6	0
2B	2	$\frac{1}{2}(1 - \sqrt{5})$	$-\frac{1}{2}(1 + \sqrt{5})$	$\frac{1}{2}(1 + \sqrt{5})$	1	$-\frac{1}{2}(1 - \sqrt{5})$	0	-2	-1
4B	4	-1	-1	-1	1	-1	0	4	1
3C	3	$\frac{1}{2}(1 - \sqrt{5})$	$\frac{1}{2}(1 + \sqrt{5})$	$\frac{1}{2}(1 + \sqrt{5})$	0	$\frac{1}{2}(1 - \sqrt{5})$	-1	3	0

integral/
non-irrational
characters

Dynkin
diagram



hence:

$$\text{im}(\beta)|_{\text{irred}} \subset \text{RO}(2I)|_{\text{irred}} \Leftrightarrow \underbrace{SU(5)}_{\text{actual GUT group}} \subset \underbrace{E_8}_{\text{stringy GUT group}}$$

Comparison B to
ordinary Cohomology
and
“discrete torsion” of fractional M5-branes

[back to Part II](#)

The comparison homomorphism B

Away from the singular locus

of a black M2-brane

$$\begin{array}{ccc}
 \text{AdS}_4 \times \underbrace{S^7 / G_A}_{\text{spherical space form}} & \xleftarrow{\ell_P \gg 1} \cdots \xrightarrow{\ell_P \ll 1} & \mathbb{R}^{2,1} \times \underbrace{\mathbb{R}^8 // G_A}_{\text{orbifold du Val singularity}}
 \end{array}$$

the orbifold is smooth and, for A-type singularities, so is the RO-degree:

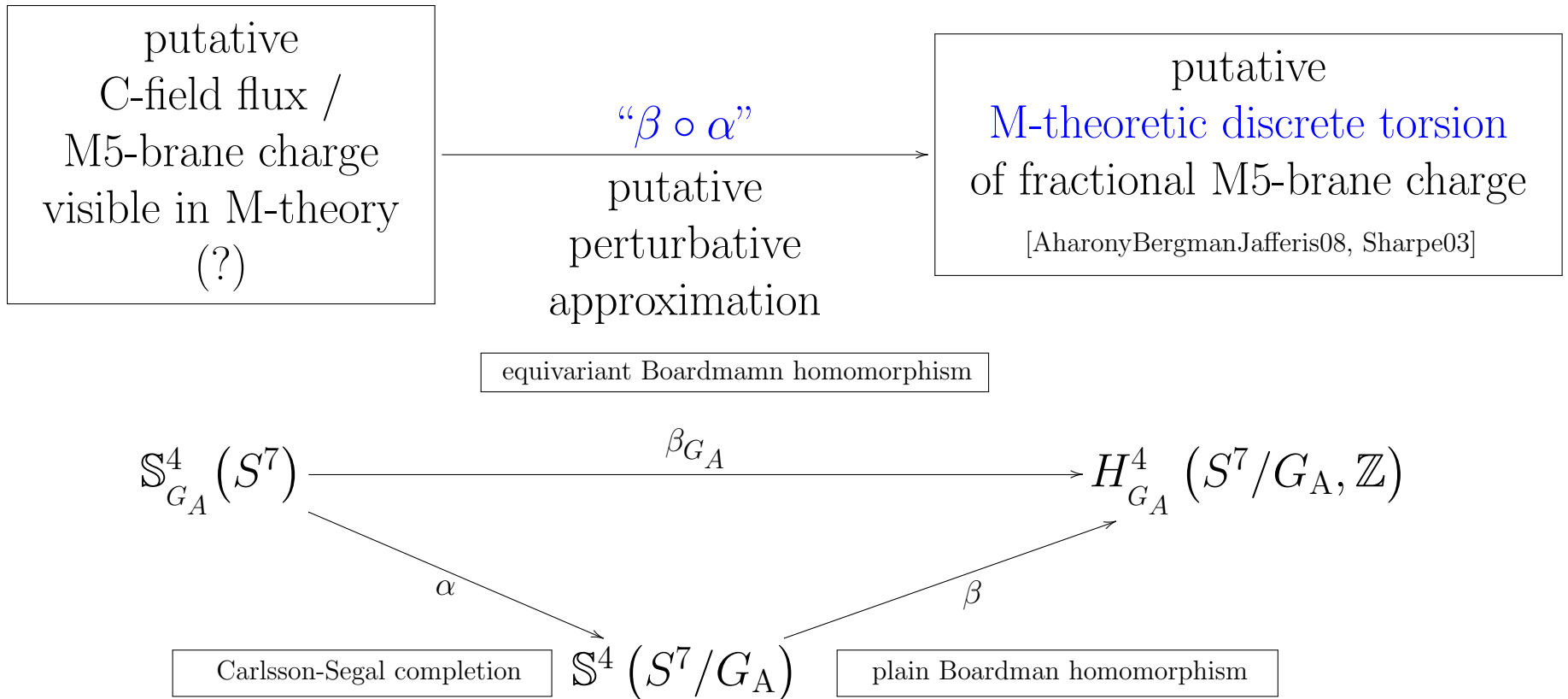
$$\begin{array}{ccc}
 \text{coefficient bundle} & ((\mathbb{R}^{10,1} - \mathbb{R}^{2,1}) \times \mathbb{H}) // G & (\mathbb{R}^{10,1} - \mathbb{R}^{2,1}) / G \times \mathbb{R}^4 \\
 & \downarrow & \downarrow \\
 \text{spacetime orbifold} & (\mathbb{R}^{10,1} - \mathbb{R}^{2,1}) // G_A & (\mathbb{R}^{10,1} - \mathbb{R}^{2,1}) / G_A
 \end{array} \quad \simeq$$

⇒ relevant comparison morphism is [equivariant Boardman homomorphism](#)

$$\underbrace{S_{G_A}^4(S^7)}_{\text{equivariant stable Cohomotopy}} \xrightarrow{\beta_{G_A}} \underbrace{HZ_{G_A}^4(S^7)}_{\text{equivariant ordinary cohomology}} \simeq \underbrace{H^4(S^7 / G_A, \mathbb{Z})}_{\text{Borel equivariance}}$$

The comparison homomorphism B

Theorem 3 i): factors through plain Boardman homomorphism:



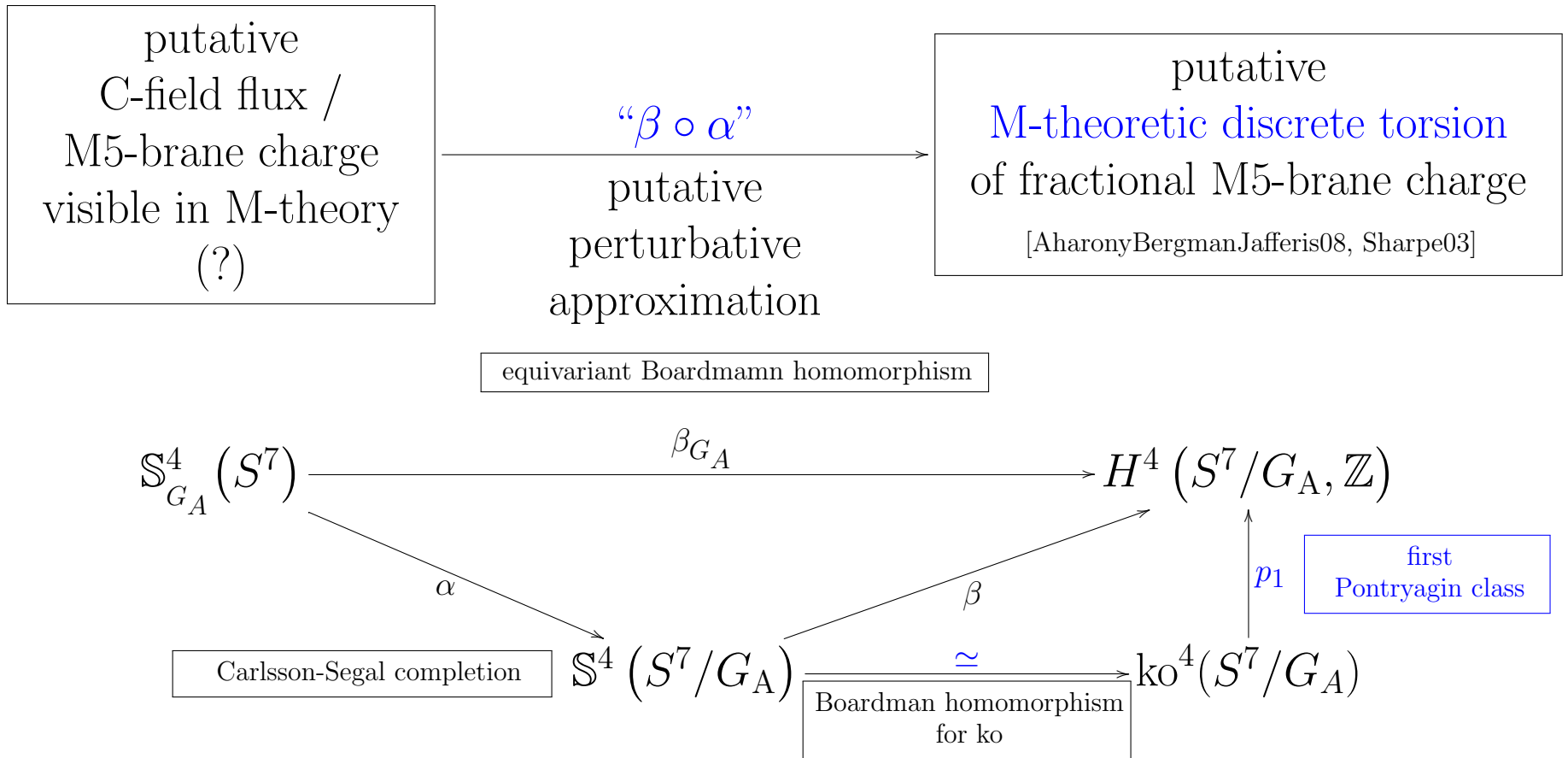
Proof. Use [Schwede18, Example 4.5.19]. \square .

Theorem 3 ii): $4 \text{ coker}(\beta) = 0$

Proof. By [Arlettaz04, Theorem 1.2 b)]. \square

The comparison homomorphism B

Theorem 3 iii): factors isomorphically through ko^4 :



Proof. By the AHSS and using $\pi_{\bullet \leq 2}(\mathbb{S}) = \pi_{\bullet \leq 2}(ko)$ \square .

Physically reasonable?

This $\text{coker}(\beta)$ is KO-version of what was argued for KU in [DiaconescuMooreWitten00].

Conclusion

1. Part I – Motivation of differential equivariant Cohomotopy:

- (a) [Derivation](#) of equivariant cohomotopy/ \mathbb{Q} [from first principles](#) via rational super homotopy theory [FSS13, FSS16a, FSS16b, BSS18, HSS18]
- (b) actual Cohomotopy is the minimal non-rational lift – differential equivariant Cohomotopy of super-orbifolds exists in global equivariant super homotopy theory
- (c) Hypothesis **H**:
the observables of M-theory are the differential equivariant real cohomology of the moduli stack of supertorsion-free differential equivariant Cohomotopy of spacetime $\text{Pin}(5)^b$ -orbi $\mathbb{R}^{10,1|32}$ -folds

2. Part II – Consistency checks of Hypothesis **H**:

- (a) reproduces fractional D-brane charge in equivariant K-theory
 - i. excluding exactly the spurious irrational RR-charges,
 - ii. which may correspond, via McKay, to breaking E_8 to $SU(5)$ GUT
- (b) reproduces discrete torsion of fractional M5-branes with DMW-correction.

In particular, equivariant stable cohomotopy somehow [unifies ordinary cohomology of the C-field with K-theory of D-branes.](#)

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