

Twisted, Equivariant, Differential (TED) K-Theory via Diffeological Stacks

Urs Schreiber on joint work with Hisham Sati



NYU AD Science Division, Program of Mathematics
& Center for Quantum and Topological Systems
New York University, Abu Dhabi



talk via:

Global Diffeology Seminar

@

DIFFEOLOGY.NET

04 Nov 2022

Motivation:

Motivation: understand the real world

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start simple: understand quantum physics

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Tool: [Mathematical foundations.](#)

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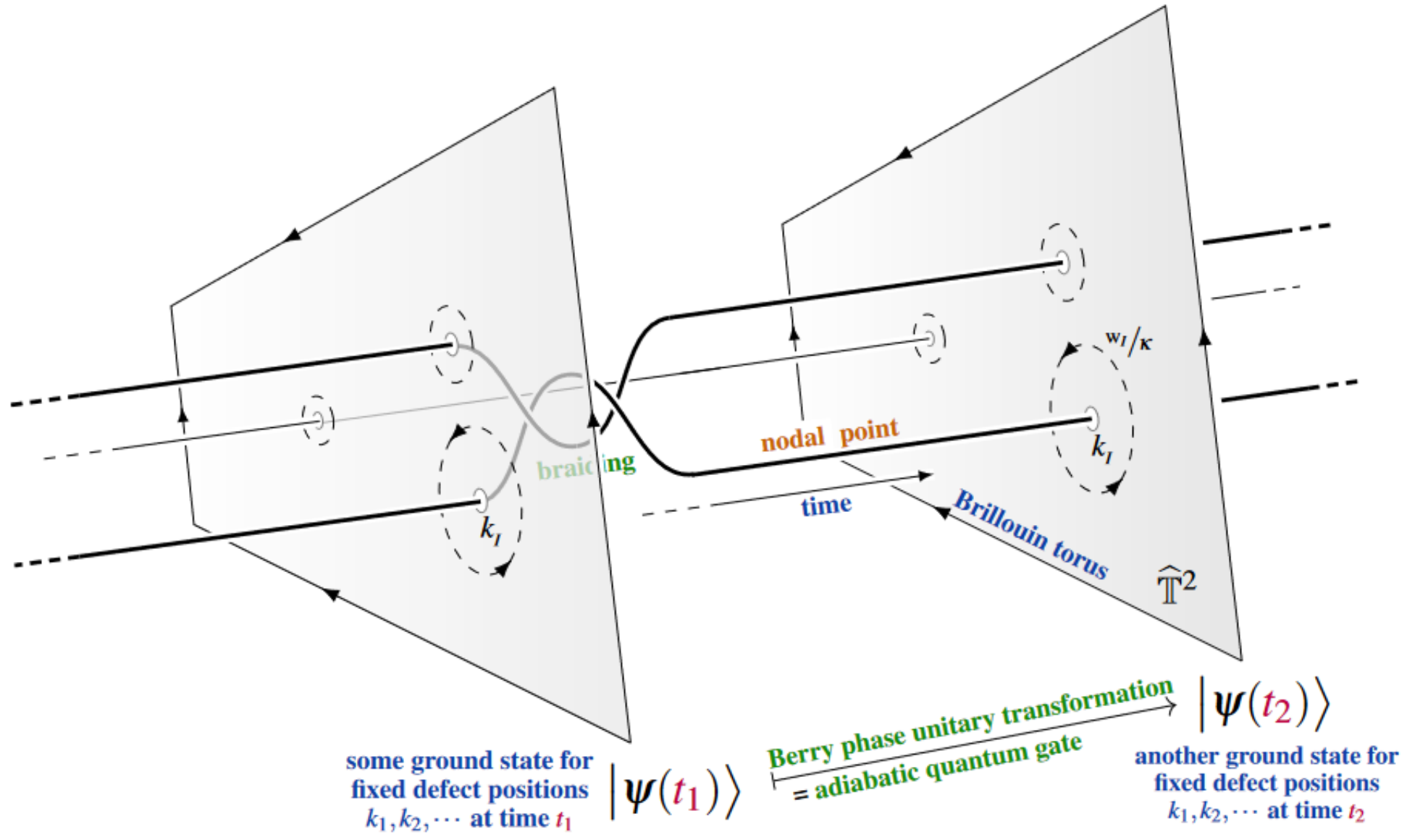


Figure 1 – Adiabatic braid quantum gate. Schematically indicated is the unitary transformation induced on the topologically ordered ground state (as discussed below in §3.3) of an effectively 2-dimensional topological semi-metal (as in §3.1) under adiabatic braiding (Rem. 1.1) of nodal points in the Brillouin torus (Rem. 3.9).

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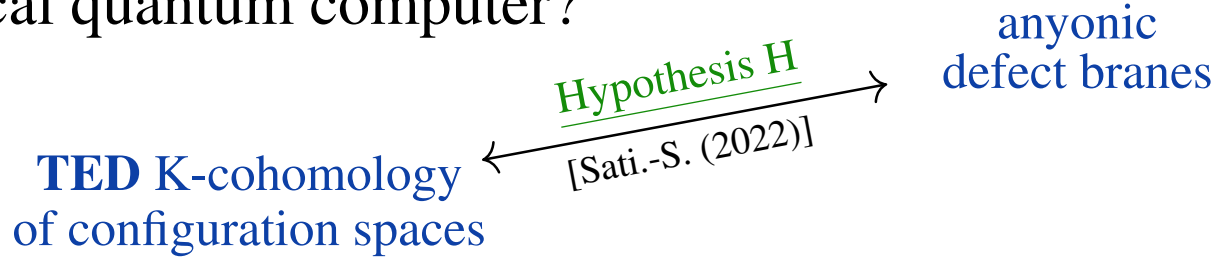
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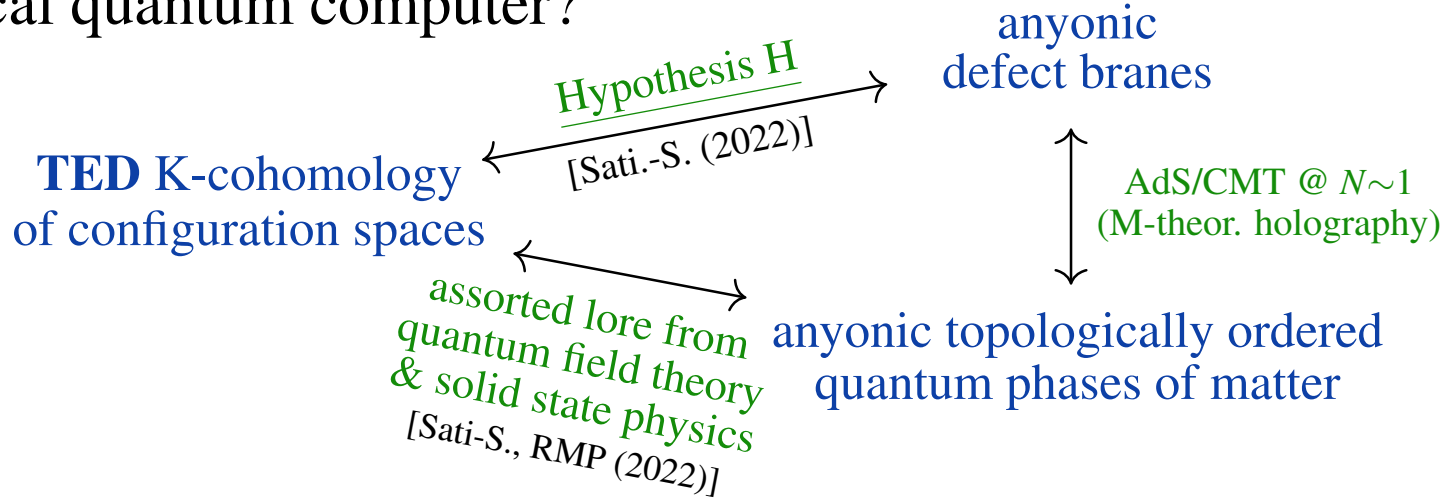
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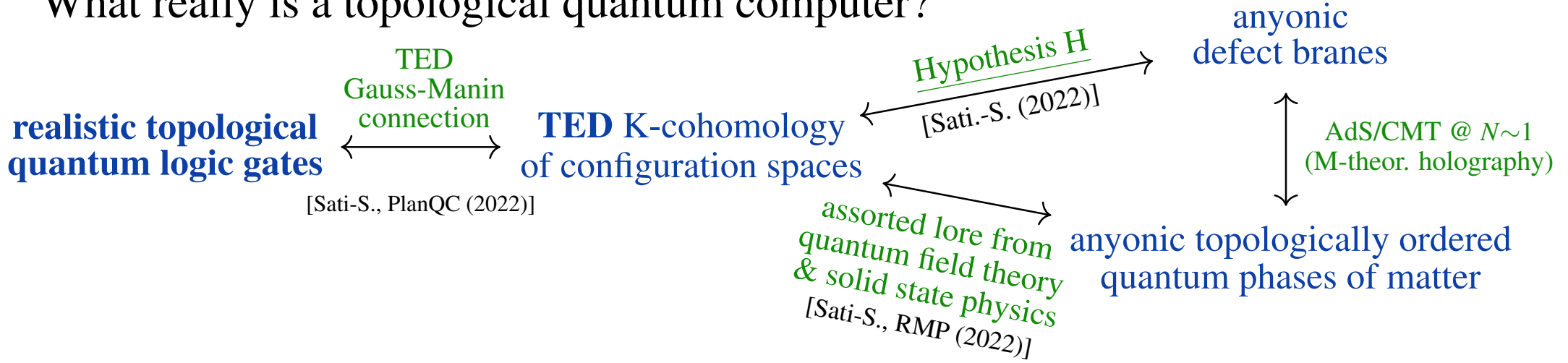
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I – Equivariant ∞ -Bundles

II – TED-K-Theory

III – Anyonic Defect Branes

IV – Quantum Computation

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This part is
a lightning indication of
how diffeological spaces fit
into the scheme of things,

following:

<i>Equivariant Principal ∞-bundles</i>	[arXiv:2112.13654]
<i>Proper Orbifold Cohomology</i>	[arXiv:2008.01101]
<i>Diff. Cohomology in a Cohesive ∞-Topos</i>	[arXiv:1310.7930]

Diffeological Spaces in Perspective.

Souriau (1979), Iglesias-Zemmour (1985) – paraphrased:

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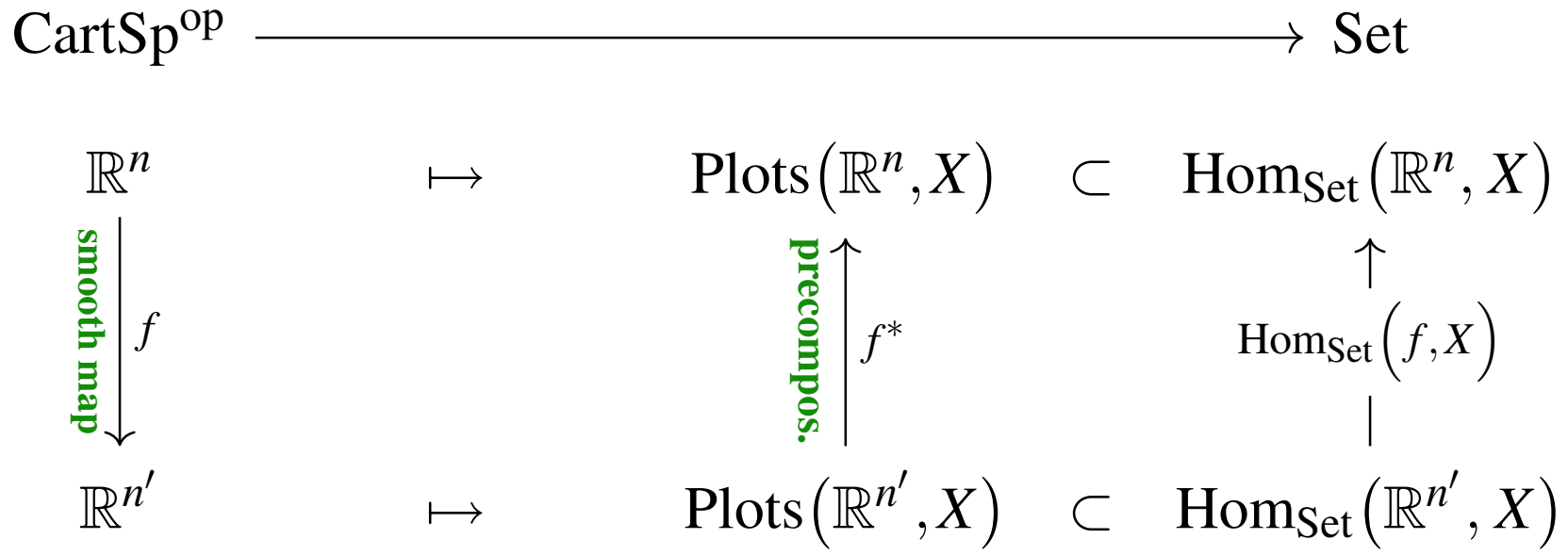
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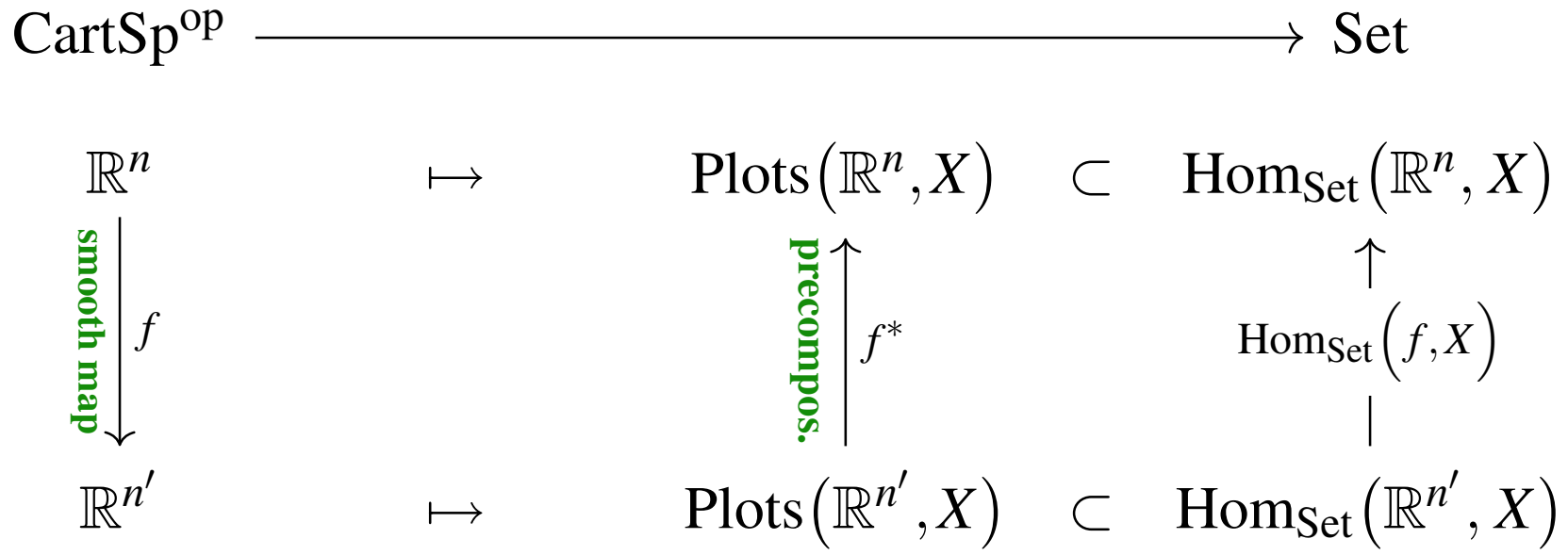


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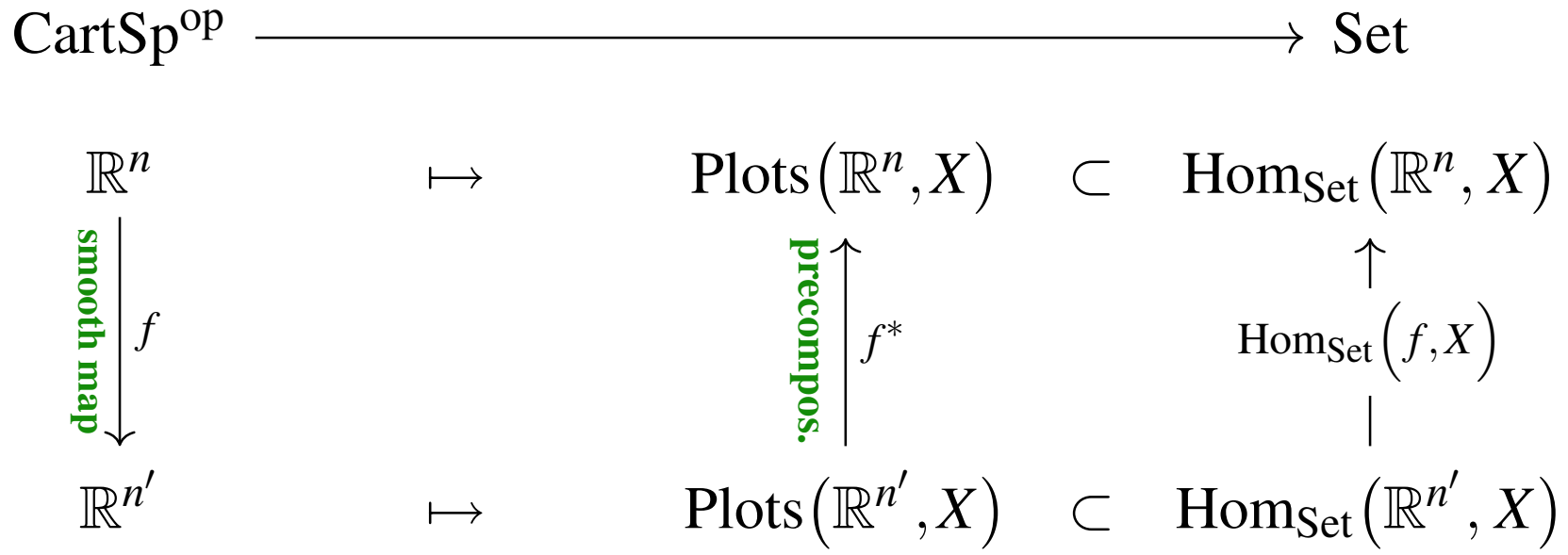
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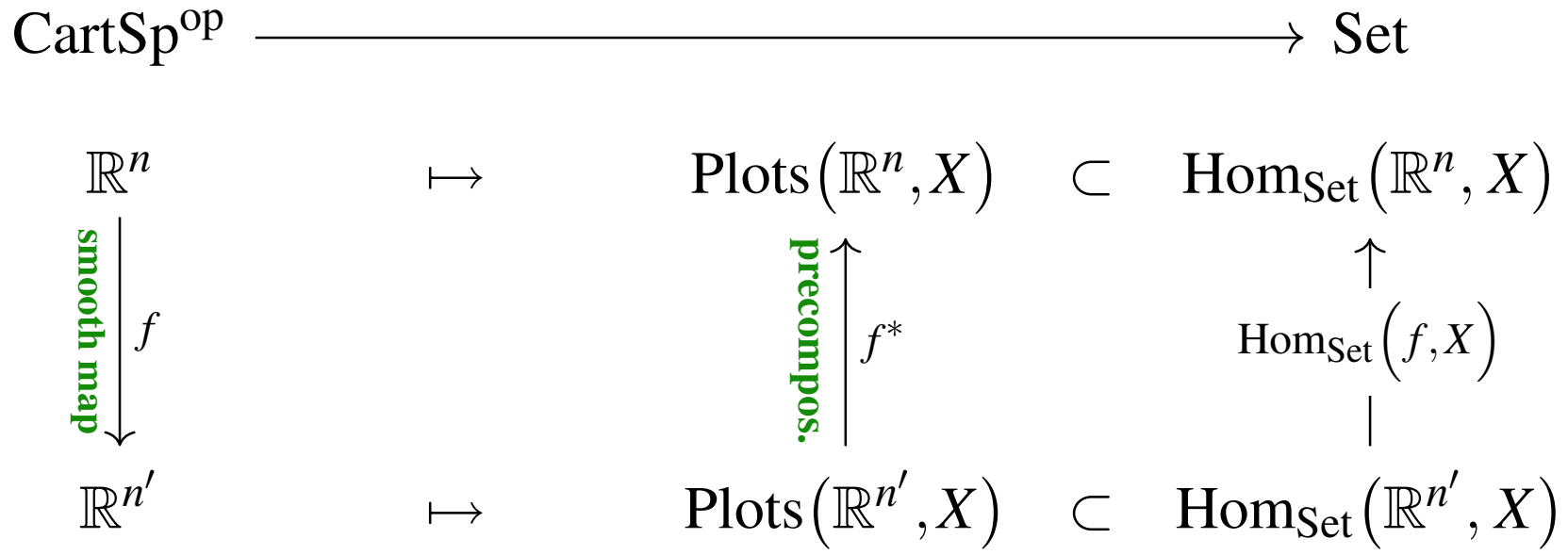
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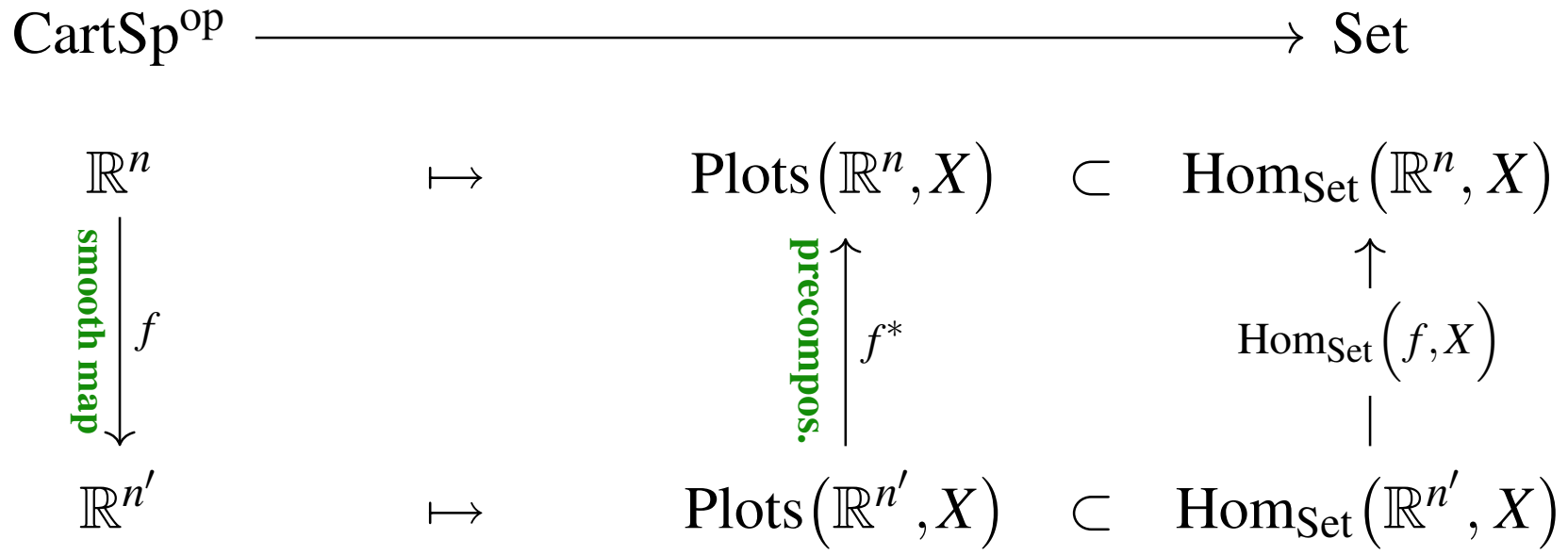
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TopSpc



form path ∞ -groupoids
(sing. simpl. compl.) Pth

Grpd $_{\infty}$

base ∞ -topos of
bare ∞ -groupoids

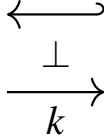
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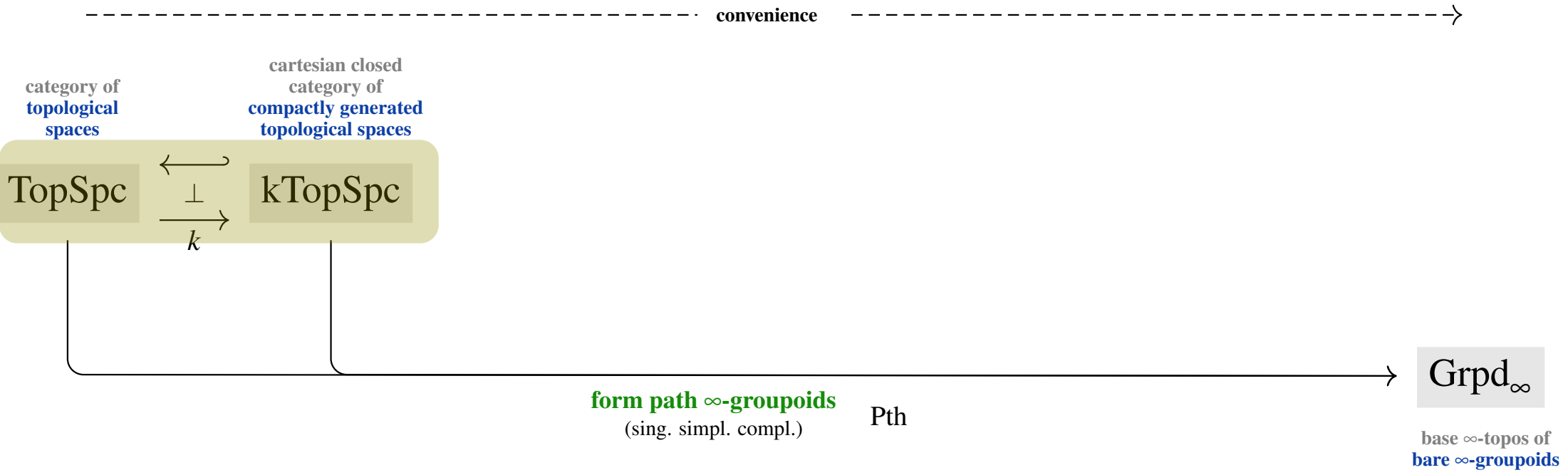


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Hurewicz (1948)
Gale (1950)
Steenrod (1967)

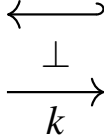
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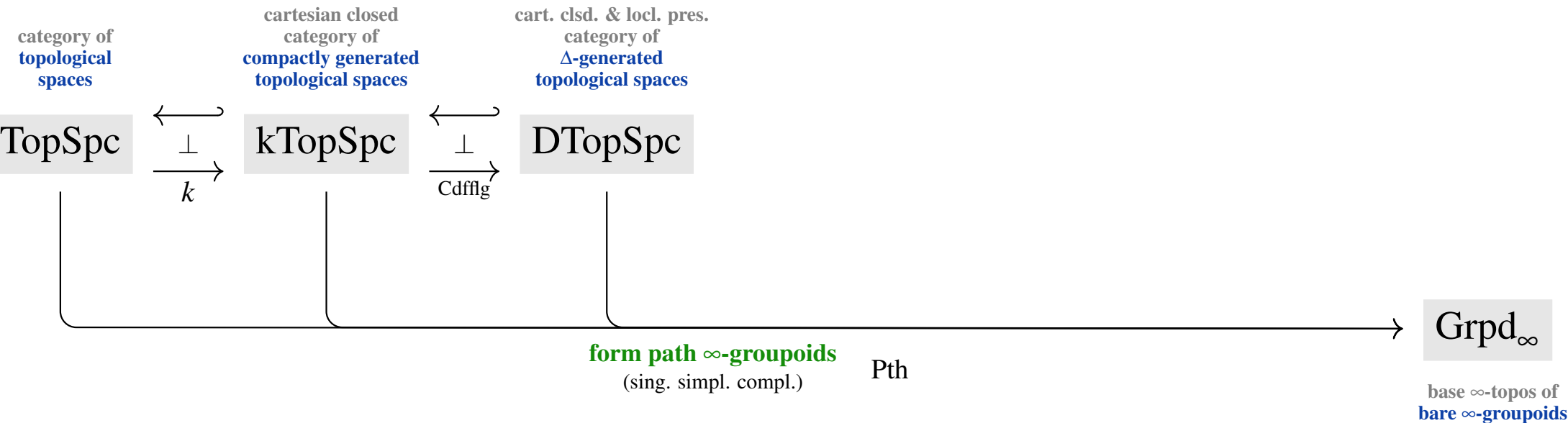
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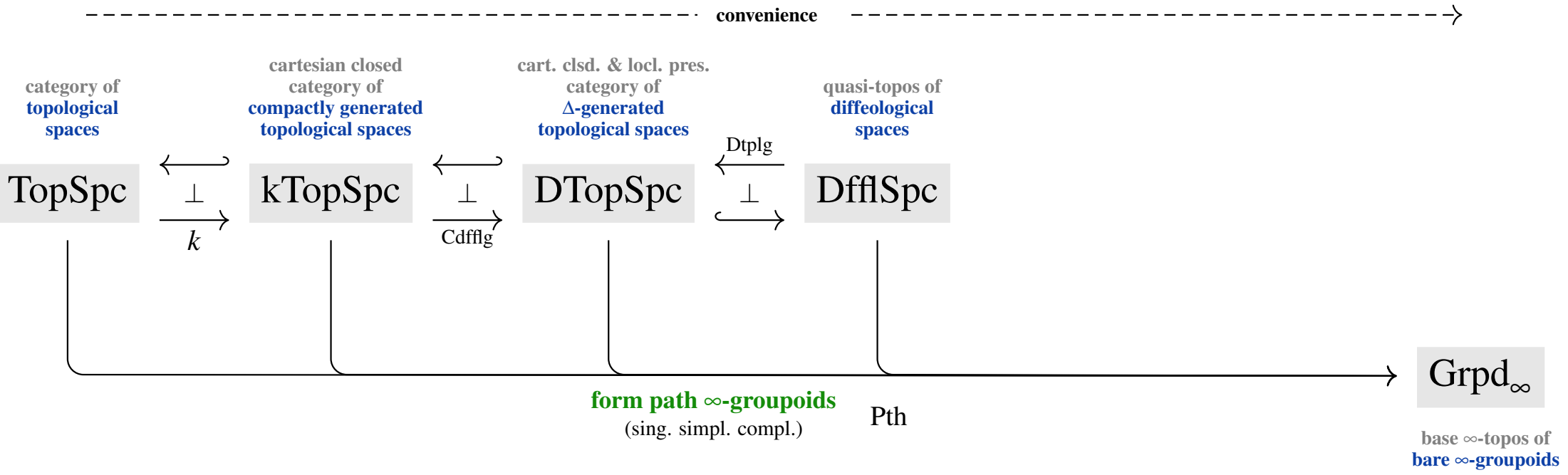
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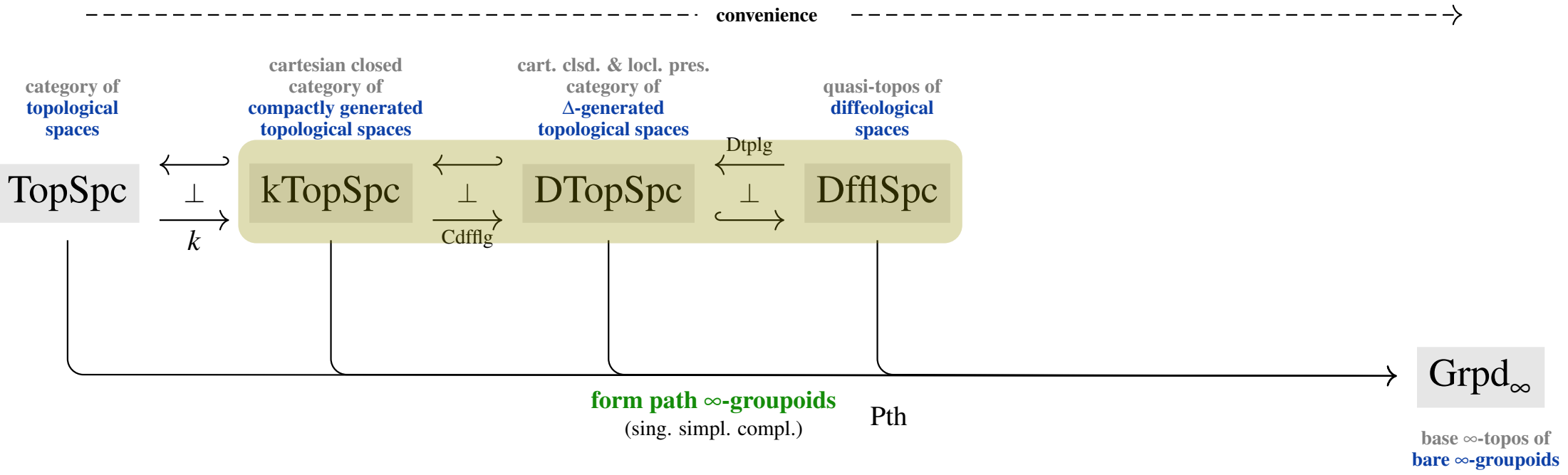
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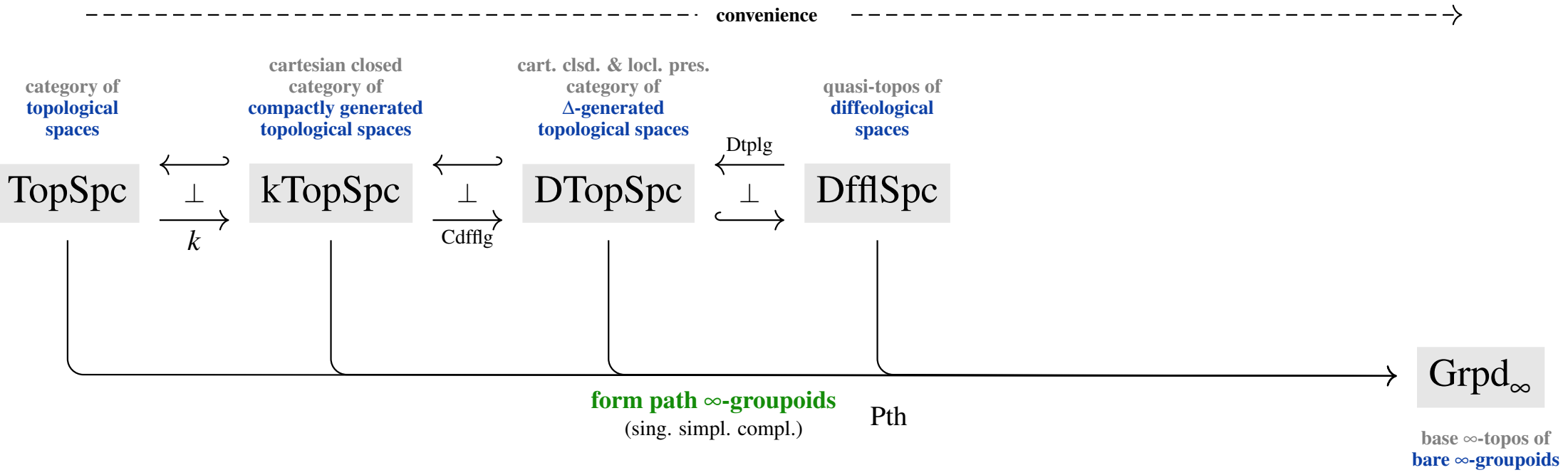
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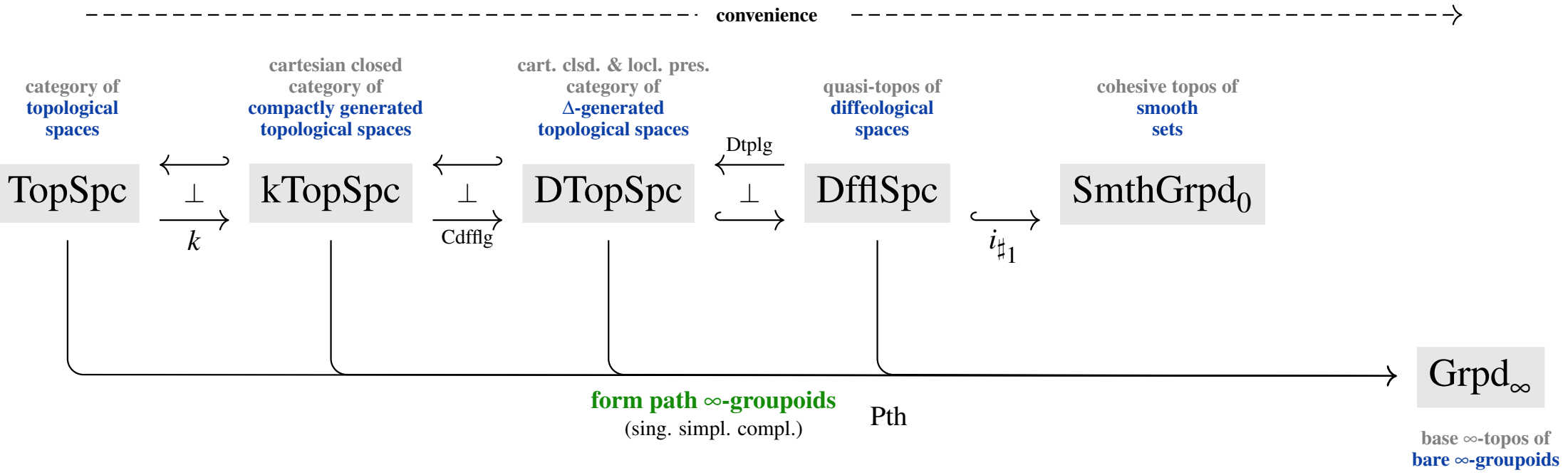
Shimakawa, Yoshida & Haraguchi (2010)

Christensen, Sinnamon & Wu (2014)

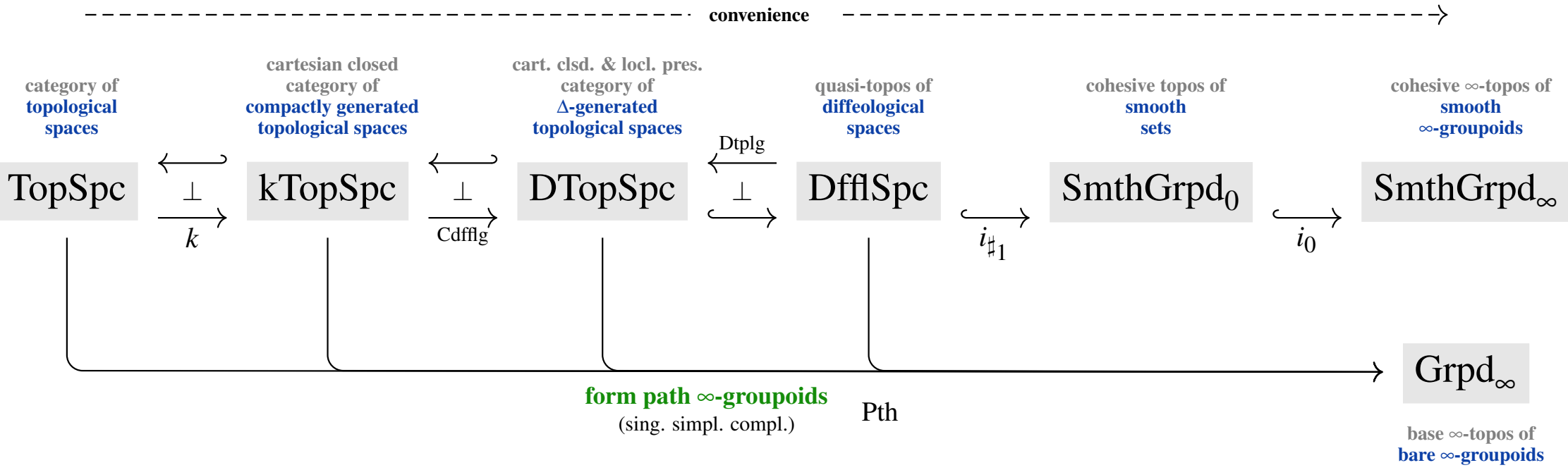
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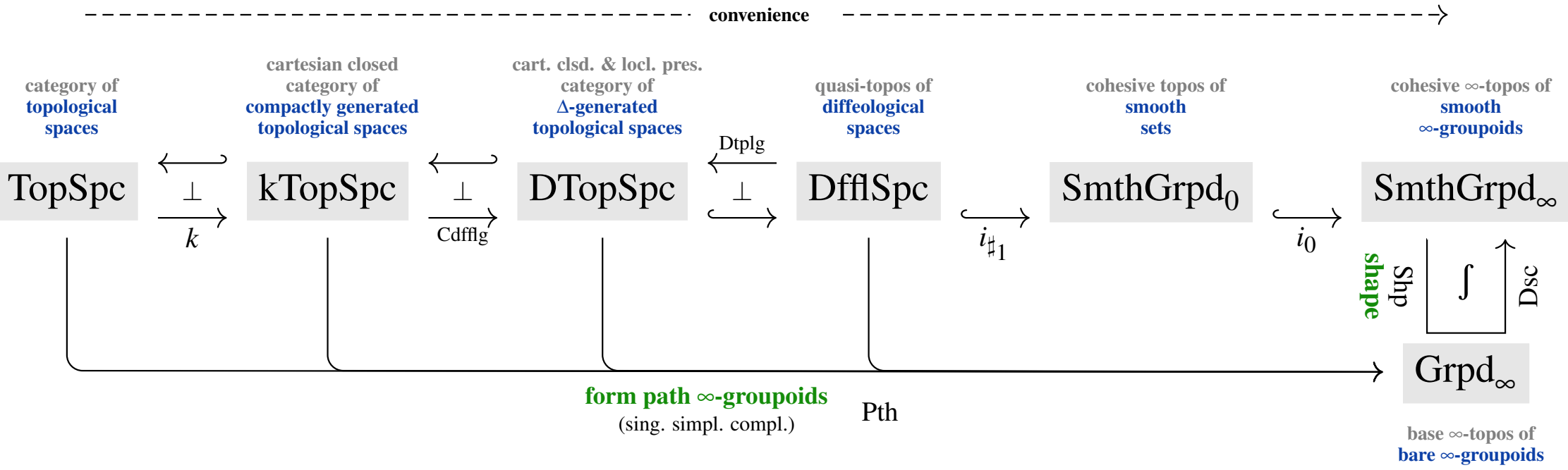
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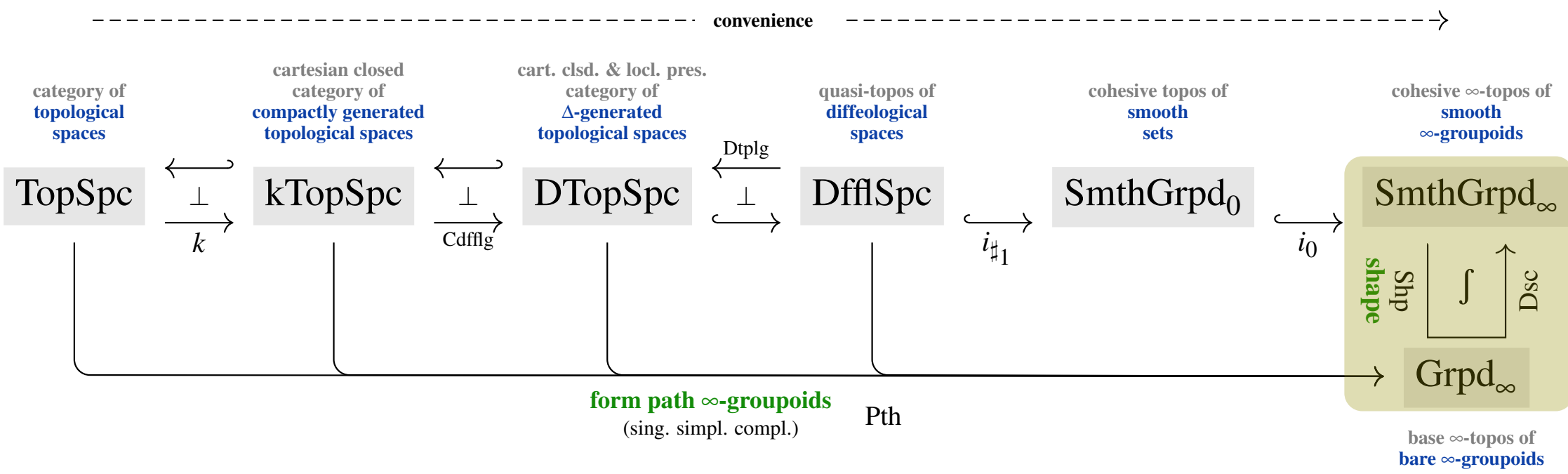
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Dugger (1998)

Cisinski (2003)

Schreiber (2013) – cohesive terminology borrowed from Lawvere (1994)

Pavlov (2014)

Pavlov et al. (2019)

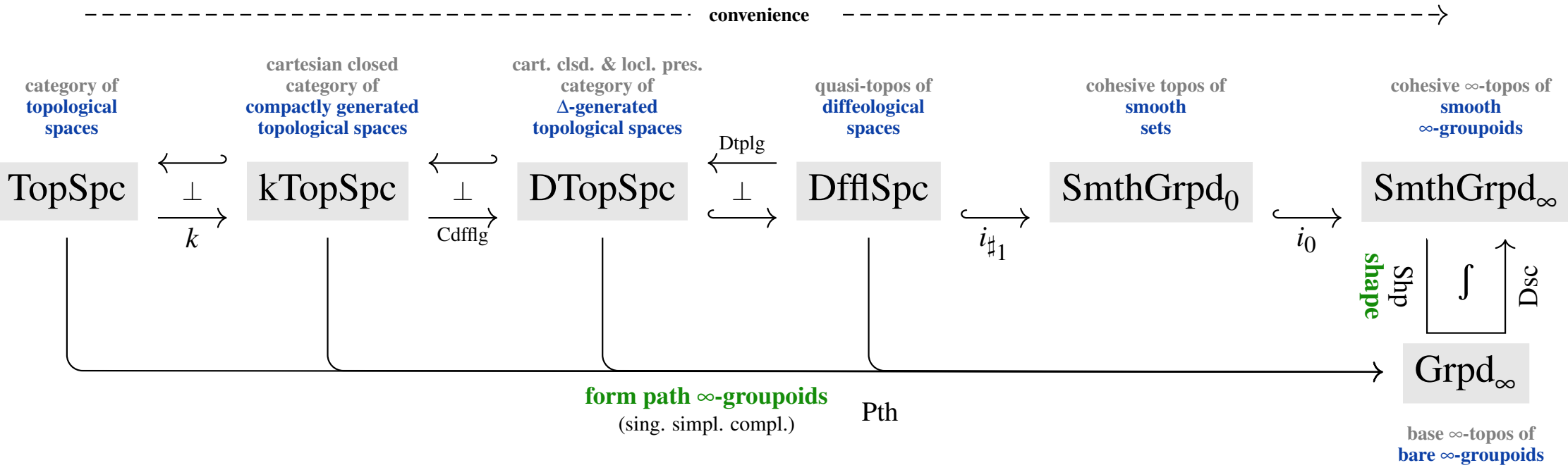
Bunk (2020)

Sati & Schreiber (2020)

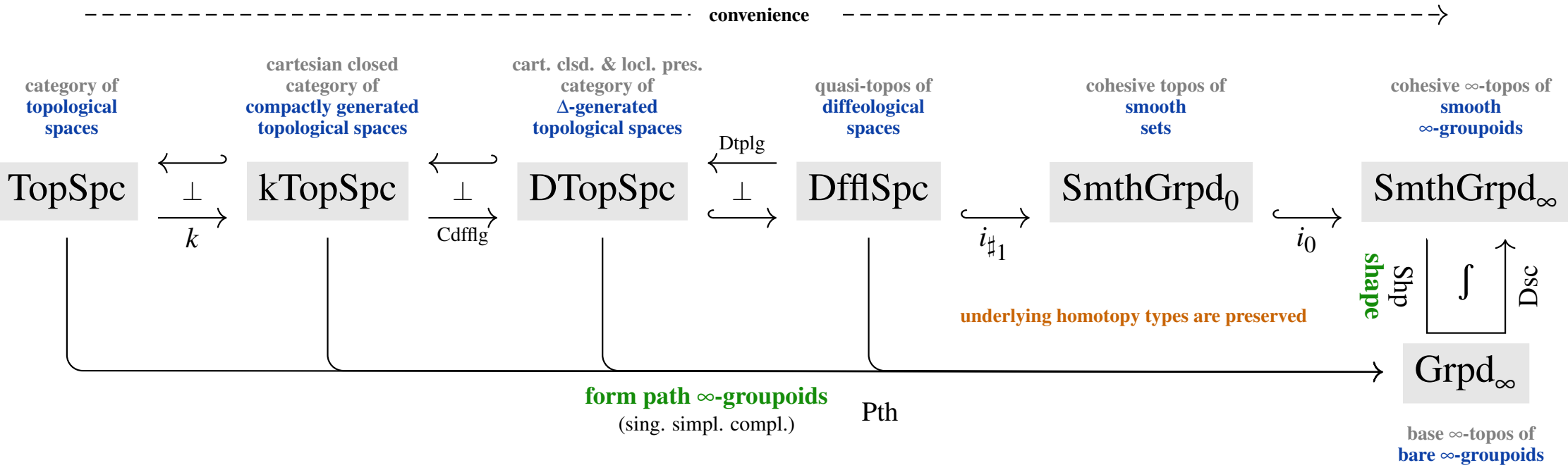
Sati & Schreiber (2021)

Clough (2021)

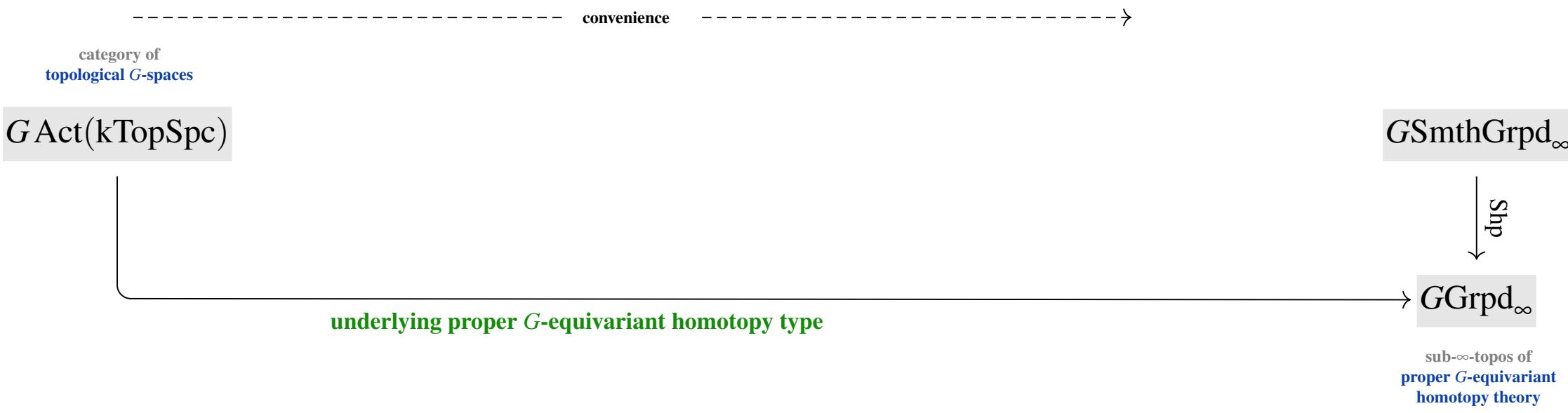
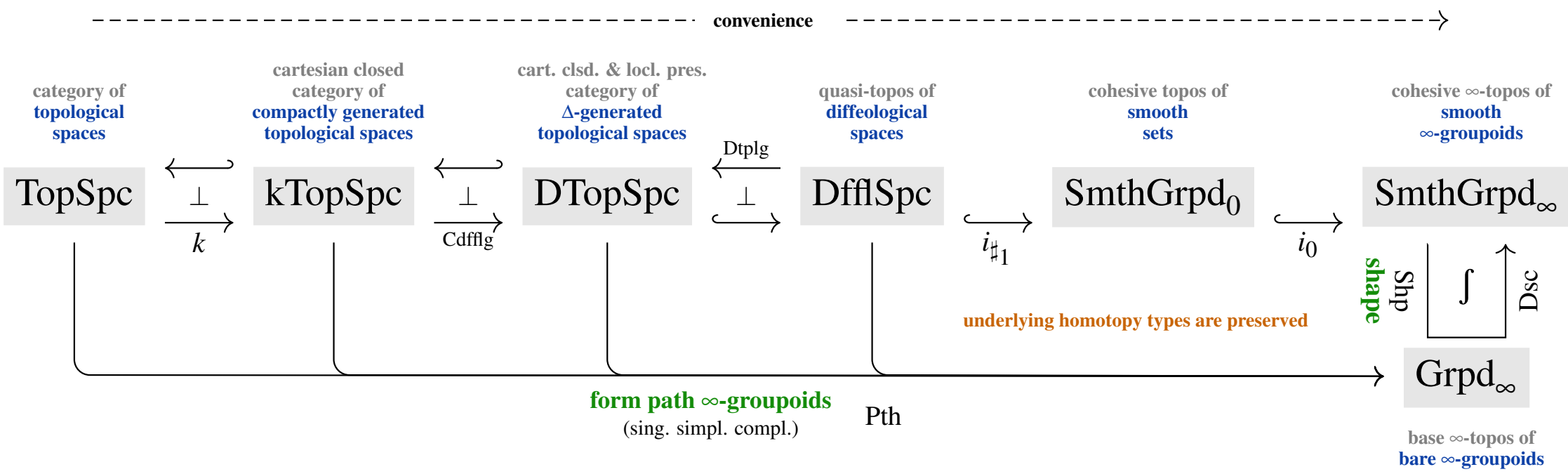
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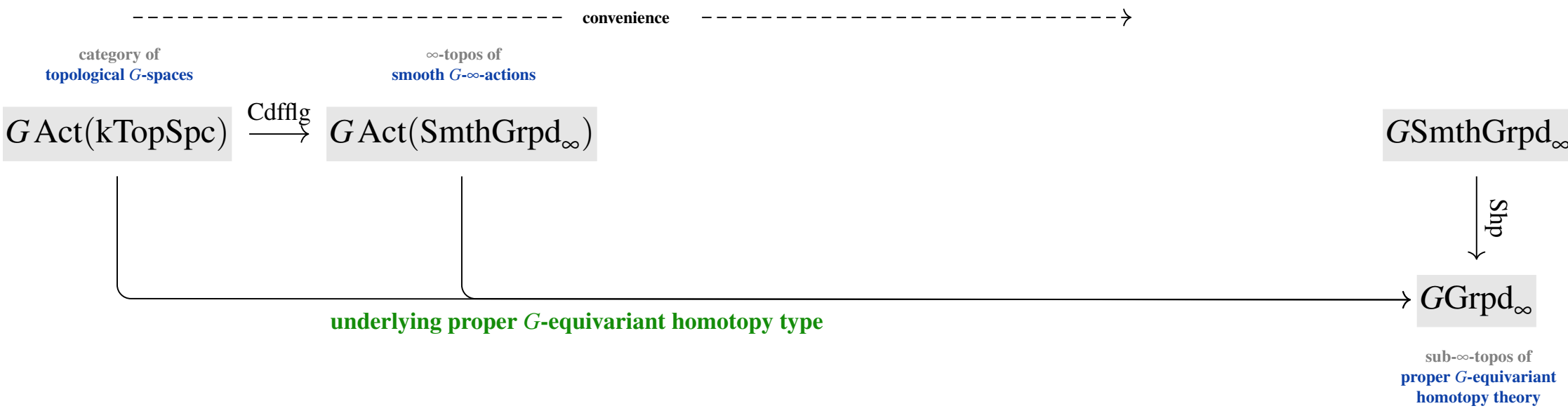
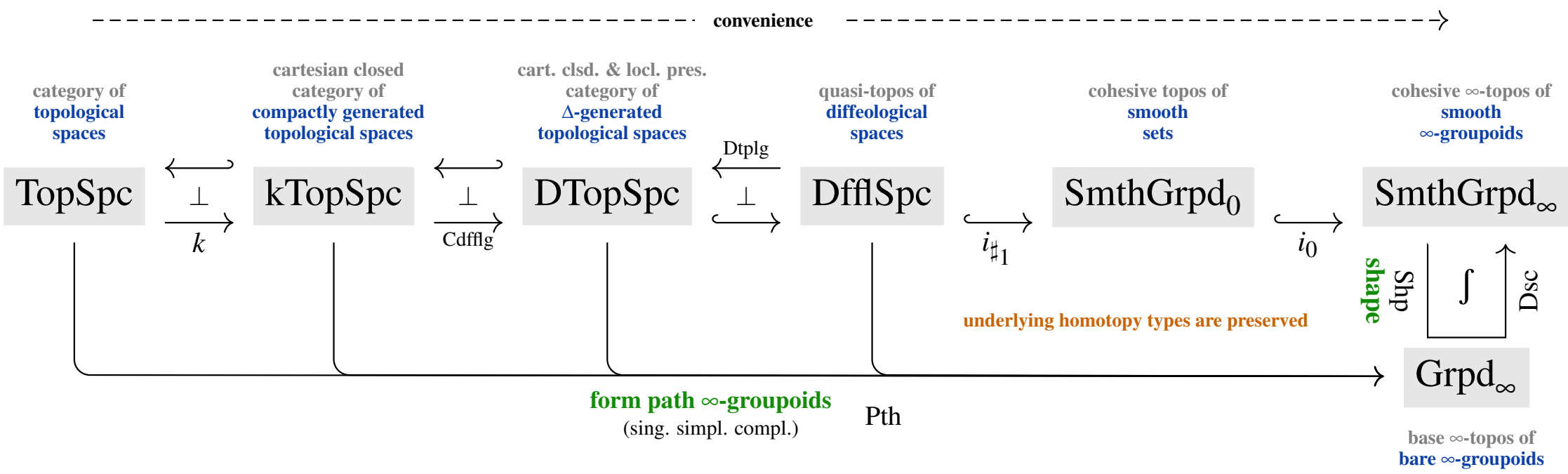
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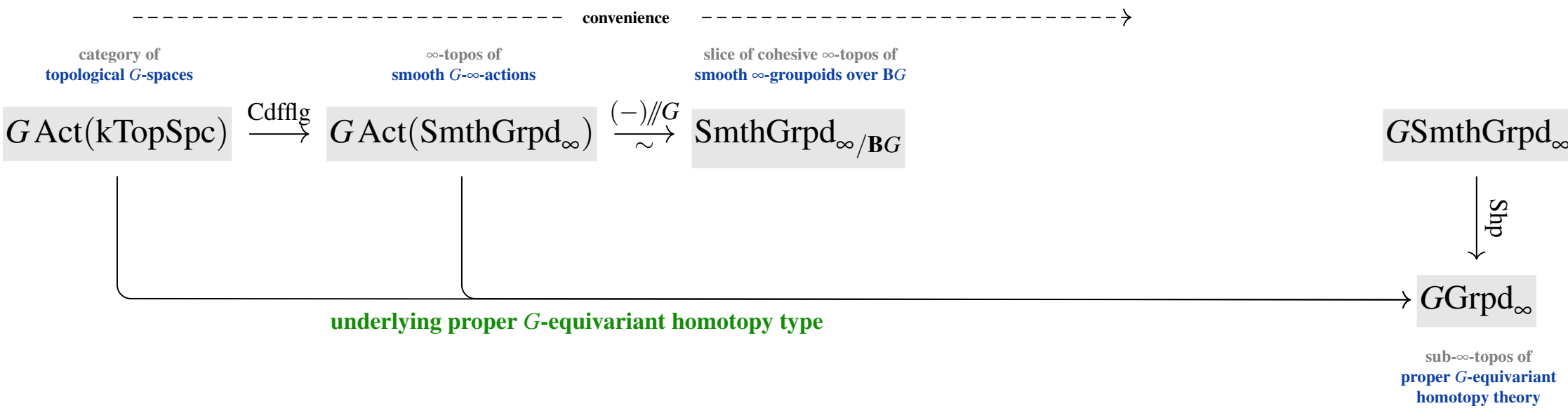
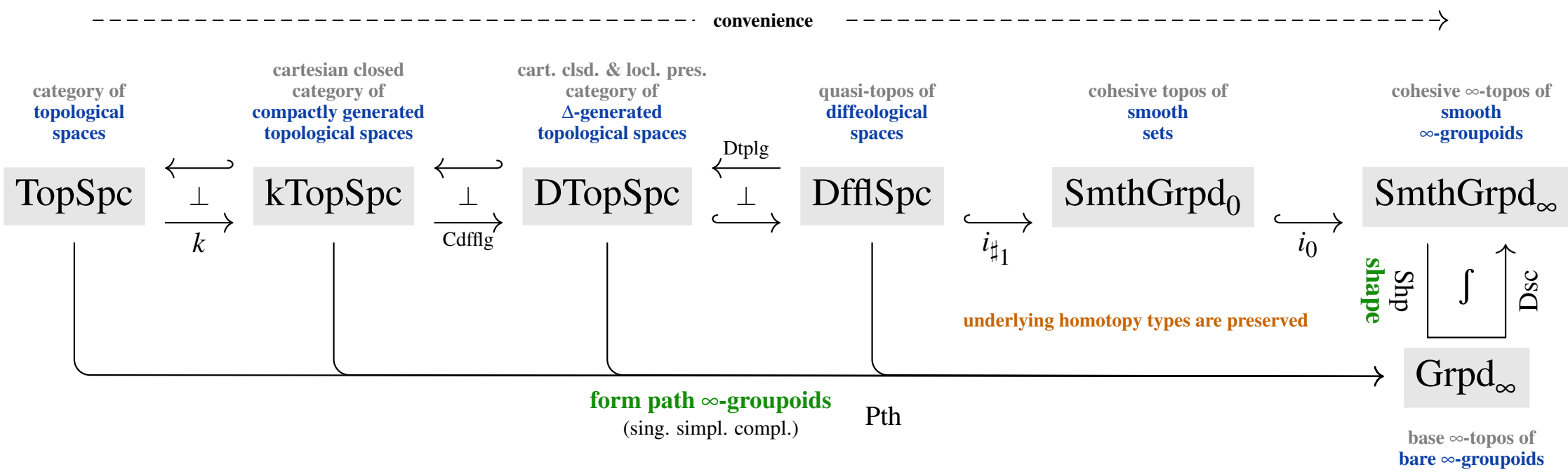
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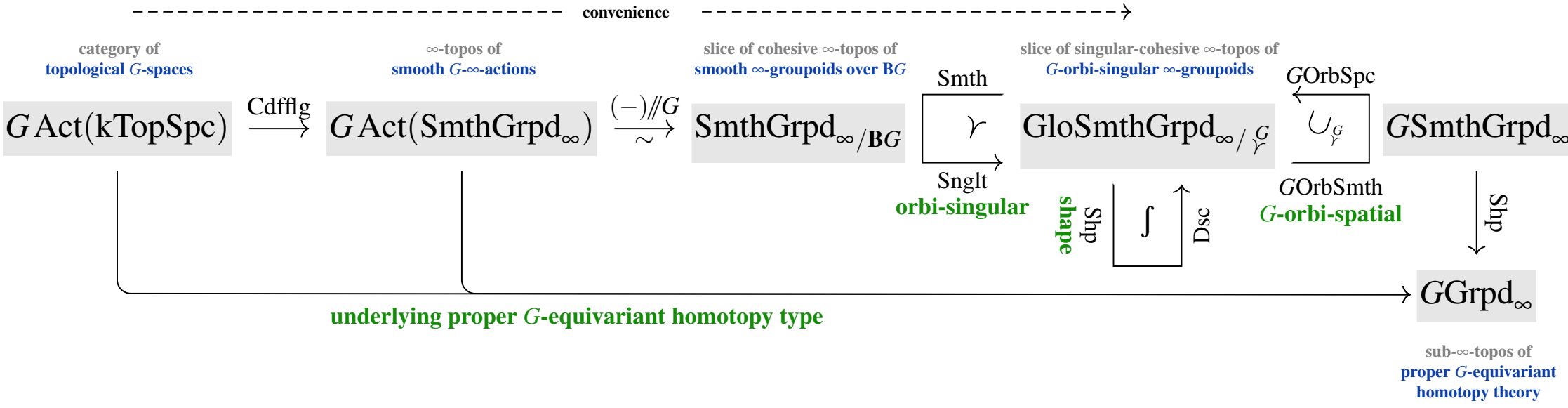
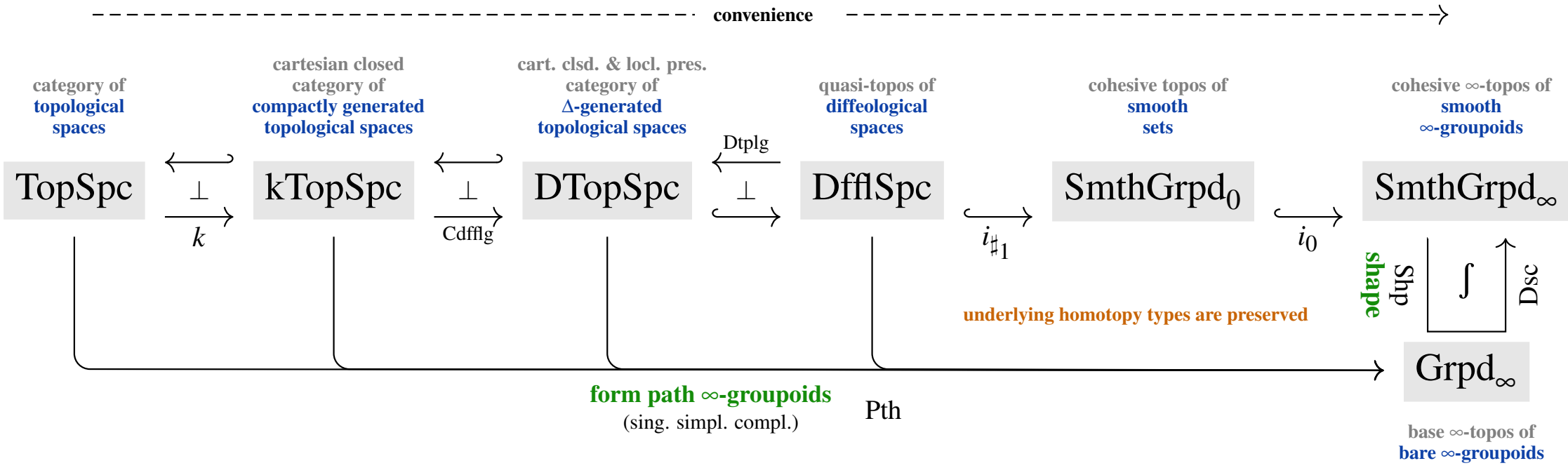
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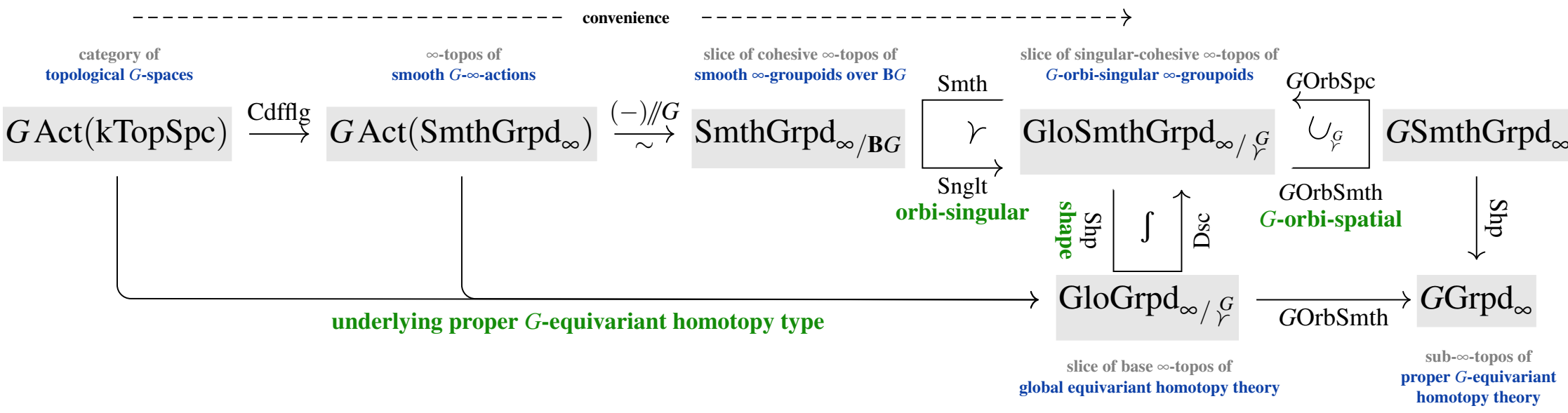
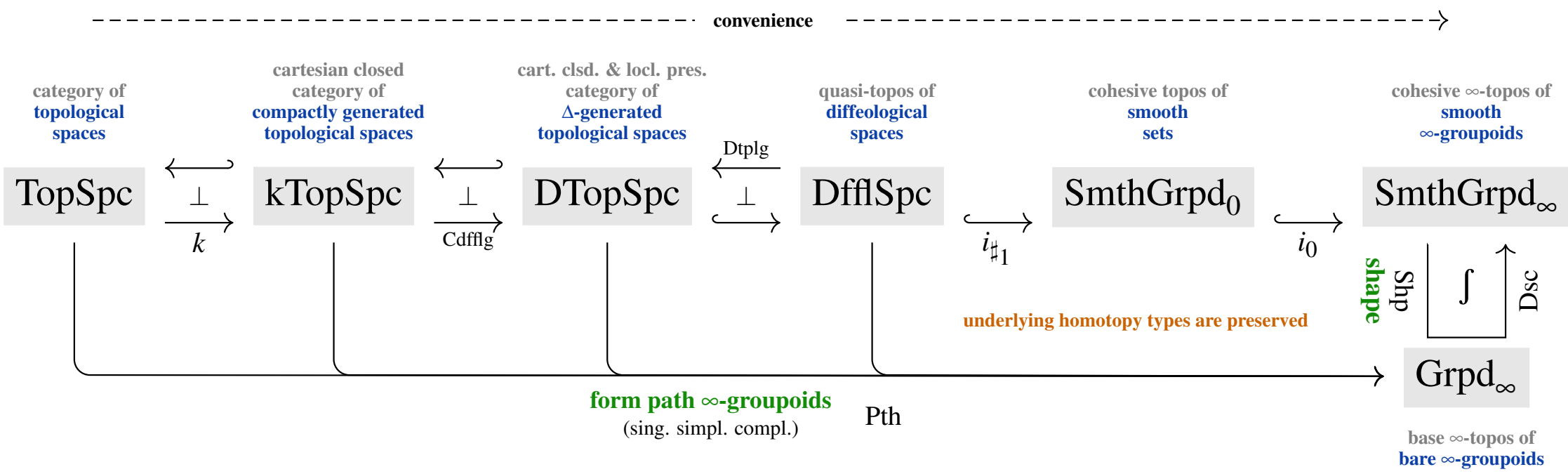
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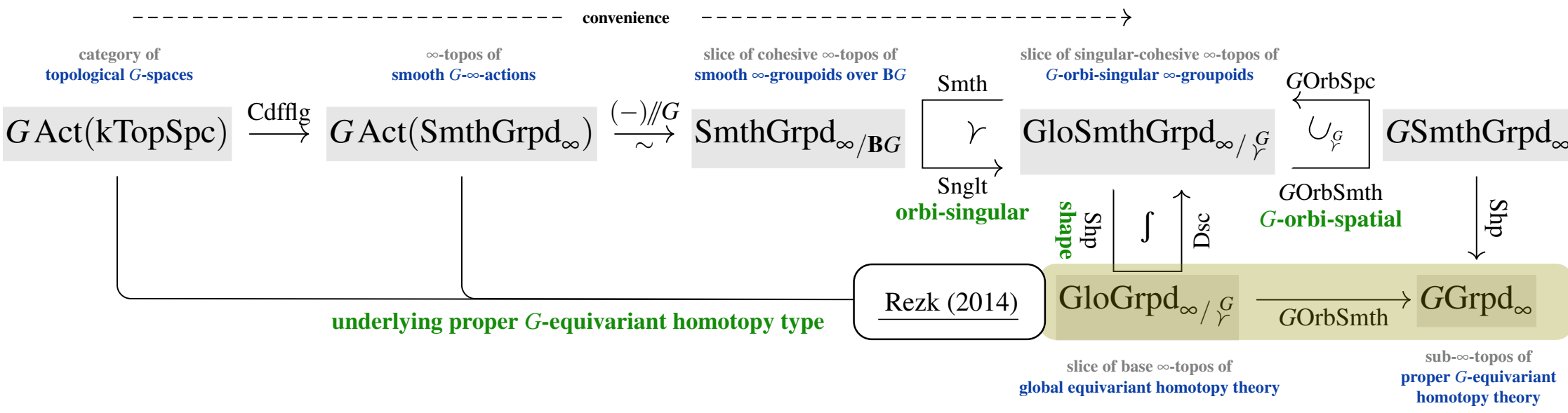
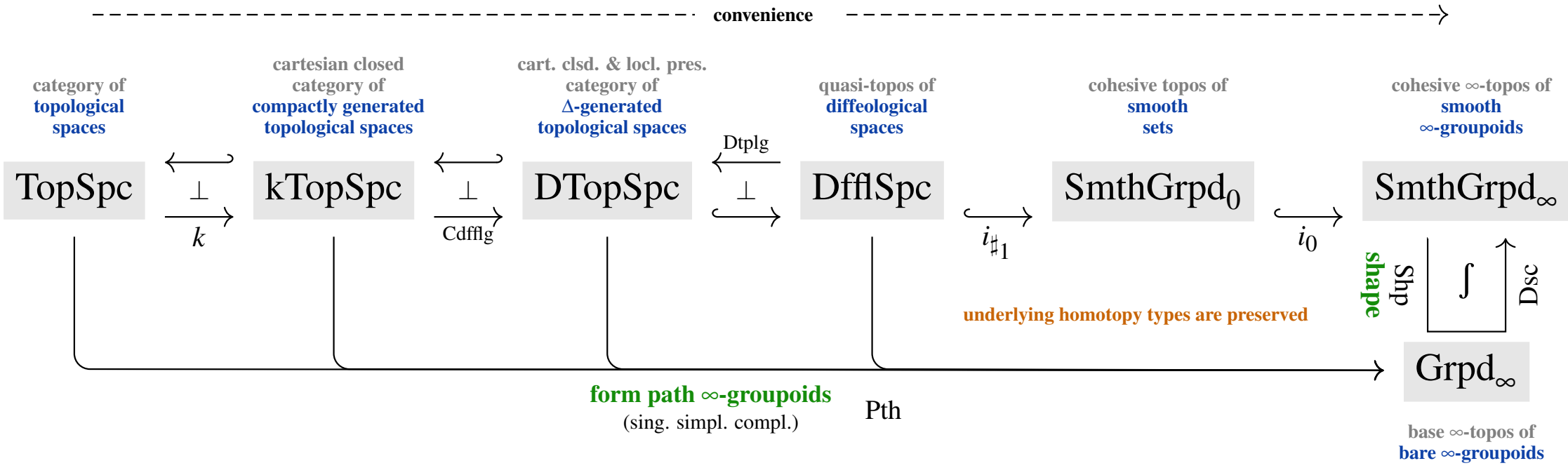
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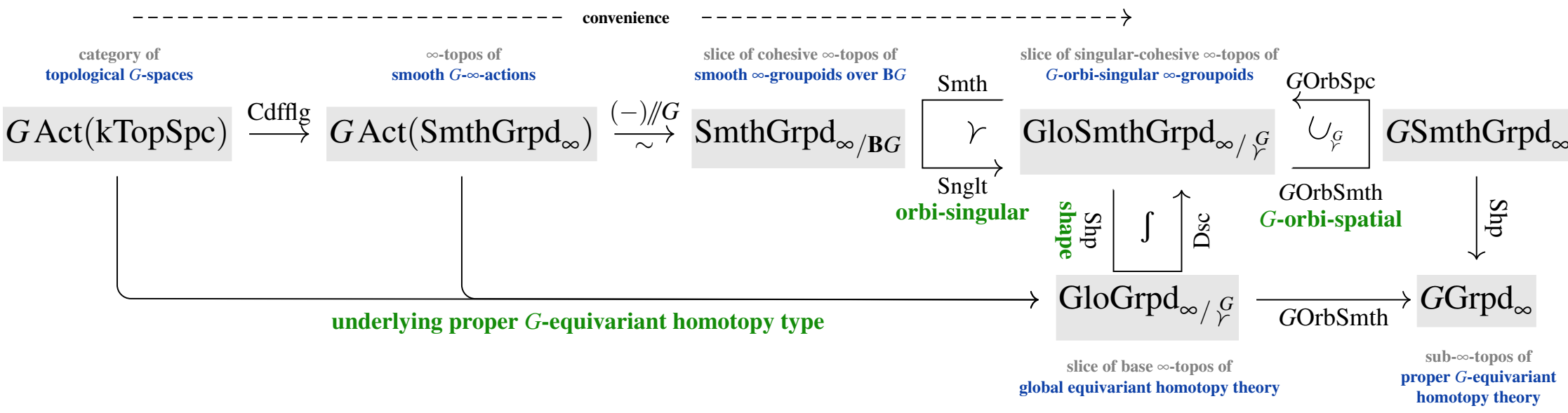
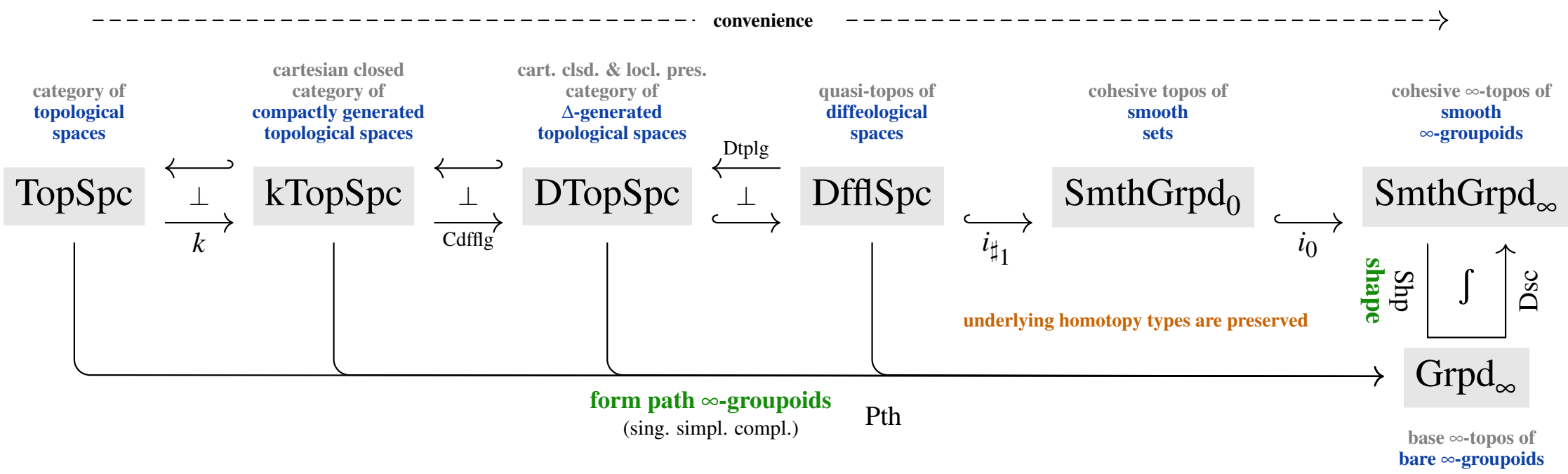
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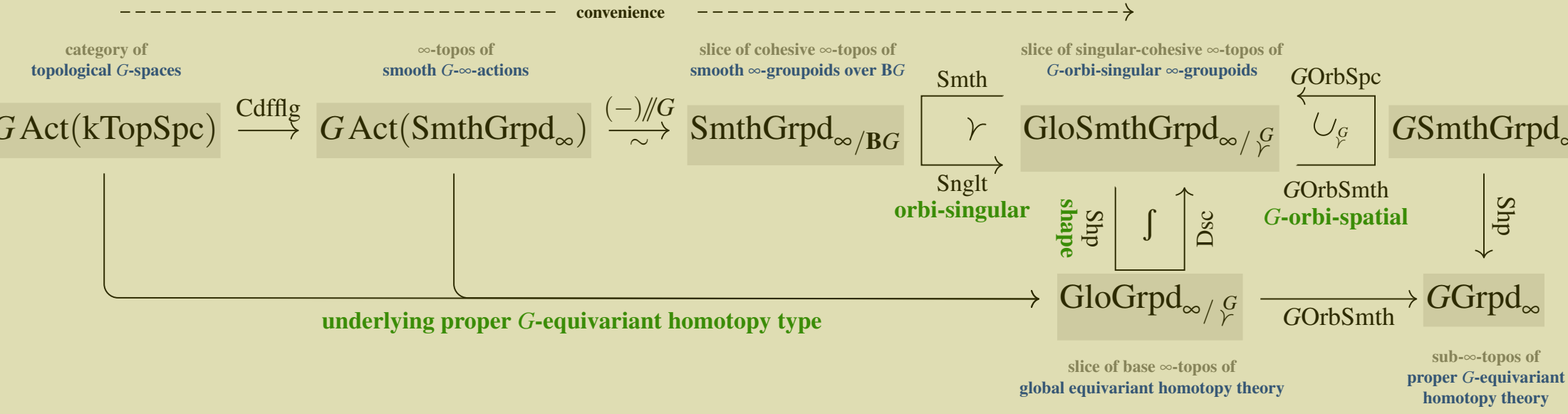
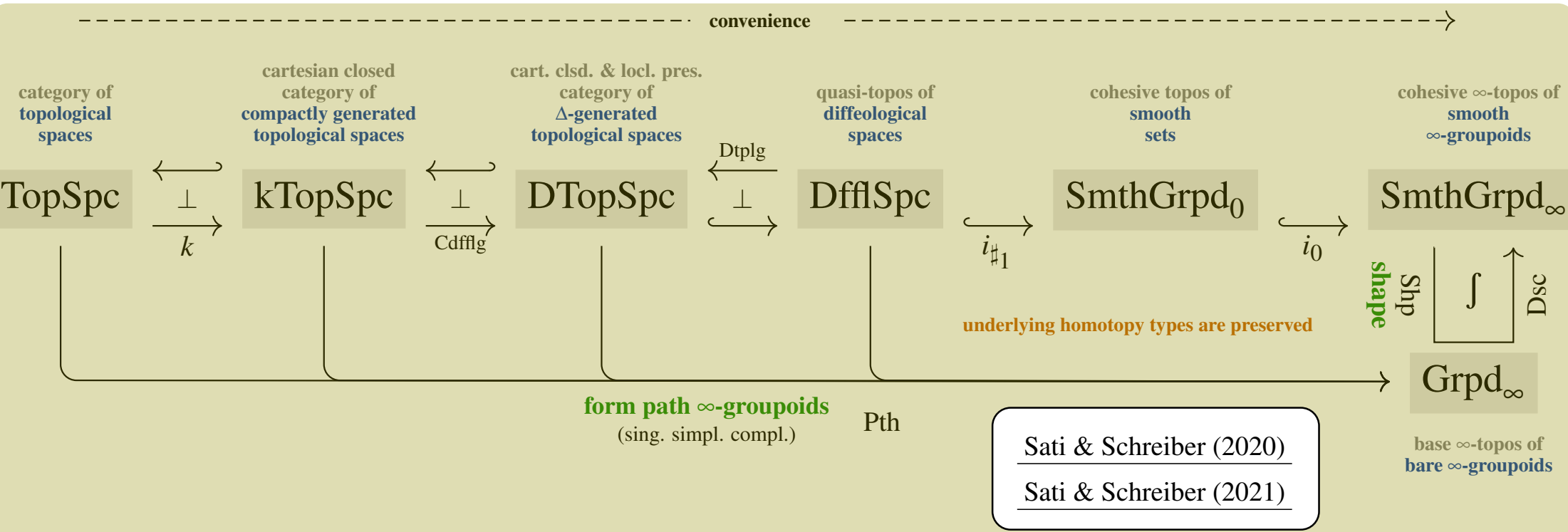
Diffeological Spaces in Perspective.



Diffeological Spaces in Perspective.



Diffeological Spaces in Perspective.



0 – Cohesive ∞ -Topoi

I – Equivariant ∞ -Bundles

II – TED-K-Theory

III – Anyonic Defect Branes

IV – Quantum Computation

This part is
a gentle exposition of
the most basic concept
underlying these articles:

<i>Principal ∞-bundles</i>	[arXiv:1207.0248/49]
<i>Equivariant Principal ∞-bundles</i>	[arXiv:2112.13654]
<i>Proper Orbifold Cohomology</i>	[arXiv:2008.01101]

following

<i>Diff. Cohomology in a Cohesive ∞-Topos</i>	[arXiv:1310.7930]
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Motivation, Overview, Summary and Outlook – in one single slide:

Generalized Cohomology Theories \leftrightarrow **Cohesive Higher Fiber Bundles**

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Here “generalized” subsumes “Whitehead-generalized cohomology” (\leftrightarrow spectra) but goes further:

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non-abelian	\leftrightarrow	general fibers
twisted	\leftrightarrow	associated
differential	\leftrightarrow	cohesive
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A major phenomenon/subtlety is that the last two aspects go hand-in-hand:

Proper G -equivariance corresponds to the **cohesive slice** over **BG**,
while

Borel equivariance corresponds just to the **slice of shapes**.

2-Groupoids

2-Groupoids are the algebra of 2-dimensional pasting,
such that all composition is associative and invertible:

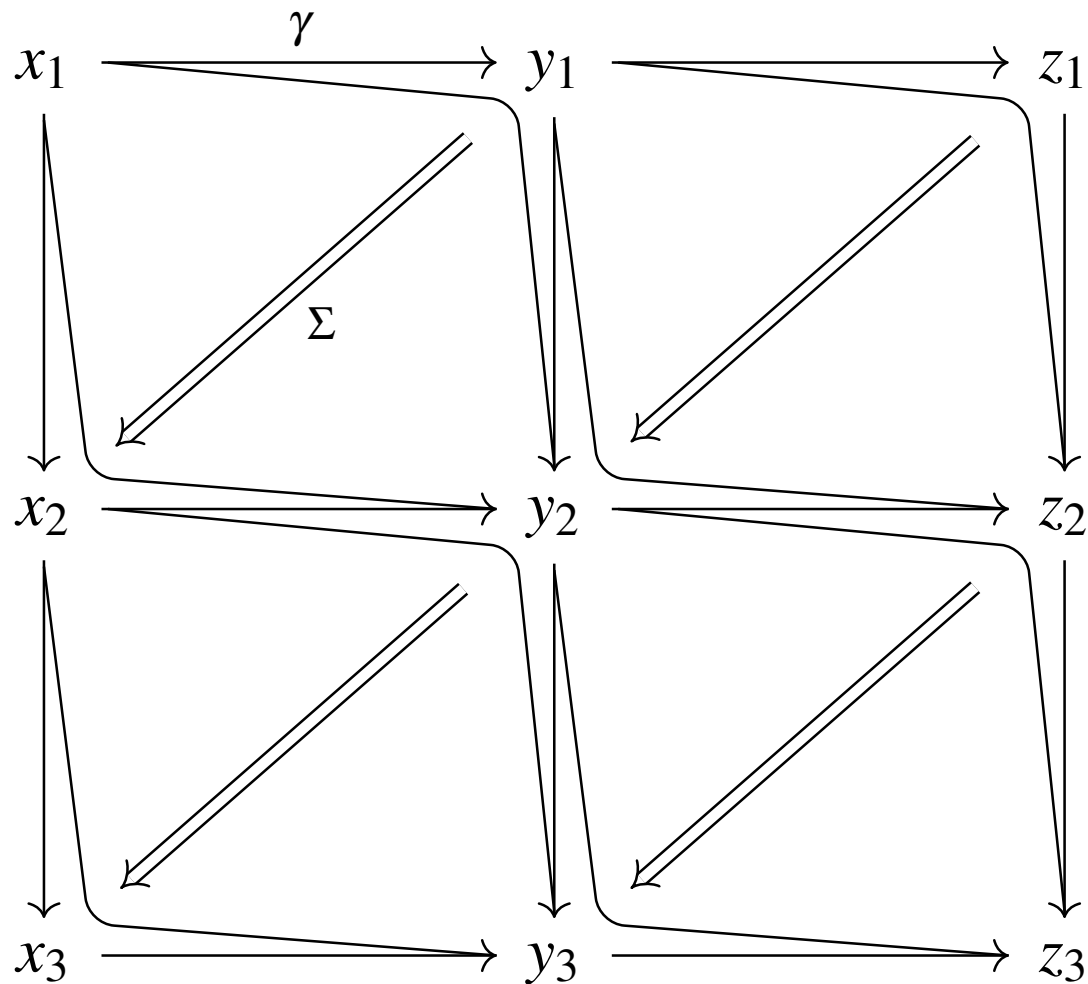
In general we need n -groupoids for $n \in \{1, 2, 3, \dots, \infty\}$
but for sake of exposition we may focus on $n = 2$.

2-Groupoids

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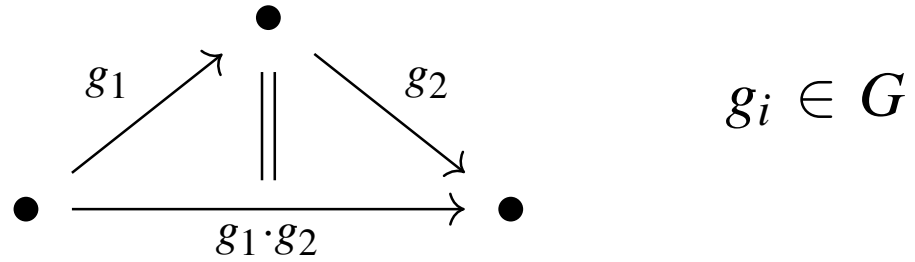
Accurate intuition:

homotopy classes of surfaces Σ relative boundary paths γ
in a topological space:



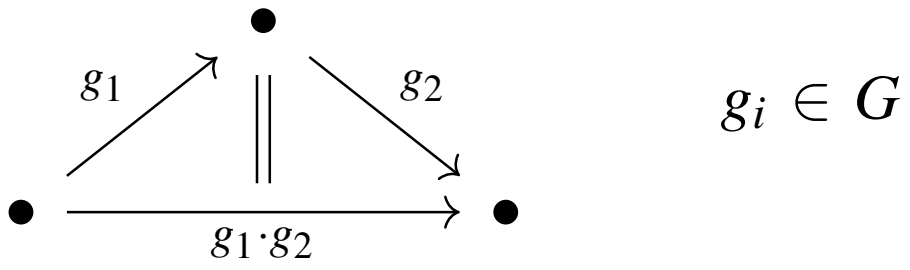
2-Groupoids – Examples.

For G a discrete group, there is its *delooping 1-groupoid* – $\boxed{\mathbf{BG}}$:

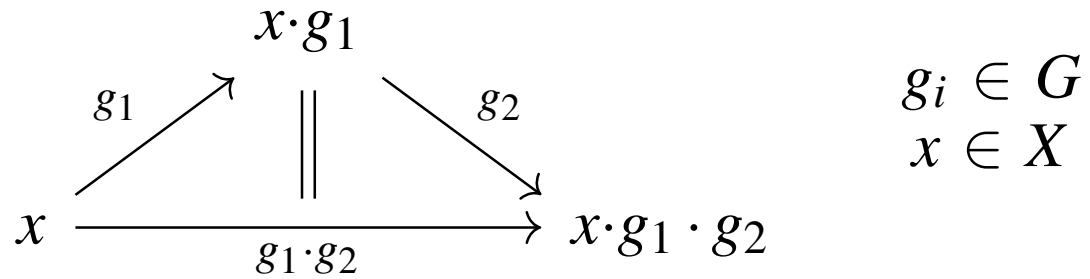


2-Groupoids – Examples.

For G a discrete group, there is its *delooping 1-groupoid* – \mathbf{BG} :



For $X \curvearrowright G$ a G -action on a set X there is its *action groupoid* or *homotopy quotient* – $X // S$:

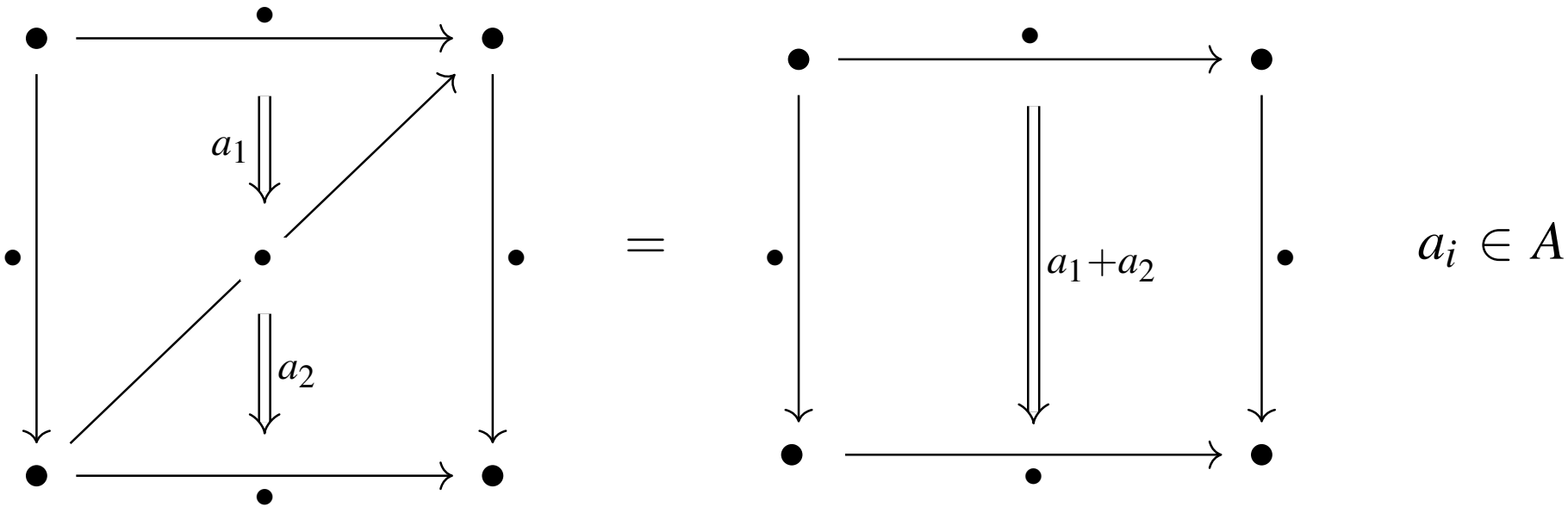


Hence: $\mathbf{BG} \simeq * // G$.

2-Groupoids – Examples.

For A an *abelian* group there is the *double delooping 2-groupoid*

$$\mathbf{B}^2 A = \mathbf{B} \left(\overset{\text{“2-group”}}{\mathbf{B}A} \right)$$

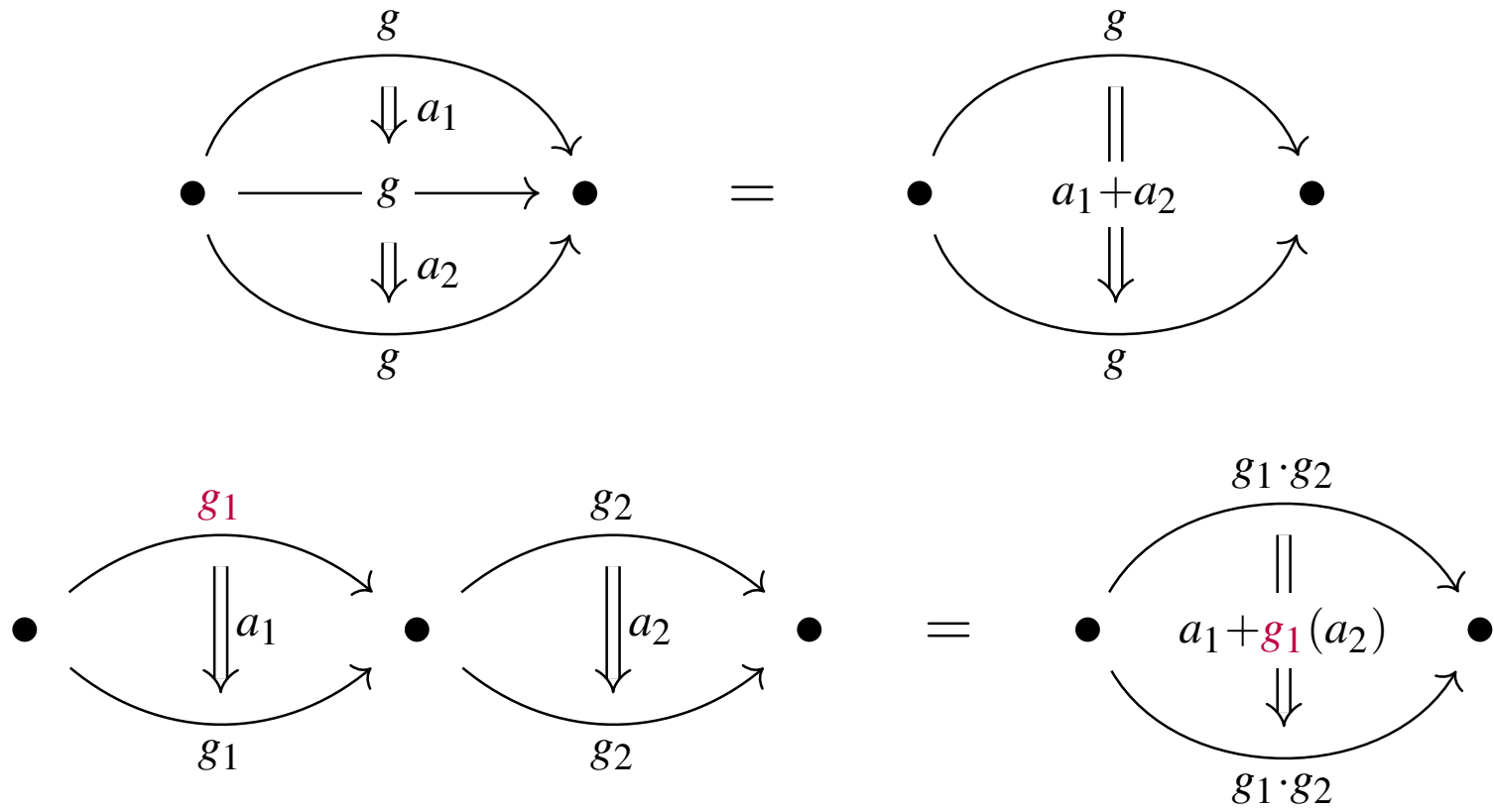


2-Groupoids – Examples.

For $A \curvearrowright G$ a *linear* action, i.e. by group automorphisms,

there is the delooping 2-groupoid $\mathbf{B}(\underbrace{(\mathbf{B}A) \rtimes G}) \simeq (\mathbf{B}^2A) // G$

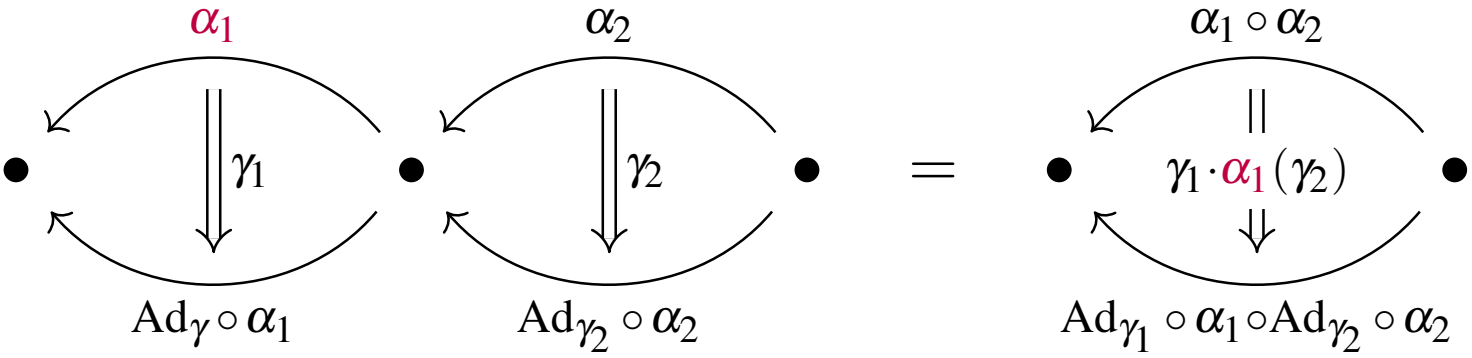
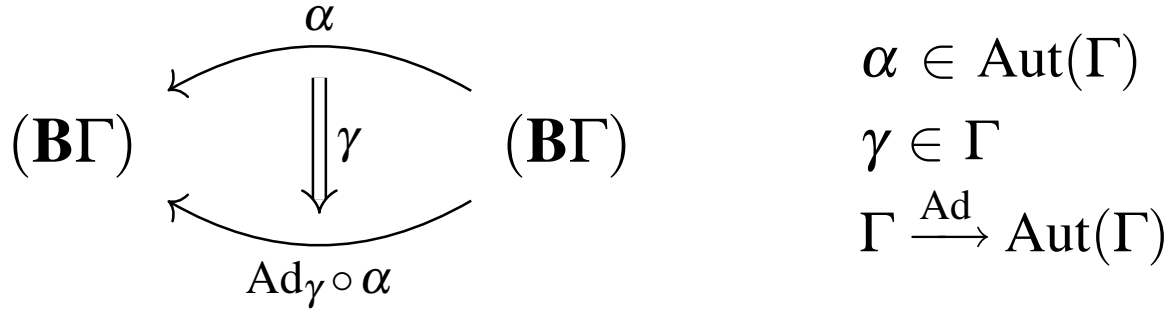
of the *semidirect product 2-group*:



2-Groupoids – Examples.

This is a special case of the delooping of the *automorphism 2-group* of a group Γ :

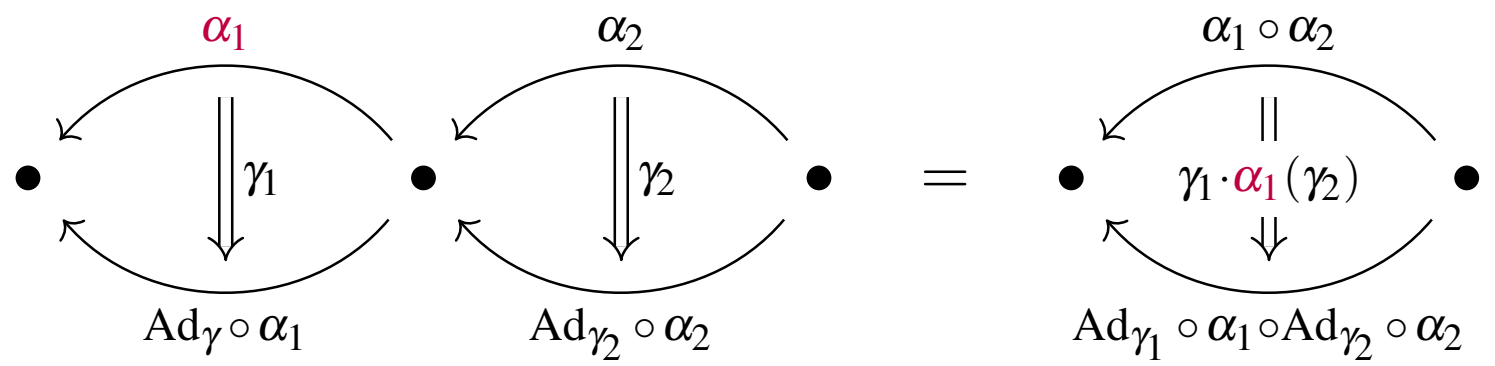
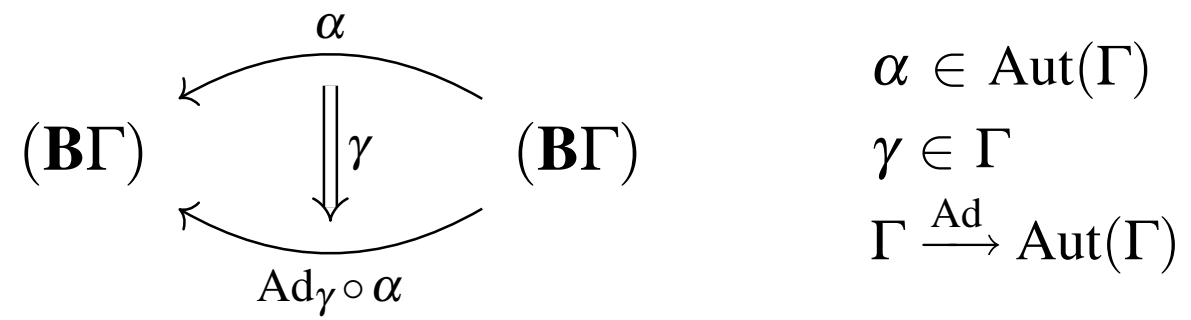
$$\mathbf{B}\left(\text{Aut}(\mathbf{B}\Gamma)\right) = \mathbf{B}\left(\overbrace{\text{Aut}(\Gamma) // \Gamma}\right)$$



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NB: Always need to choose whether actions are right- or left-actions, hence whether group multiplication is opposite or aligned to arrow composition. Before long we want *structure groups* to act *from the left* and *equivariance groups* to act *from the right*.

2-Groupoids – Examples.

Notice:

(1) $(\mathbf{BA}) \rtimes G$ is a non-abelian 2-group iff G is a non-abelian group;

2-Groupoids – Examples.

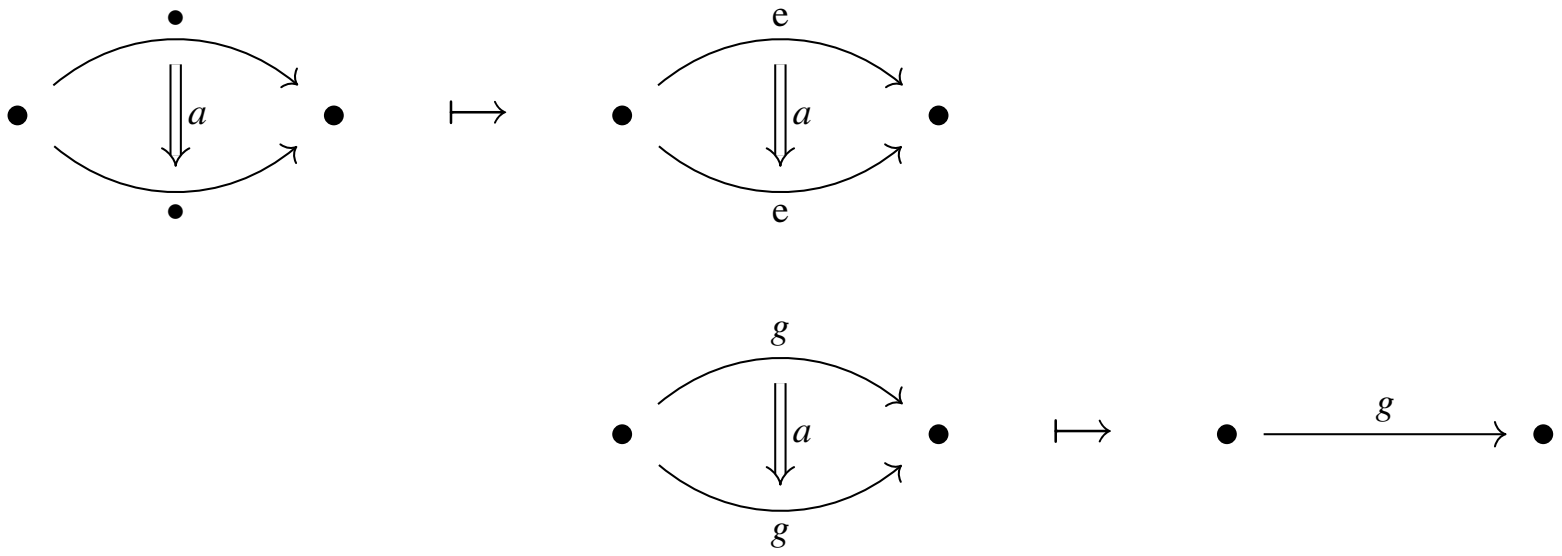
Notice:

(1) $(\mathbf{B}A) \rtimes G$ is a non-abelian 2-group iff G is a non-abelian group;

(2) its delooping sits in this fiber sequence:

$$\mathbf{B}^2 A \xrightarrow{\text{fib}(p)} \mathbf{B}((\mathbf{B}A) \rtimes G) \xrightarrow[p \in \text{KanFib}]{p} \mathbf{B}G$$

\parallel
 $(\mathbf{B}^2 A) // G$



2-Groupoids – 2-Functors.

A *2-functor* is a map between 2-groupoids respecting identities and composition.

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E.g.: if $\mathbb{Z} \curvearrowright \mathbb{Z}_2$ by sign inversion, and $G \xrightarrow{\sigma} \mathbb{Z}_2$ a homomorphism then

2nd group cohomology of G with coefficients in $G \curvearrowright \mathbb{Z}$ is 2-functors:

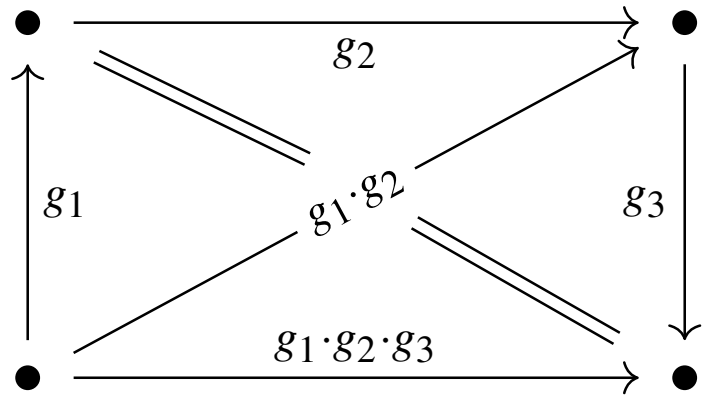
$$\begin{array}{ccc}
 \mathbf{B}G & \overset{\text{2-functor} = \text{2-cocycle}}{\dashrightarrow} & (\mathbf{B}^2 A) // \mathbb{Z}_2 \\
 \searrow \mathbf{B}\sigma & & \swarrow \\
 & \mathbf{B}\mathbb{Z}_2 &
 \end{array}$$

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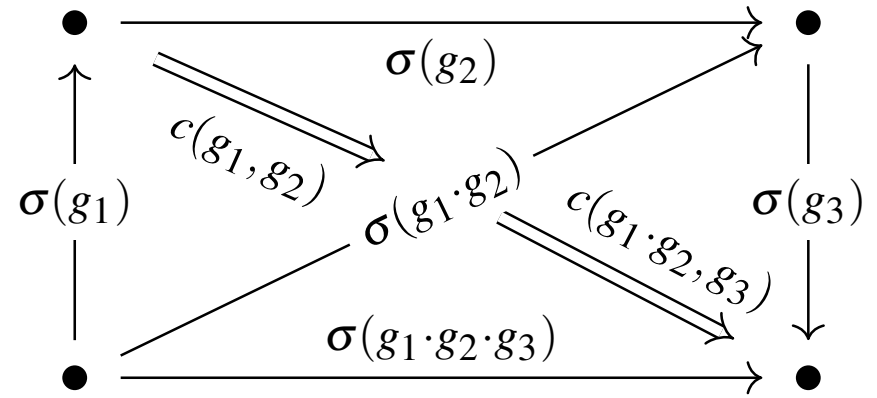
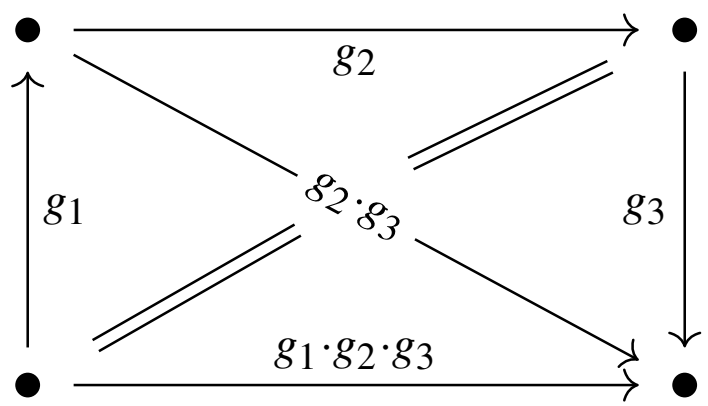
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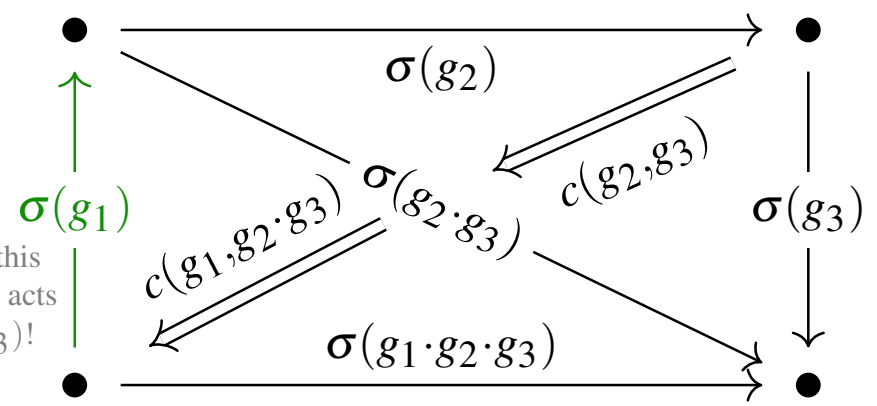
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 \end{array}$$



||| associativity



||| cocycle condition



recall that this 1-morphism acts on $c(g_2, g_3)$!

2-Groupoids with smooth structure.

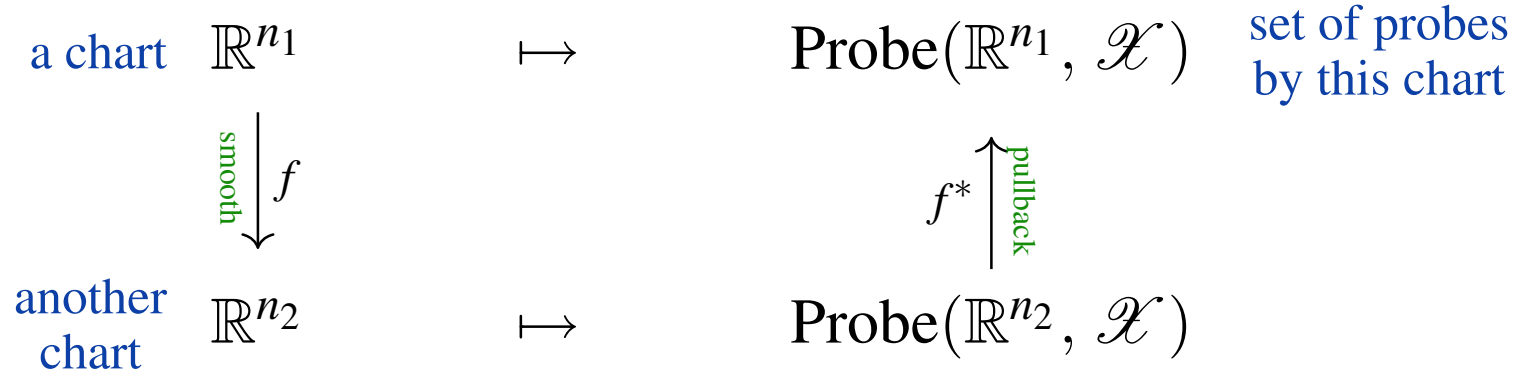
A *smooth 2-groupoid* \mathcal{X} is given by a rule

which to each chart \mathbb{R}^n , $n \in \mathbb{N}$, assigns the plain 2-groupoid $\text{Probe}(\mathbb{R}^n, \mathcal{X})$

of ways of smoothly mapping \mathbb{R}^n into the would-be \mathcal{X}

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$$\begin{array}{ccc} \text{a chart } \mathbb{R}^{n_1} & \mapsto & \text{Probe}(\mathbb{R}^{n_1}, \mathcal{X}) \quad \text{set of probes by this chart} \\ \text{smooth } \downarrow f & & f^* \uparrow \text{pullback} \\ \text{another chart } \mathbb{R}^{n_2} & \mapsto & \text{Probe}(\mathbb{R}^{n_2}, \mathcal{X}) \end{array}$$

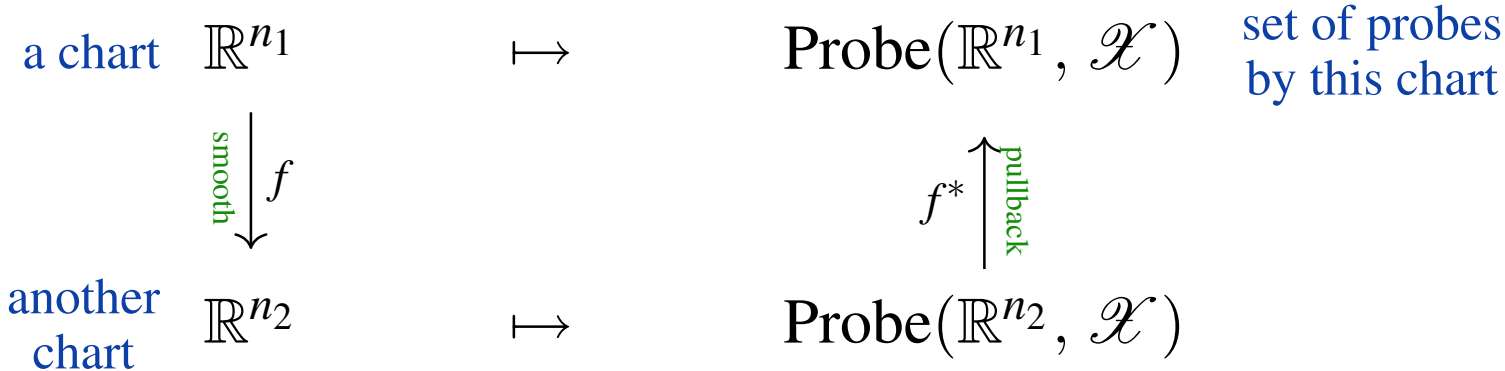
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So $\text{Probe}(*, \mathcal{X}) = \text{Probe}(\mathbb{R}^0, \mathcal{X})$ is the underlying 2-groupoid

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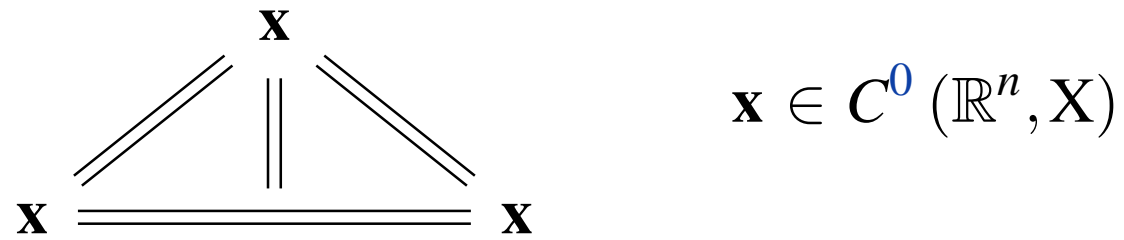
Grothendieck (1965): “functorial geometry”

common jargon: “pre-2-stacks on the site of Cartesian spaces”

2-Groupoids with smooth structure – Examples.

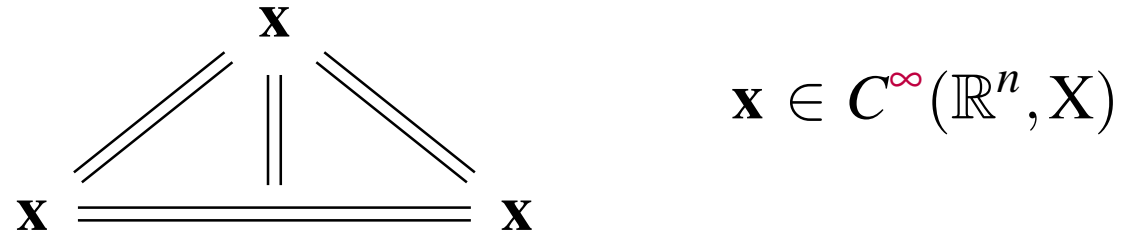
If X is a **topological space**, then as a smooth 2-groupoid it's this assignment:

$$X : \mathbb{R}^n \mapsto \text{Probe}(\mathbb{R}^n, X) := C^0(\mathbb{R}^n, X)$$



If X is a **smooth manifold**, then as a smooth 2-groupoid it's this assignment:

$$X : \mathbb{R}^n \mapsto \text{Probe}(\mathbb{R}^n, X) := C^\infty(\mathbb{R}^n, X)$$

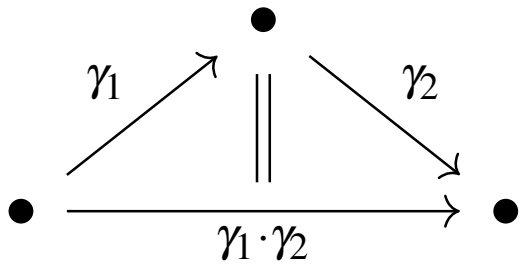


(also known as X in its incarnation as a *diffeological space*).

2-Groupoids with smooth structure – Examples.

If Γ a *Lie* group, then the sets of smooth functions $C^\infty(\mathbb{R}^n, \Gamma)$ are plain groups, and the *smooth delooping groupoid* $\mathbf{B}\Gamma$ is:

$$\mathbf{B}\Gamma : \mathbb{R}^n \mapsto \text{Probe}(\mathbb{R}^n, \mathbf{B}\Gamma) := \mathbf{B}(C^\infty(\mathbb{R}^n, \Gamma))$$

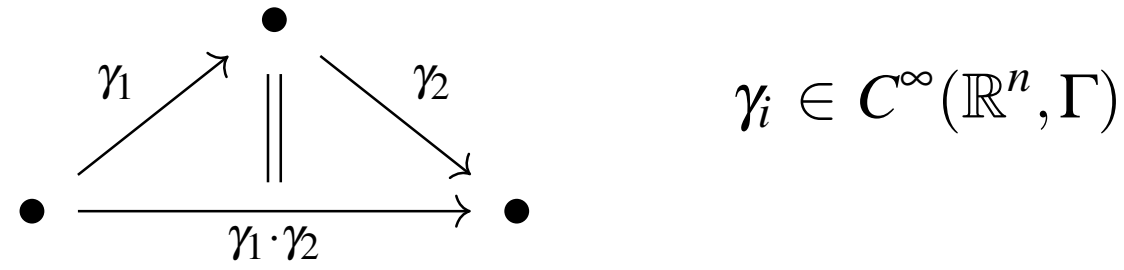


$$\gamma_i \in C^\infty(\mathbb{R}^n, \Gamma)$$

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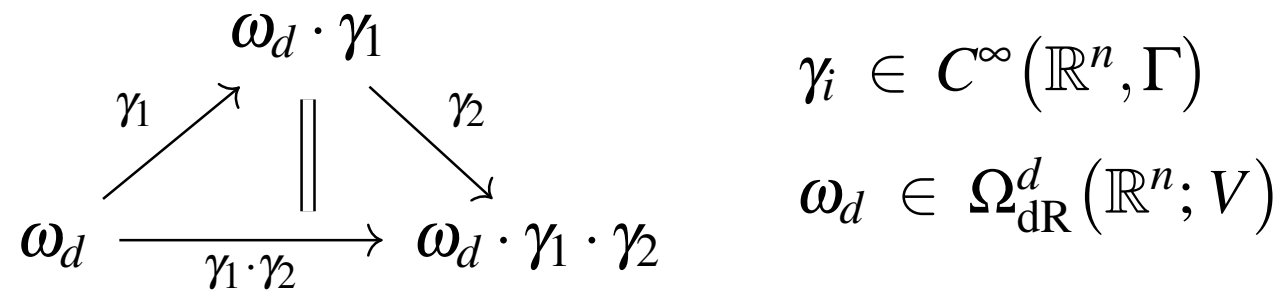
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If V is a Γ -representation, then the *smooth moduli space of V -valued differential forms* is

$$\Omega_{\text{dR}}^d(-; V) // \Gamma : \mathbb{R}^n \mapsto \Omega_{\text{dR}}^d(\mathbb{R}^n; V) // \Gamma$$



2-Groupoids with smooth structure – As smooth homotopy types.

Smooth 2-groupoids are *models* for *smooth 2-stacks* aka **smooth homotopy 2-types**

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A smooth 2-functor $\mathcal{X} \xrightarrow{f} \mathcal{Y}$ is called:

PrjFib	<i>projective fibration</i>	iff for each \mathbb{R}^n , every $k + 1$ -morphism in $\text{Probe}(\mathbb{R}^n, \mathcal{Y})$ that starts at k -morphisms which come from $\text{Probe}(\mathbb{R}^n, \mathcal{X})$ lifts compatibly to a $k + 1$ -morphism in $\text{Probe}(\mathbb{R}^n, \mathcal{X})$
LWEq	<i>local weak equivalence</i>	iff for every \mathbb{R}^n there exists an open ball $0 \in \mathbb{D}_\varepsilon^n \xrightarrow{i} \mathbb{R}^n$ such that $\text{Probe}(\mathbb{R}^n, f) _i$ is a weak homotopy equivalence namely an iso on the evident homotopy groups
PrjCof	<i>projective cofibration</i>	if (Dugger's sufficient condition): for all k , the spaces of k -morphisms are disjoint unions of charts \mathbb{R}^n (for any n -s)

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Fact/Def.: *Maps of 2-stacks* are modeled by

$$\emptyset \xrightarrow[\in \text{PrjCof}]{\text{cof. domain}} \widehat{\mathcal{X}} \xrightarrow[\text{smooth 2-functor}]{\phi} \mathcal{Y} \xrightarrow[\in \text{PrjFib}]{\text{fib. co-domain}} *$$

2-Groupoids with smooth structure – As smooth homotopy types.

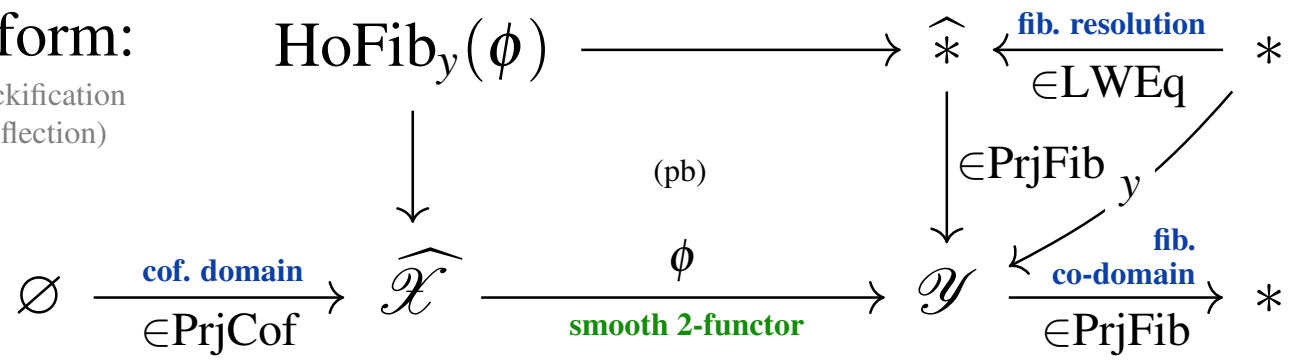
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Fact/Def.: *Maps of 2-stacks* and their *homotopy fibers* are modeled by pullbacks of this form:

(because 2-stackification is an ∞ -lex reflection)



2-Groupoids with smooth structure – As smooth homotopy types.

Smooth 2-groupoids are *models* for *smooth 2-stacks* aka **smooth homotopy 2-types**

A smooth 2-functor $\mathcal{X} \xrightarrow{f} \mathcal{Y}$ is called:

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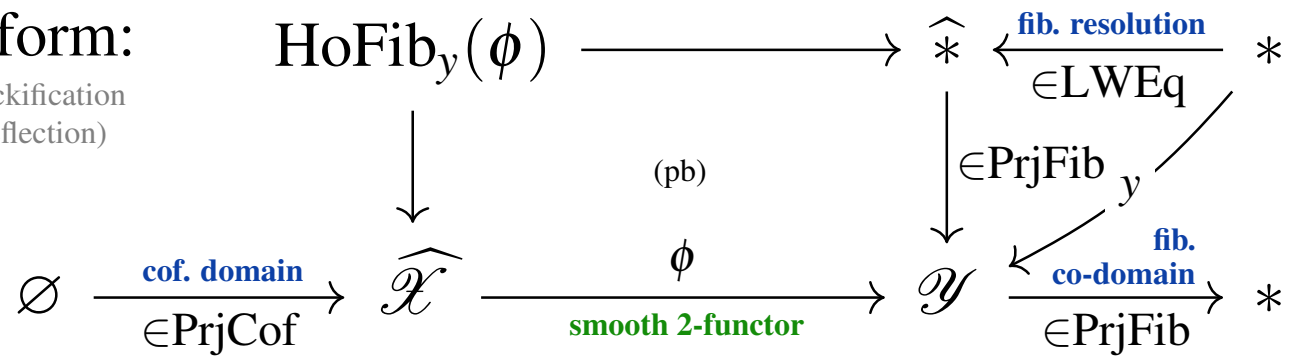
Definition/Example:
 For $\widehat{\mathcal{X}}$ projectively cofibrant and \mathcal{Y} projectively fibrant the smooth *mapping 2-stack* is modeled by

$$\text{Map}(\widehat{\mathcal{X}}, \mathcal{Y}) : \mathbb{R}^n \mapsto \text{Hom}(\widehat{\mathcal{X}} \times \mathbb{R}^n, \mathcal{Y})$$

2-groupoid of smooth 2-functors

Fact/Def.: *Maps of 2-stacks and their homotopy fibers are modeled by pullbacks of this form:*

(because 2-stackification is an ∞ -lex reflection)



2-Groupoids with smooth structure – Homotopy fiber sequences.

$$\begin{array}{ccccccc}
 \mathbf{U}(1) & \hookrightarrow & \Gamma & \twoheadrightarrow & \Gamma/\mathbf{U}(1) & \text{smooth group} & \\
 & & \text{locally trivial} & & & & \\
 & & \text{circle-extension} & & \uparrow \in \text{LWEq} & & \\
 & & & & \Gamma//\mathbf{U}(1) & \xrightarrow{\in \text{PrjFib}} & \mathbf{BU}(1) \longrightarrow \mathbf{B}\Gamma \longrightarrow \mathbf{B}(\Gamma/\mathbf{U}(1)) \\
 & & & & & & \uparrow \in \text{LWEq} \\
 & & & & & & \mathbf{B}\Gamma//\mathbf{BU}(1) \xrightarrow{\text{2-cocyle}} \mathbf{B}^2\mathbf{U}(1) \\
 & & & & & & \text{classifying} \\
 & & & & & & \text{the extension}
 \end{array}$$

2-Groupoids with smooth structure – Homotopy fiber sequences.

$$U(1) \hookrightarrow \Gamma \twoheadrightarrow \Gamma/U(1) \text{ smooth group}$$

locally trivial
circle-extension

$$\uparrow \in \text{LWEq}$$

$$\Gamma//U(1) \xrightarrow{\in \text{PrjFib}} \mathbf{BU}(1) \longrightarrow \mathbf{B}\Gamma \longrightarrow \mathbf{B}(\Gamma/U(1))$$

$$\uparrow \in \text{LWEq}$$

$$\begin{array}{ccc} c & \gamma & [\gamma] \\ \parallel & \parallel & \parallel \\ c & \gamma & [\gamma] \end{array}$$

$$\mathbf{B}\Gamma//\mathbf{BU}(1) \xrightarrow{\text{2-cocyle classifying the extension}} \mathbf{B}^2U(1)$$

$$\uparrow$$

$$\begin{array}{ccccccc} \gamma & & \bullet & & \bullet & & \bullet \\ c \downarrow & \mapsto & c \downarrow & \mapsto & \gamma \downarrow & \mapsto & [\gamma] \downarrow \\ c \cdot \gamma & & \bullet & & \bullet & & \bullet \end{array}$$

$$\uparrow$$

$$\begin{array}{ccc} \bullet & & \bullet \\ \gamma \left(\begin{array}{c} \xrightarrow{c} \\ \downarrow \end{array} \right) c \cdot \gamma & \mapsto & \left(\begin{array}{c} \xrightarrow{c} \\ \downarrow \end{array} \right) \\ \bullet & & \bullet \end{array}$$

2-Groupoids with smooth structure – Dixmier-Douady class.

For example, write $U(n)$, $n \in \mathbb{N} \sqcup \{\omega\}$

for the unitary group on a countably-dimensional complex Hilbert space and regard this as a smooth group by its “continuous diffeology”:

$$U(n) : \mathbb{R}^k \mapsto \text{Probe}(\mathbb{R}^k, U(n)) := C^0(\mathbb{R}^k, U(n))$$

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Then we have the following long fiber sequence of smooth 2-groupoids:

$$\begin{array}{ccccccc}
 U(1) \hookrightarrow U(n) & \longrightarrow & \text{PU}(n) & & & & \\
 & & \uparrow \in \text{LWEq} & & & & \\
 & & U(n) // U(1) & \rightarrow & \mathbf{BU}(1) & \rightarrow & \mathbf{BU}(n) \longrightarrow \mathbf{BPU}(n) \\
 & & & & & & \uparrow \in \text{LWEq} \\
 & & & & & & \mathbf{BU}(n) // \mathbf{BU}(1) \rightarrow \mathbf{B}^2\mathbf{U}(1)
 \end{array}$$

2-Groupoids with smooth structure – Dixmier-Douady class.

For example, write $U(n)$, $n \in \mathbb{N} \sqcup \{\omega\}$ for the unitary group on a countably-dimensional complex Hilbert space and regard this as a smooth group by its “continuous diffeology”:

$$U(n) : \mathbb{R}^k \mapsto \text{Probe}(\mathbb{R}^k, U(n)) := C^0(\mathbb{R}^k, U(n))$$

Then we have the following long fiber sequence of smooth 2-groupoids:

$$\begin{array}{ccccccc}
 U(1) \hookrightarrow U(n) & \longrightarrow & \text{PU}(n) & & & & \\
 & & \uparrow \in \text{LWEq} & & & & \\
 & & U(n) // U(1) & \rightarrow & \mathbf{BU}(1) & \rightarrow & \mathbf{BU}(n) \longrightarrow \mathbf{BPU}(n) \\
 & & & & & & \uparrow \in \text{LWEq} \\
 & & & & & & \mathbf{BU}(n) // \mathbf{BU}(1) \rightarrow \mathbf{B}^2\mathbf{U}(1)
 \end{array}$$

This is compatible with complex conjugation, so we have a map of 2-stacks like this:

universal Dixmier-Douady class

$$\mathbf{BPU}(n) // \mathbb{Z}_2 \xleftarrow{\in \text{LWEq}} \xrightarrow{\text{DD}} \mathbf{B}^2\mathbf{U}(1) // \mathbb{Z}_2$$

2-Groupoids with smooth structure – Čech groupoids.

For X a smooth manifold with $\{U_i \hookrightarrow X\}_{i \in I}$ a **good open cover**, in that

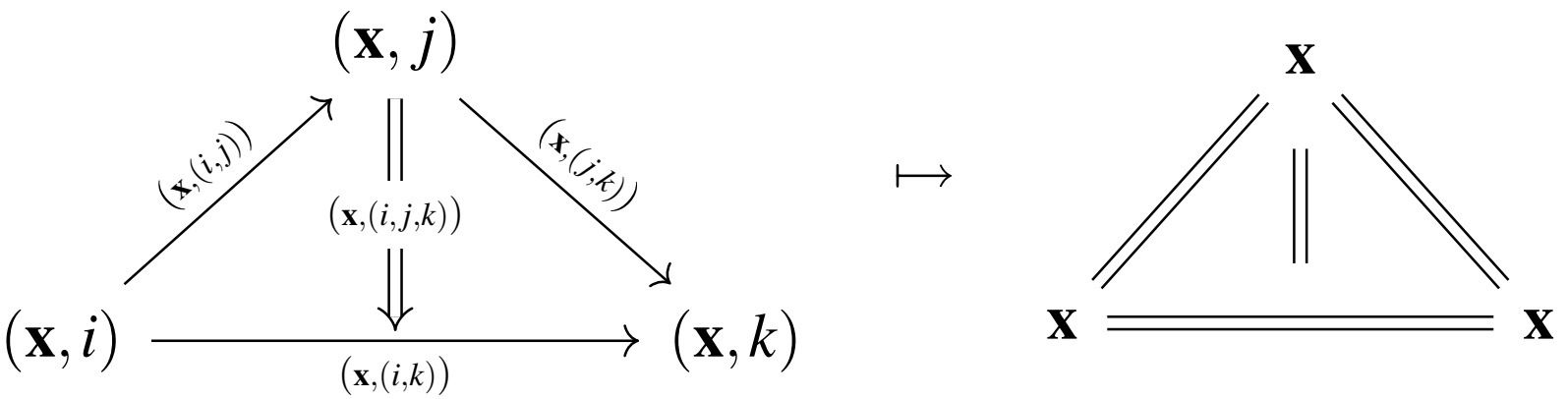
$$(\mathbf{x}, (i_1, \dots, i_n)) \in C^\infty(\mathbb{R}^m, U_{i_1} \cap \dots \cap U_{i_n}) \quad \Rightarrow \quad U_{i_1} \cap \dots \cap U_{i_n} \underset{\text{diff}}{\simeq} \mathbb{R}^{\dim(X)},$$

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we have the *smooth Čech 2-groupoid*:



which is a *projectively cofibrant resolution* of X .

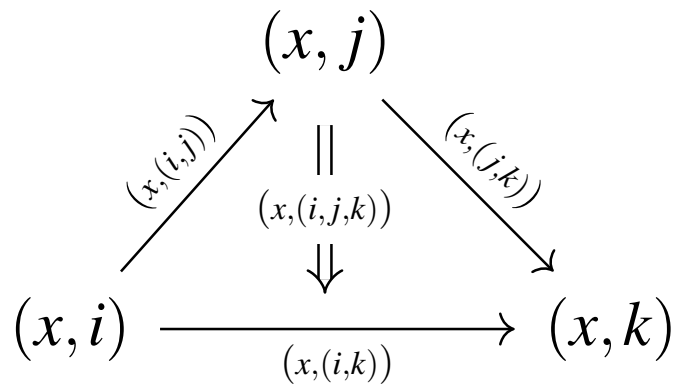
2-Groupoids with smooth structure – Čech cocycles.

Smooth 2-functors from such a Čech resolution $\widehat{X} \rightarrow X$

to the delooping $\mathbf{B}\Gamma$ of a Lie group

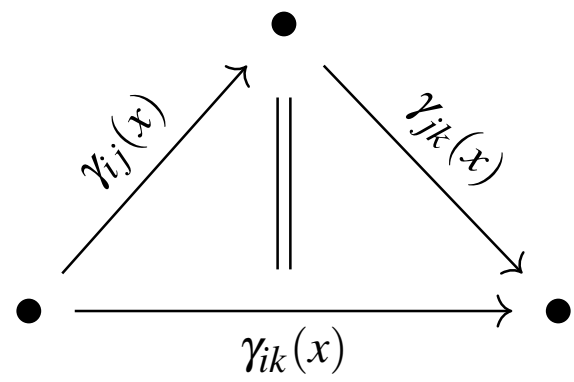
are *cocycles* in the *Čech cohomology* of X with coefficients in Γ :

$$\widehat{X} \xrightarrow{\text{smooth functor} = \check{\text{Cech cocycle}}} \mathbf{B}\Gamma$$



Čech relations

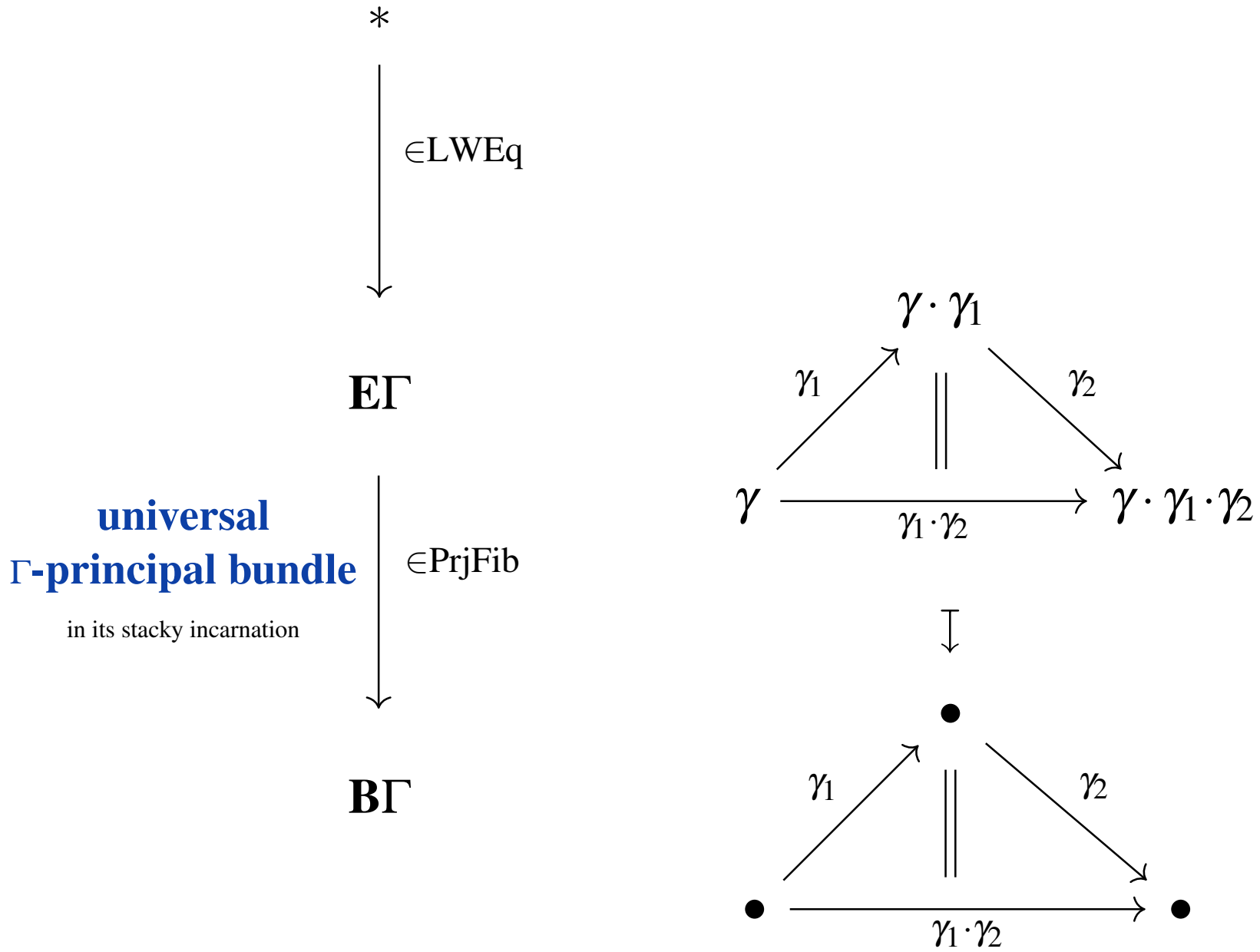
\mapsto



cocycle condition

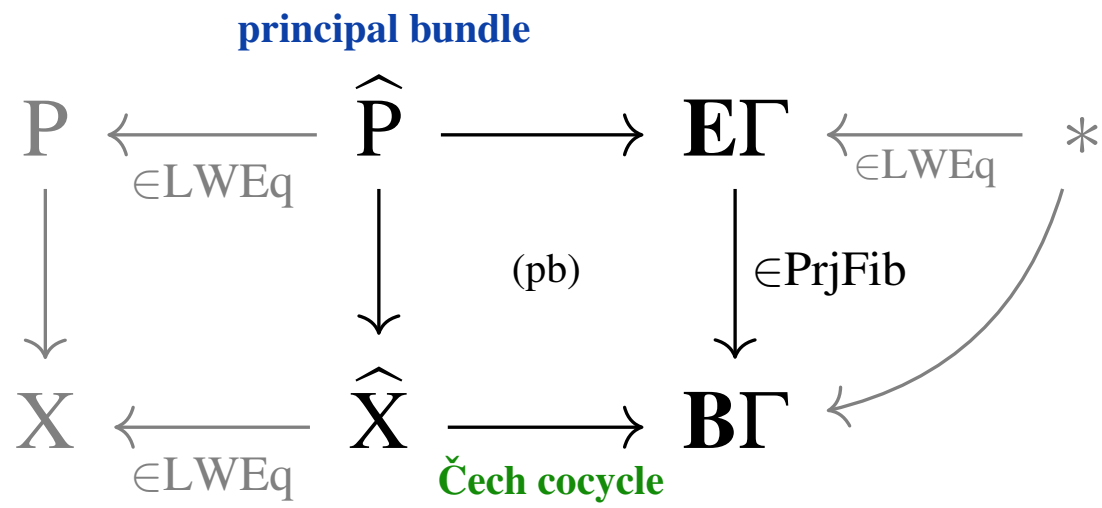
Principal bundles via smooth groupoids – Universal principal bundles.

The inclusion of the unique base point into $\mathbf{B}\Gamma$ has the following *fibrant resolution*:



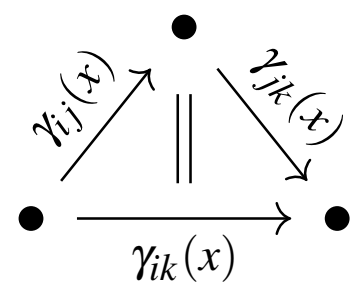
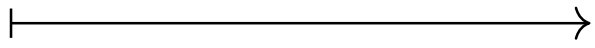
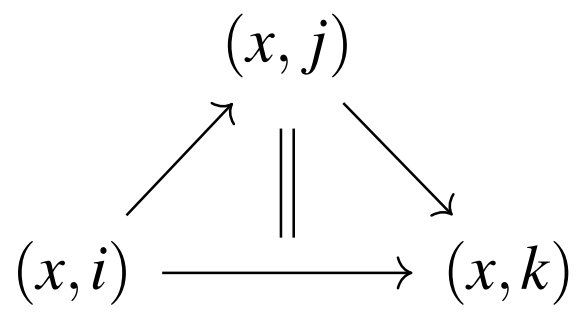
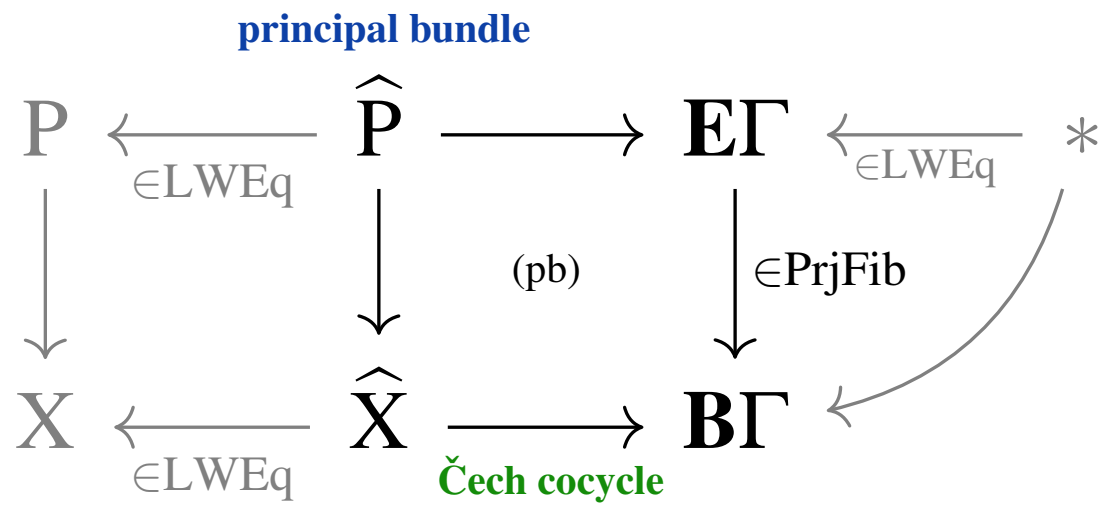
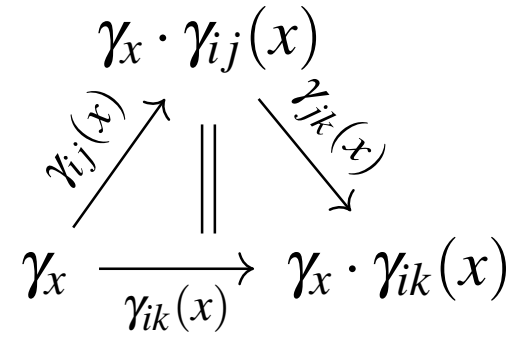
Principal bundles via smooth groupoids.

The *homotopy fiber* of a 2-functor = Čech cocycle is equivalently *the principal bundle P it classifies*:



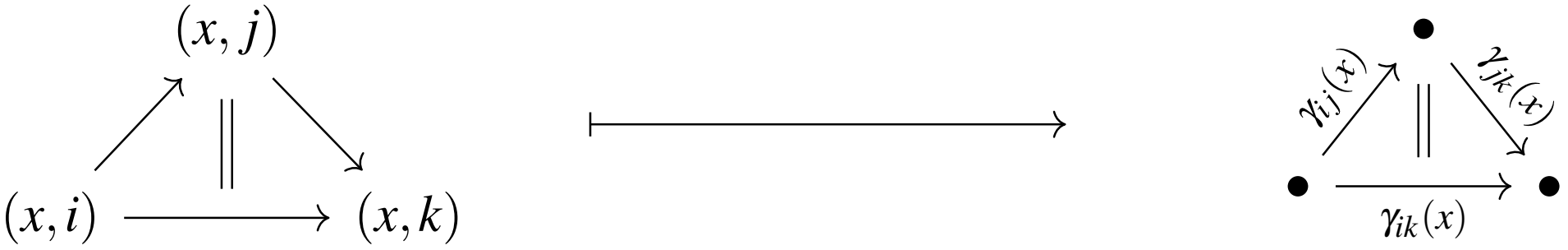
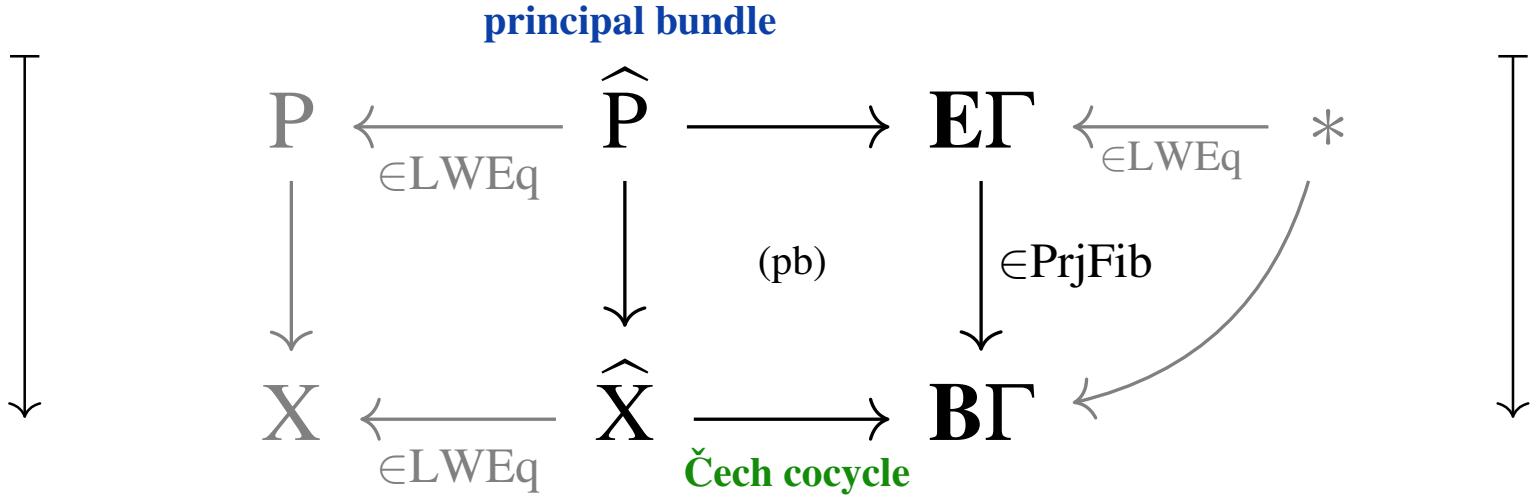
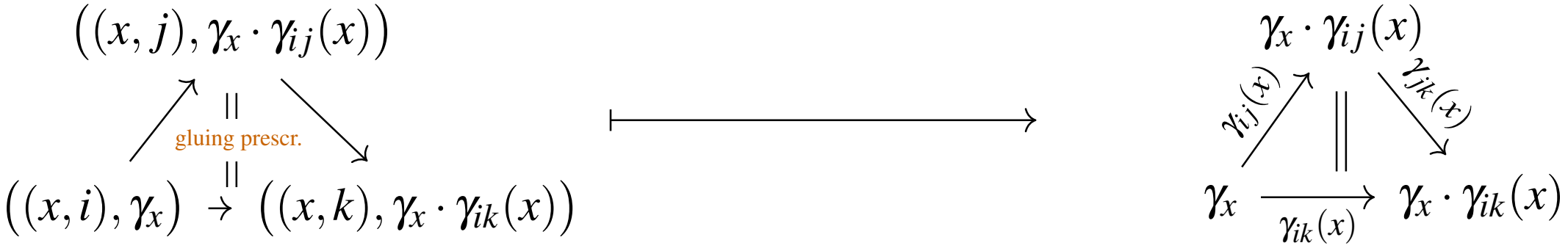
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Principal 2-bundles via smooth 2-groupoids.

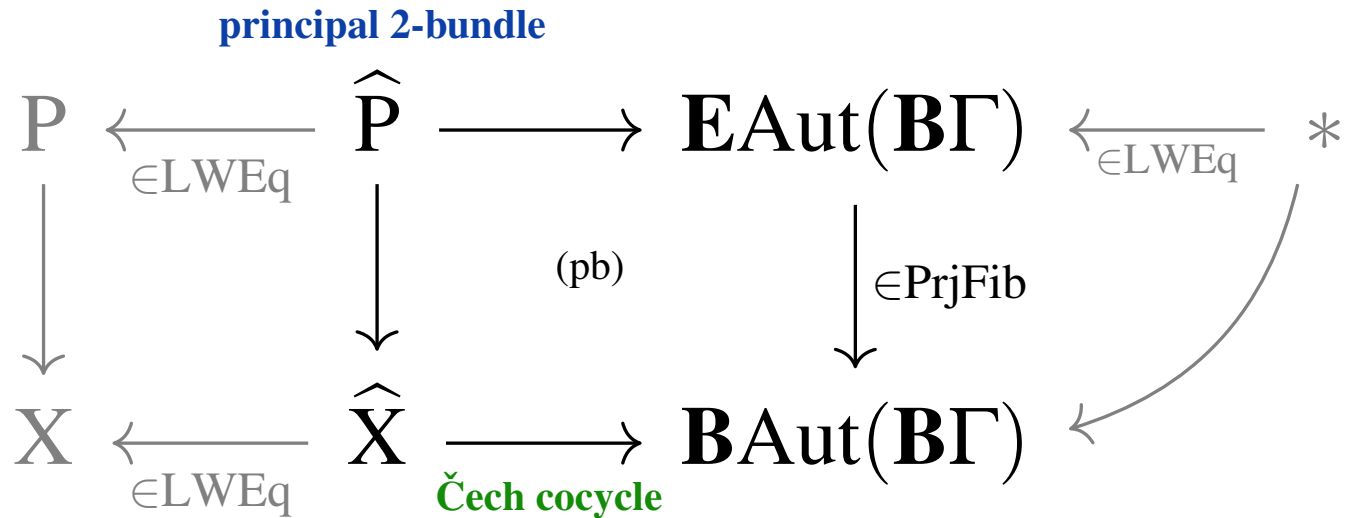
This neat formulation of ordinary principal bundles
immediatly generalizes to give [principal 2-bundles](#):

Principal 2-bundles via smooth 2-groupoids.

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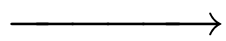
E.g. for the structure 2-group $\mathbf{Aut}(\mathbf{B}\Gamma)$

these are equivalently Giraud's *non-abelian gerbes*:



While it's tradition to be esoteric about this simple affair,

here to highlight that this is really about *twisted cohomology*:

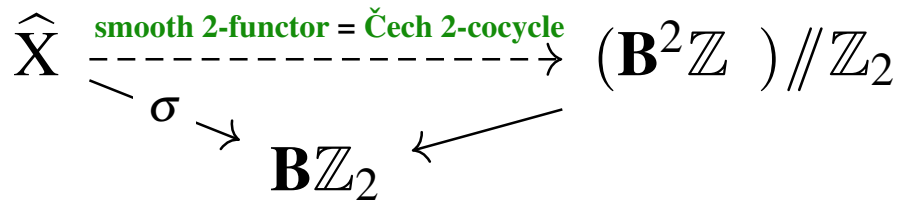


Principal 2-bundles via smooth 2-groupoids – Example: Twisted cohomology.

For structure 2-group $\text{Aut}(\mathbf{B}\mathbb{Z}) \simeq (\mathbf{B}\mathbb{Z}) \rtimes \mathbb{Z}_2$,

with $\mathbf{BAut}(\mathbf{B}\mathbb{Z}) \simeq (\mathbf{B}^2\mathbb{Z}) // \mathbb{Z}_2$ and $\widehat{X} \xrightarrow{\sigma} \mathbf{B}\mathbb{Z}_2$ a double covering, then

2nd integral cohomology of X with local coefficients is smooth 2-functors:



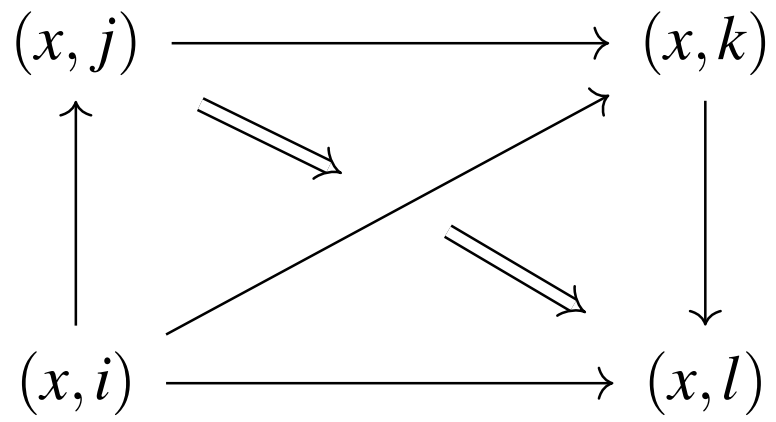
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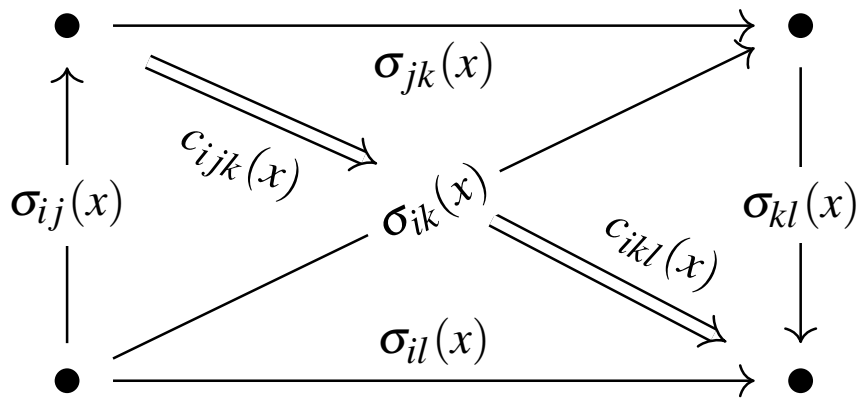
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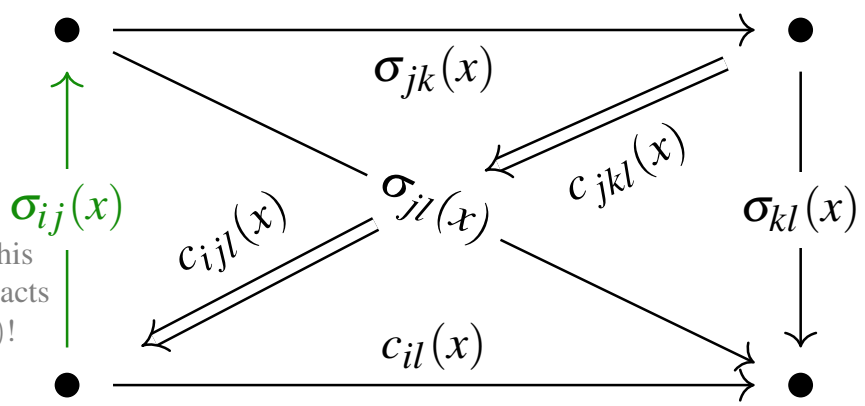
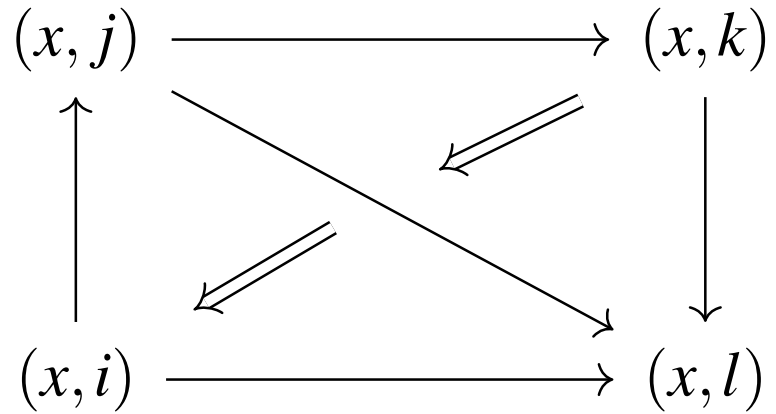
$$\widehat{X} \begin{array}{c} \xrightarrow{\text{smooth 2-functor} = \check{\text{Cech 2-cocycle}} \\ \searrow \sigma \\ \mathbf{B}\mathbb{Z}_2 \end{array} (\mathbf{B}^2\mathbb{Z}) // \mathbb{Z}_2$$



||| Čech relations



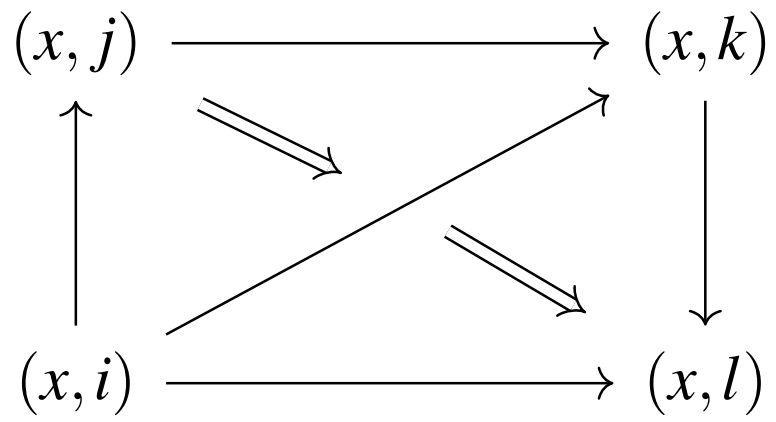
||| cocycle condition



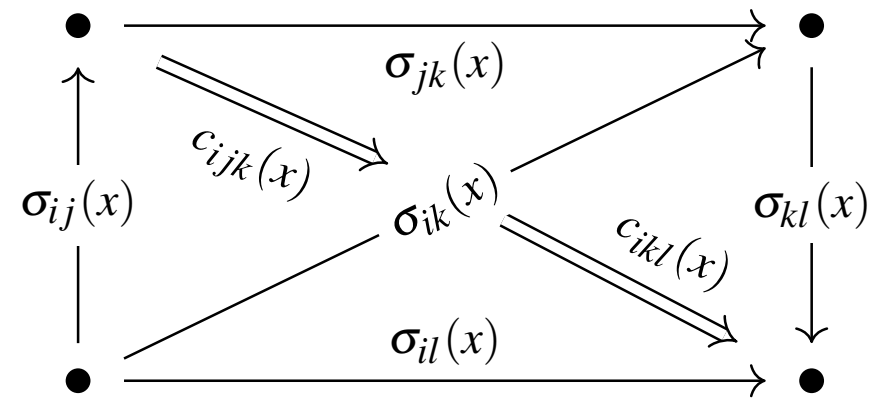
recall that this 1-morphism acts on $c_{jkl}(x)$!

Principal 2-bundles via smooth 2-groupoids – Example: Jandl bundle gerbes.

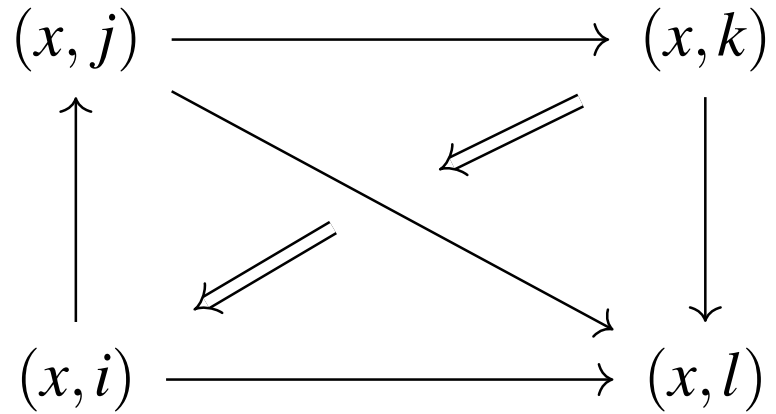
For structure 2-group $\text{Aut}(\mathbf{BU}(1)) \simeq (\mathbf{BU}(1)) \rtimes \mathbb{Z}_2$,
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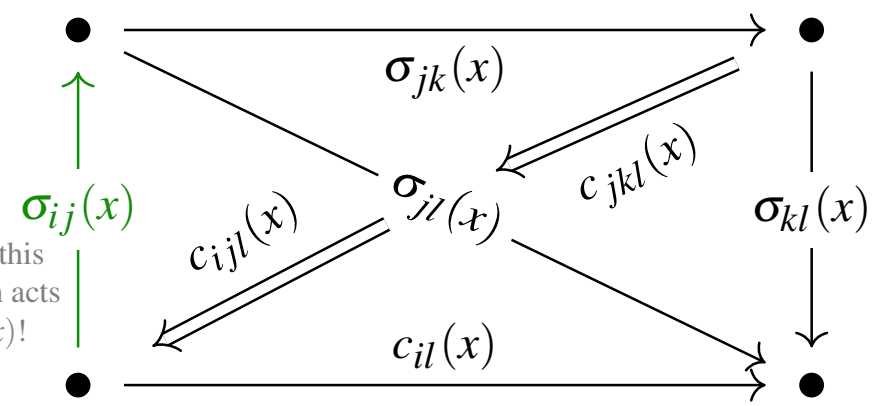
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||| cocycle condition |||

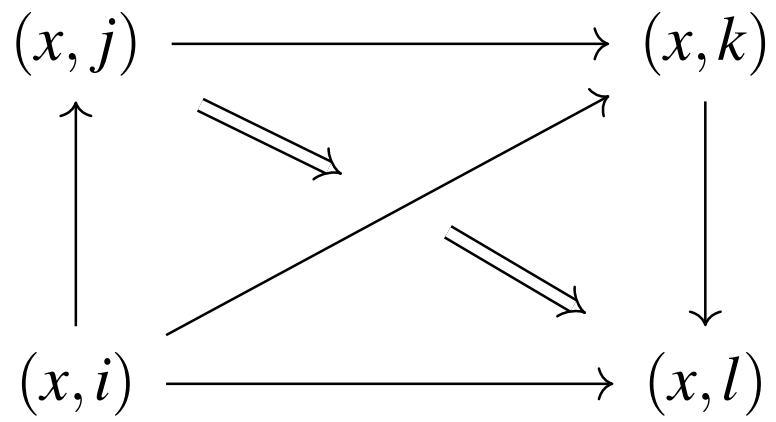
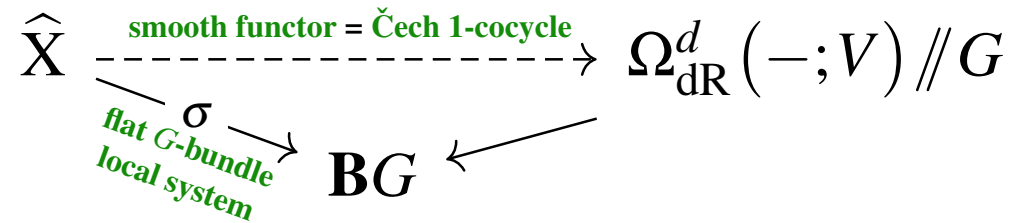


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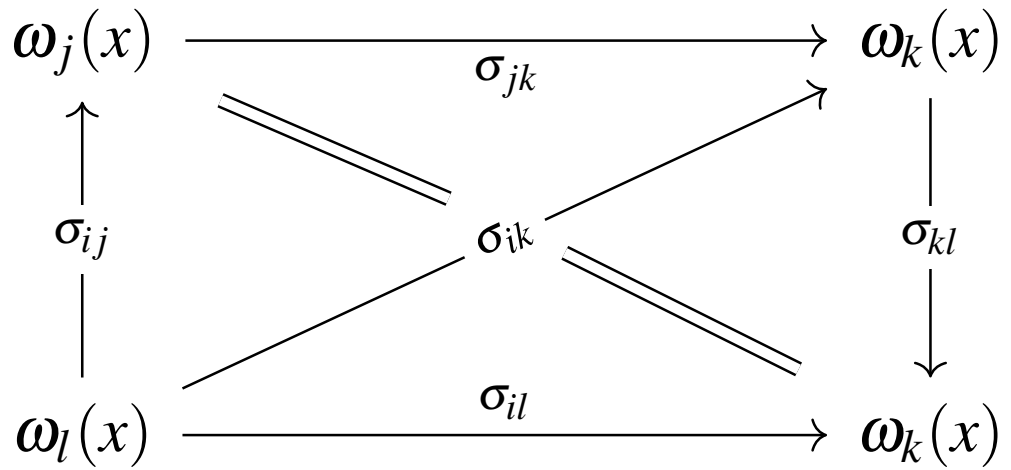
Principal 2-bundles via smooth 2-groupoids – Example: Twisted cohomology.

For equivariant de Rham coefficients with $G \curvearrowright V$ a representation of a finite group:

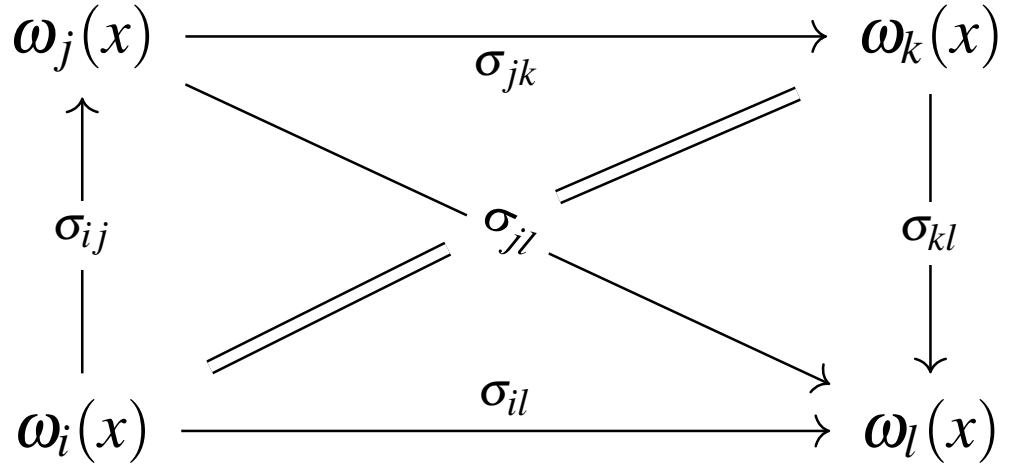
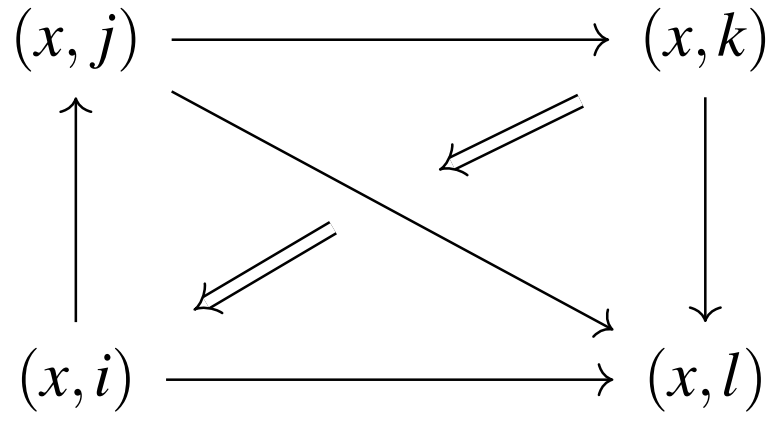


||| Čech relations

→



||| cocycle condition |||



Principal 2-bundles via smooth 2-groupoids – Punchline.

So:

Non-abelian 1-cohomology is modulated by 1-stacks $\mathbf{B}\Gamma$,
abelian 2-cohomology is modulated by 2-stacks \mathbf{B}^2A , etc.

Higher fiber/principal bundles are *bundles of such moduli stacks*,
hence are families of moduli stacks that vary over the base space,
hence locally modulate cohomology as before,
but now subject to global twists.

Principal 2-bundles via smooth 2-groupoids – Equivariance.

Finally, the “higher topos” of smooth 2-groupoids

has *equivariance* natively built into it: just let domain spaces be groupoids, too.

Principal 2-bundles via smooth 2-groupoids – Equivariance.

Finally, the “higher topos” of smooth 2-groupoids has *equivariance* natively built into it: just let domain spaces be groupoids, too:

For $X \curvearrowright G$ a smooth action of a finite group on a smooth manifold.

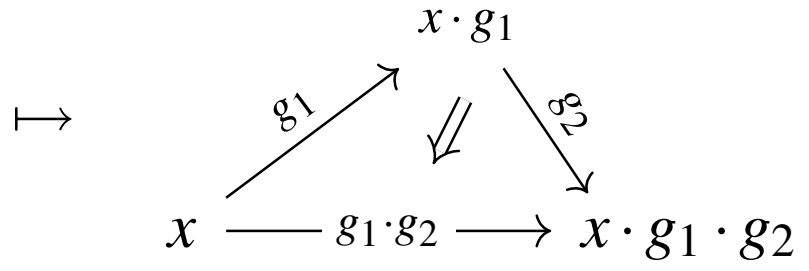
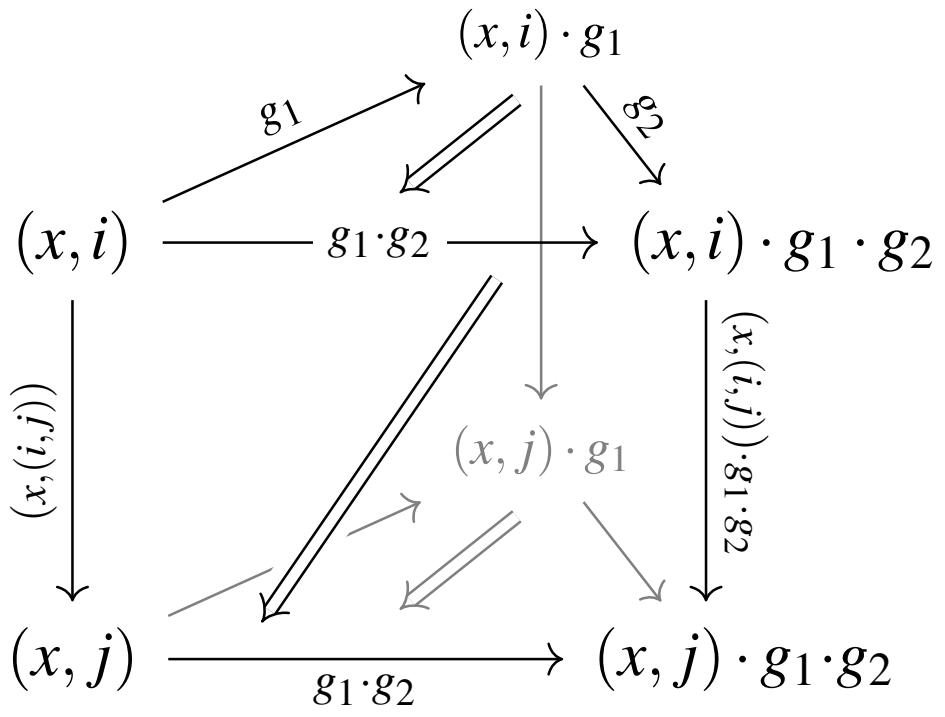
there exists an *equivariant good open cover*

$$\coprod_{i \in I} \overset{G}{U}_i \longrightarrow X$$

and its *equivariant Čech groupoid*:

$$\emptyset \xrightarrow{\in \text{PrjCof}} \widehat{X // G} \xrightarrow[\in \text{LWEq}]{\text{cofibrant resolution}} X // G$$

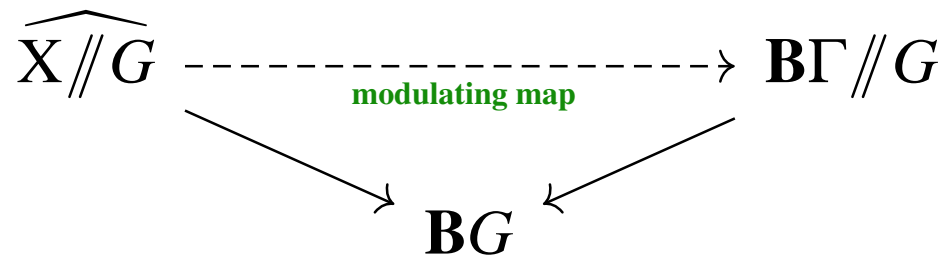
action groupoid



Principal 2-bundles via smooth 2-groupoids – Equivariance.

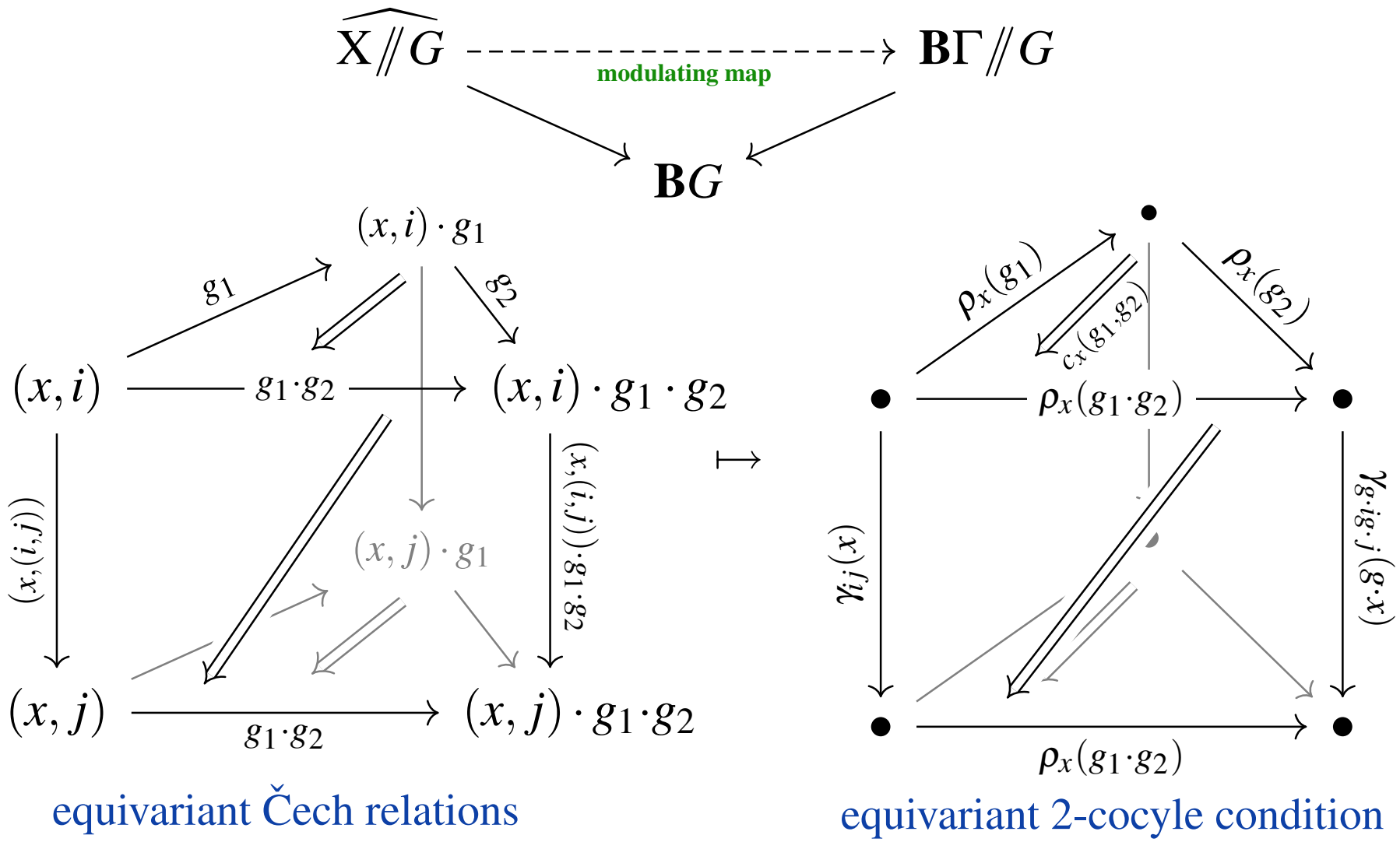
For $X \curvearrowright G$ a smooth manifold and $\Gamma \curvearrowright G$ a smooth 2-group both equipped with smooth G -action, a

G -equivariant Γ -principal 2-bundle on X is modulated by a smooth 2-functor like this:



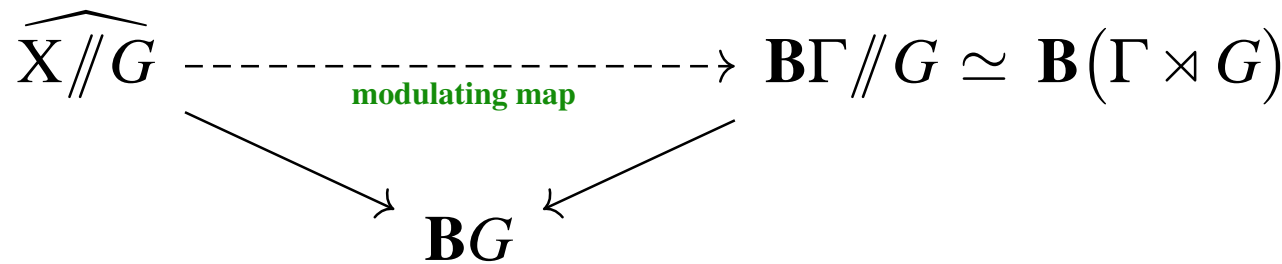
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on the right we have equivalently the **semidirect product 2-group**.

Principal 2-bundles via smooth 2-groupoids – Equivariance.

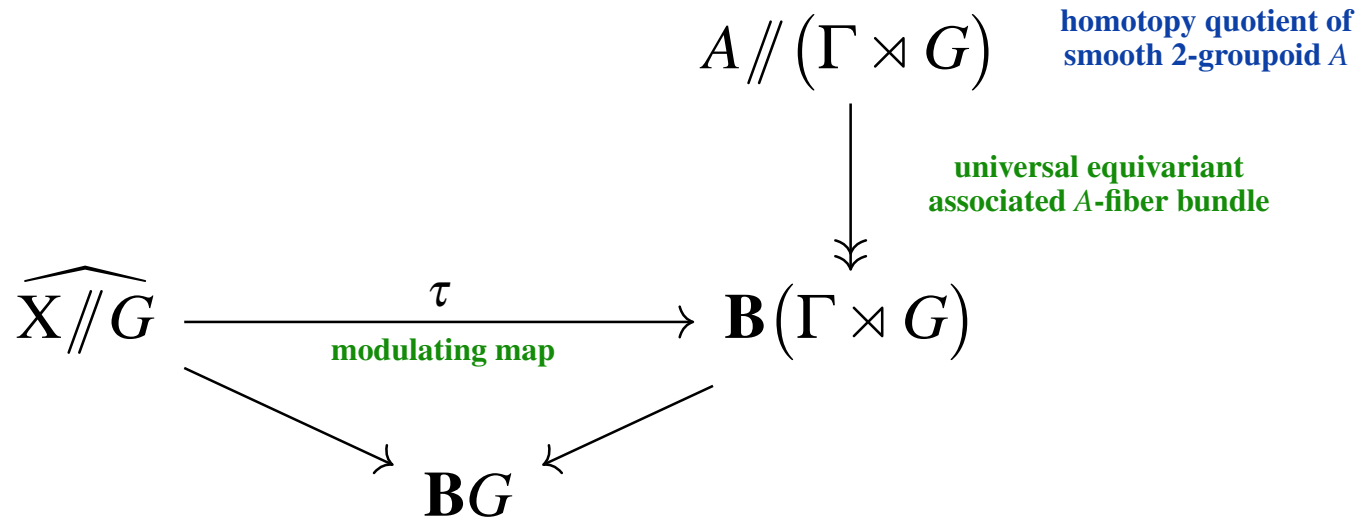
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G -equivariant Γ -principal 2-bundle on X is modulated by a smooth 2-functor like this:

$$\begin{array}{ccc}
 \widehat{X // G} & \xrightarrow[\text{modulating map}]{\text{eg. equivariant real bundle gerbe}} & \mathbf{B}^2\mathbf{U}(1) // \mathbb{Z}_2 \simeq \mathbf{B}(\mathbf{BU}(1) \rtimes \mathbb{Z}_2) \\
 & \searrow & \swarrow \\
 & \mathbf{B}\mathbb{Z}_2 &
 \end{array}$$

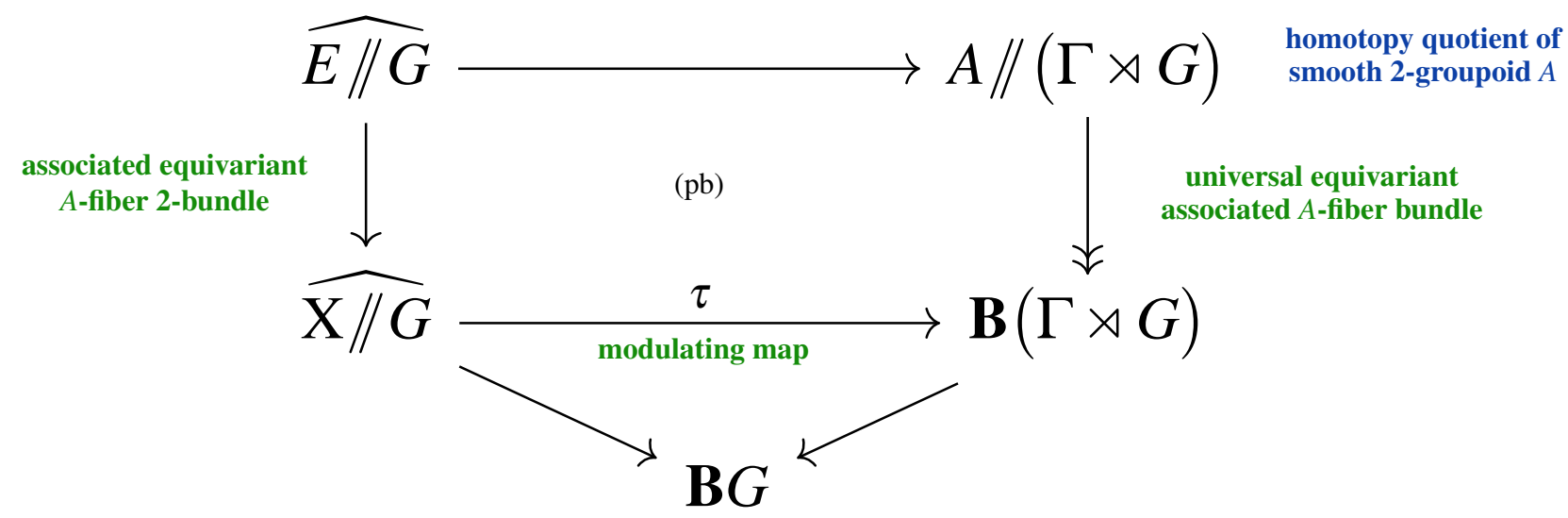
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Principal 2-bundles via smooth 2-groupoids – Associated 2-bundles.



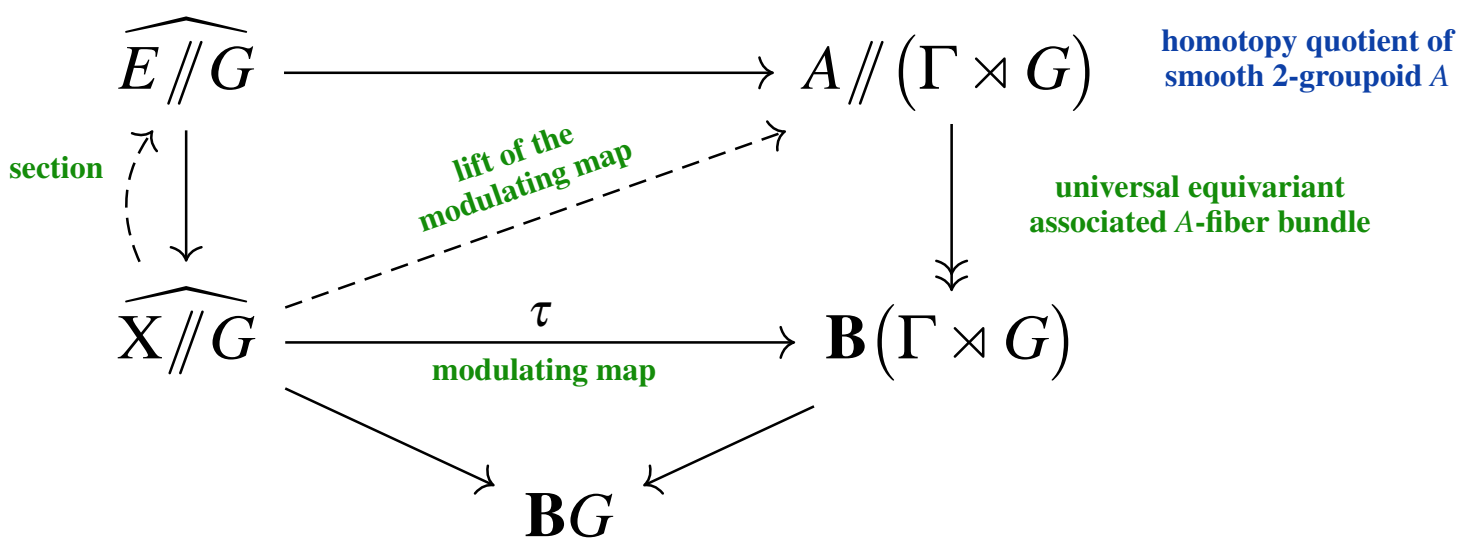
A fibration over that is equivalently an equivariant ∞ -action $(\Gamma \rtimes G) \curvearrowright A$ embodied by its universal *associated 2-bundle*.

Principal 2-bundles via smooth 2-groupoids – Associated 2-bundles.



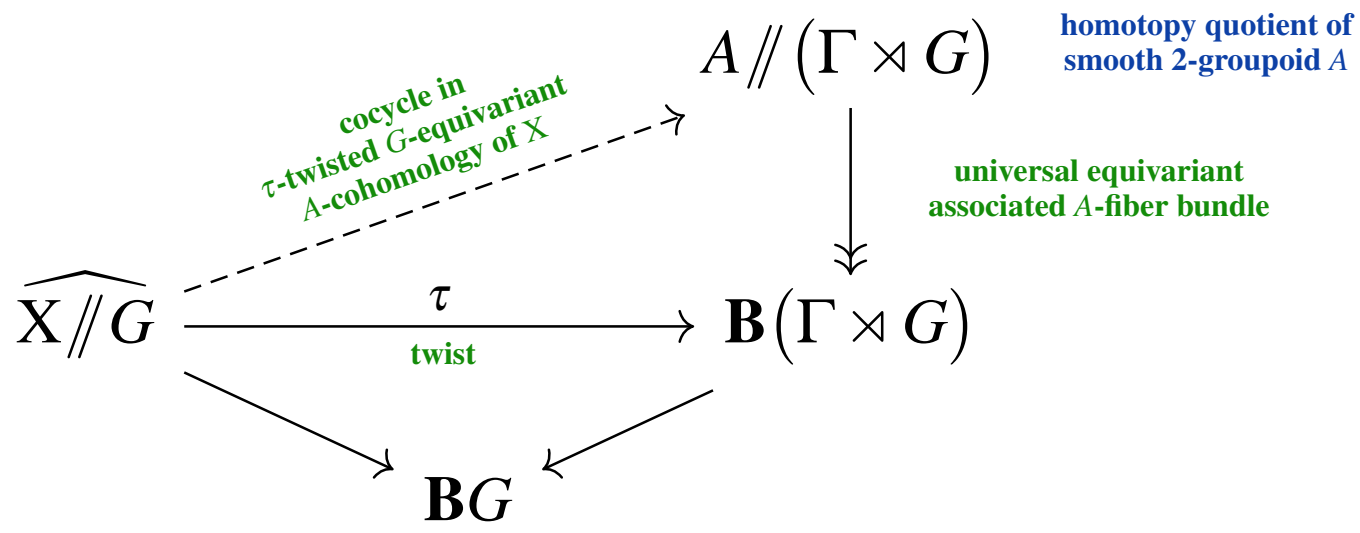
Its pullback is
the **equivariant A-fiber 2-bundle**
which is associated to
the given equivariant principal 2-bundle.

Principal 2-bundles via smooth 2-groupoids – Associated 2-bundles.



The equivariant sections are equivalently the lifts of the modulating map.

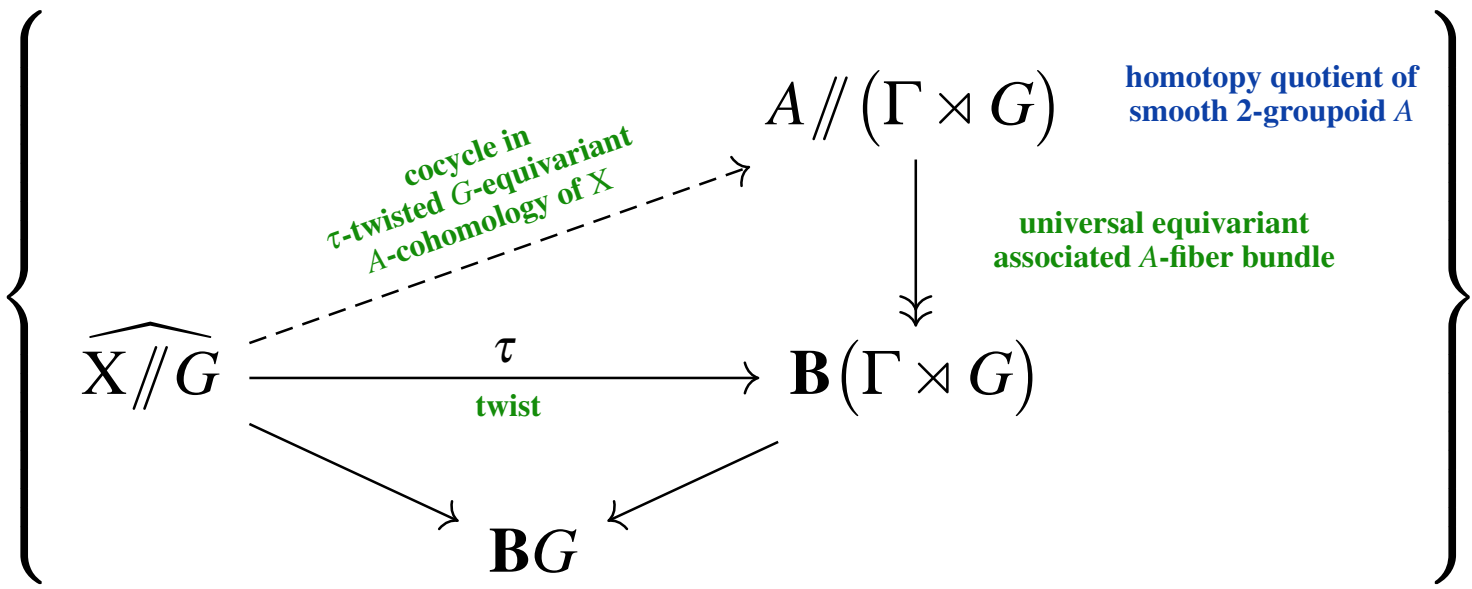
Twisted equivariant non-abelian cohomology.



Equivalently, these are the *cocycles* of τ -twisted G -equivariant A -cohomology.

Twisted equivariant non-abelian cohomology.

twisted equivariant
cohomology set
 $A_G^\tau(X) :=$



homotopy quotient of
smooth 2-groupoid A

universal equivariant
associated A -fiber bundle

\sim_{htpy}

Equivalently, these are the *cocycles* of τ -twisted G -equivariant A -cohomology.

0 – Cohesive ∞ -Topoi

I – Equivariant ∞ -Bundles

II – TED-K-Theory

III – Anyonic Defect Branes

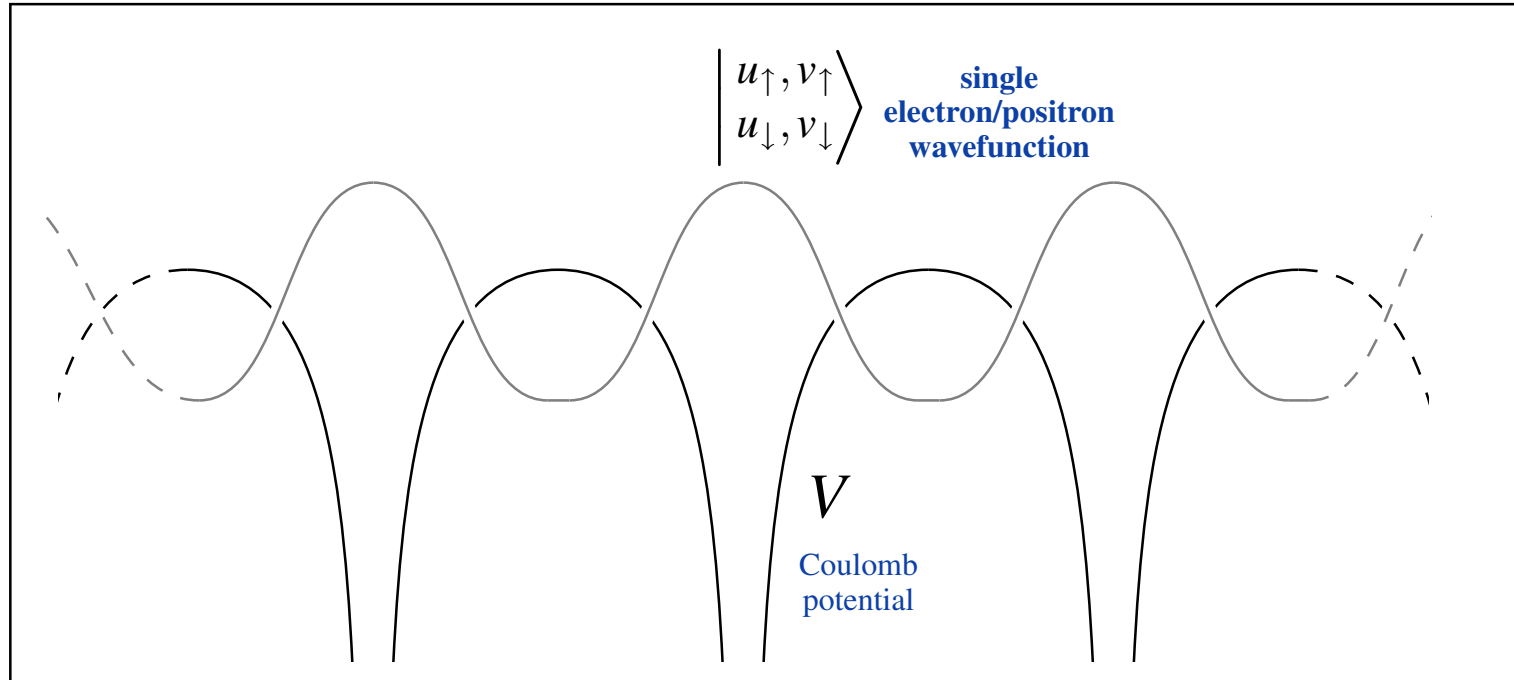
IV – Quantum Computation

This part is a
quick motivation and exposition of
twisted equivariant KR-theory
following these articles:

<i>Equivariant Principal ∞-bundles</i>	[arXiv:2112.13654]
<i>Anyonic Defect Branes in TED-K-Theory</i>	[arXiv:2203.11838]
<i>The TED character map</i>	(in preparation)

Vacua of electron/positron field in Coulomb background.

Fact ([KS77][CHO82]). The **vacua** of the free Dirac quantum field in a classical Coulomb background...



Vacua of electron/positron field in Coulomb background.

Fact ([KS77][CHO82]). The vacua of the free Dirac quantum field in a classical Coulomb background are **characterized by Fredholm operators...**

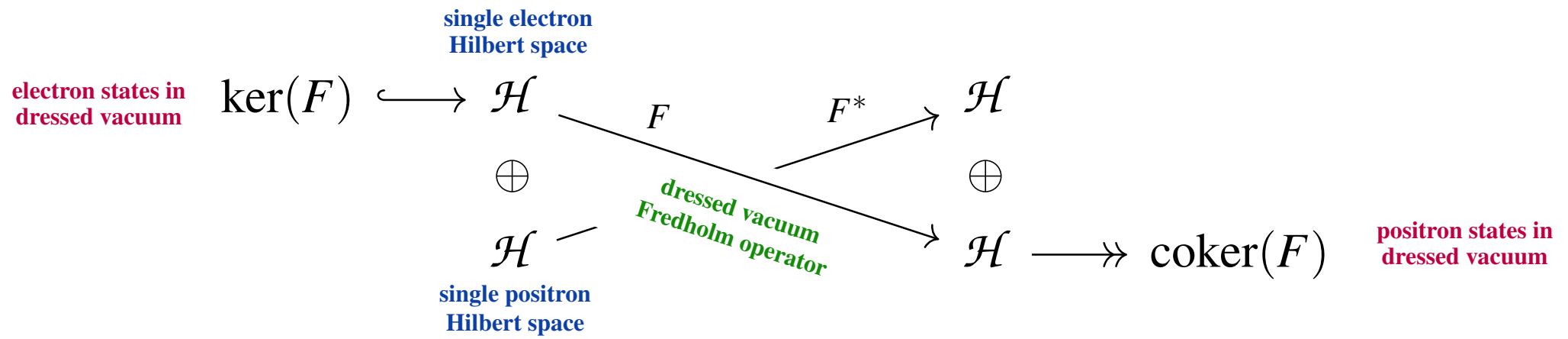
$$\begin{array}{c} \text{finite-dimensional kernel} \\ \ker(F) \hookrightarrow \mathcal{H} \xrightarrow[\text{bounded linear}]{F} \mathcal{H} \twoheadrightarrow \text{coker}(F) \\ \text{finite-dimensional cokernel} \\ \underbrace{\psi \in \mathcal{H} \mid \forall \phi \langle \phi | F | \psi \rangle = 0} \qquad \underbrace{\psi \in \mathcal{H} \mid \forall \phi \langle \psi | F | \phi \rangle = 0} \end{array}$$

Vacua of electron/positron field in Coulomb background.

Fact ([KS77][CHO82]). The vacua of the free Dirac field in a classical Coulomb background are characterized by **Fredholm operators**

$$\begin{array}{c}
 \text{finite-dimensional kernel} \qquad \qquad \text{Fredholm operator} \qquad \qquad \text{finite-dimensional cokernel} \\
 \ker(F) \hookrightarrow \mathcal{H} \xrightarrow[\text{bounded linear}]{F} \mathcal{H} \twoheadrightarrow \text{coker}(F) \\
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 \end{array}$$

on the single-electron/positron Hilbert space:



$$\begin{array}{l}
 \text{total charge in dressed vacuum} \\
 \text{ind}(F) = \dim(\ker(F)) - \dim(\text{coker}(F)) \\
 = \dim(\text{coker}(F^*)) - \dim(\ker(F^*))
 \end{array}$$

Quantum symmetries.

On these dressed vacua of electron/positron states
the following *CPT-twisted projective group*

$$\frac{\text{even projective unitary group}}{\text{U}(\mathcal{H}) \times \text{U}(\mathcal{H})} \times \left(\underbrace{\mathbb{Z}_2}_{\{e,P\}} \times \underbrace{\mathbb{Z}_2}_{\{e,T\}} \right)$$

group of quantum symmetries

$$C := PT, \quad P \cdot [U_+, U_-] := [U_-, U_+] \cdot P, \quad T \cdot [U_+, U_-] := [\bar{U}_+, \bar{U}_-] \cdot T$$

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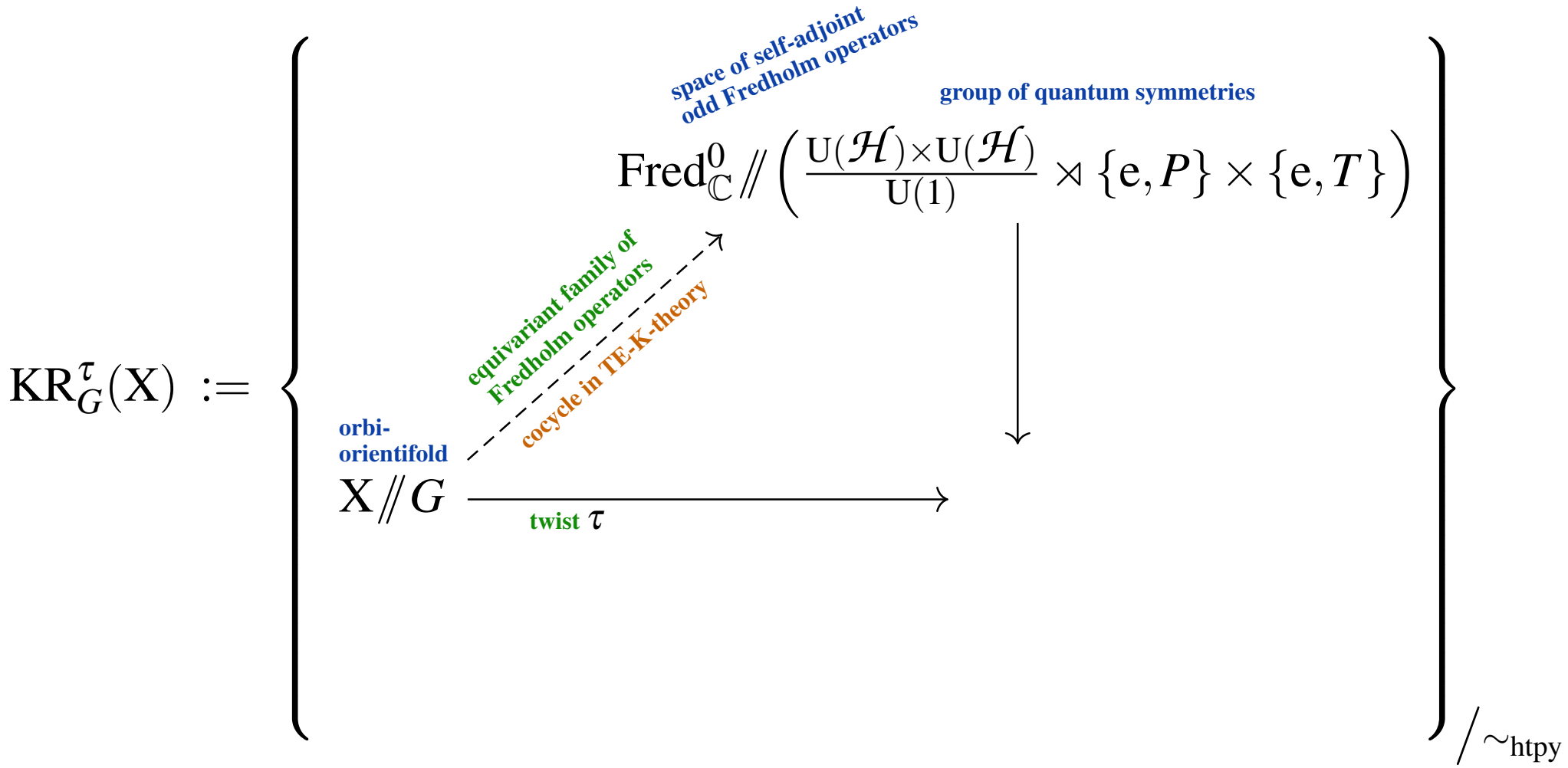
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naturally acts by conjugation:

$$\begin{aligned} [U_+, U_-] &: F \longmapsto U_+^{-1} \circ F \circ U_- \\ C \cdot [U_+, U_-] &: F \longmapsto U_-^{-1} \circ F^t \circ U_+ \\ P \cdot [U_+, U_-] &: F \longmapsto U_-^{-1} \circ F^* \circ U_+ \\ T \cdot [U_+, U_-] &: F \longmapsto U_+^{-1} \circ \bar{F} \circ U_- \end{aligned}$$

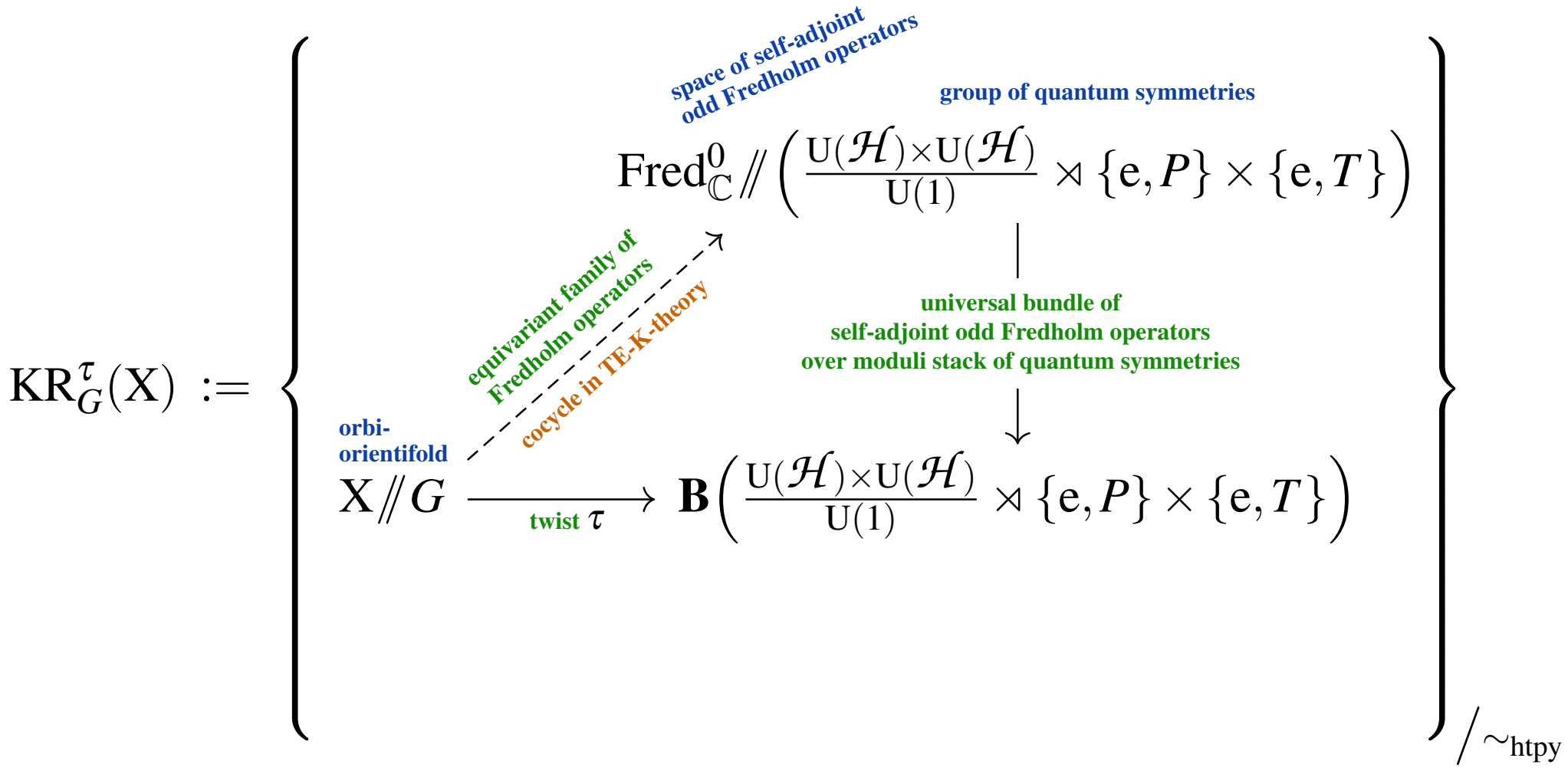
Twisted equivariant KR-theory – As a single diagram of smooth groupoids.

Homotopy classes of quantum-symmetry equivariant families of such self-adjoint odd Fredholm operators constitute *twisted equivariant KR-cohomology*:



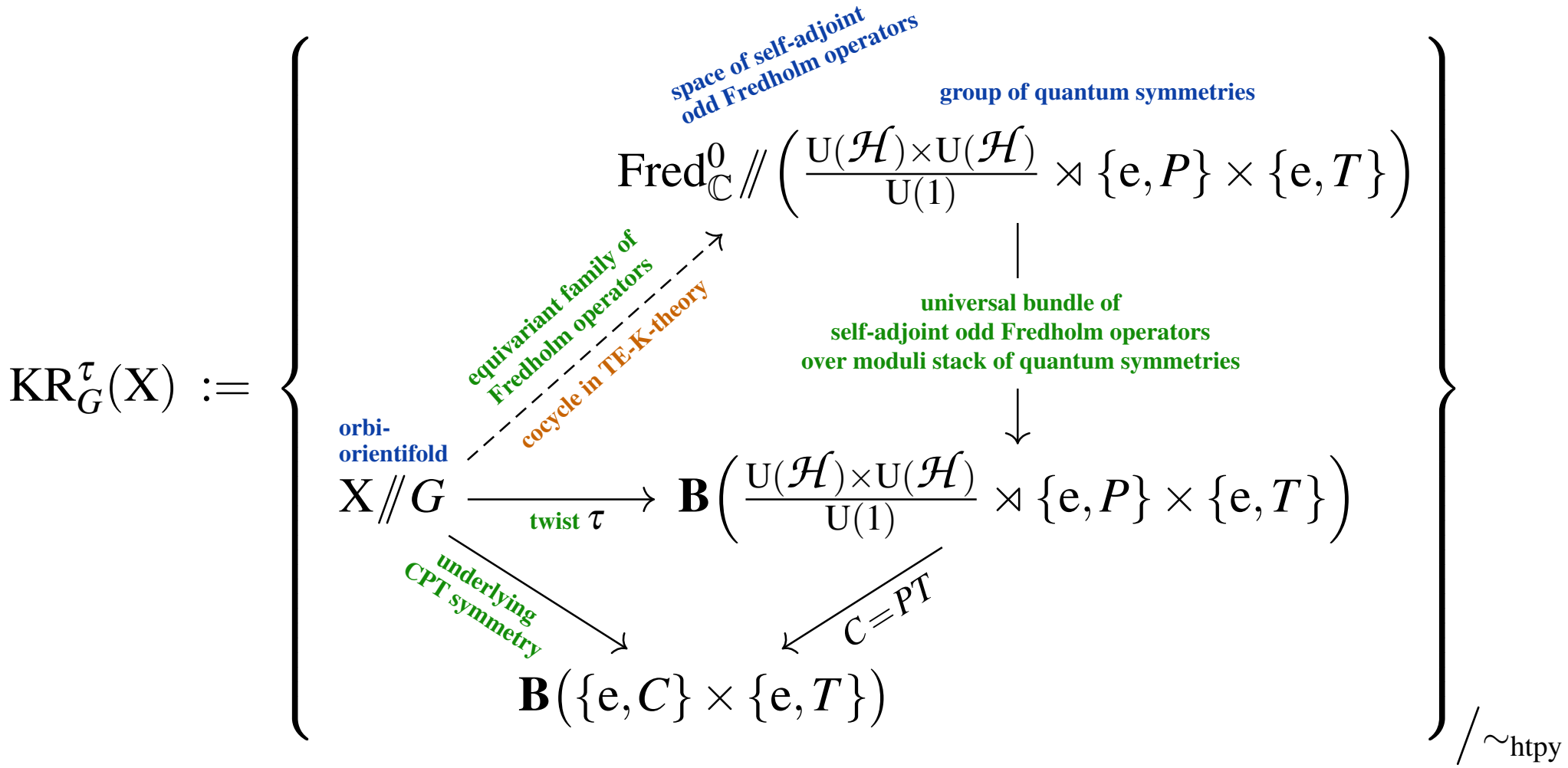
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CPT Quantum symmetries.

$$\begin{array}{ccc}
 \mathbf{B}(\{e, T\}) & \xrightarrow[\text{pure quantum T-symmetry}]{T \mapsto \hat{T}} & \mathbf{B}\left(\frac{\mathbf{U}(\mathcal{H}) \times \mathbf{U}(\mathcal{H})}{\mathbf{U}(1)} \rtimes \{e, T\}\right) \longrightarrow \mathbf{B}(\mathbf{BU}(1) \rtimes \{e, T\}) \\
 \searrow & & \swarrow \\
 & & \mathbf{B}(\{e, P\} \times \{e, T\})
 \end{array}$$

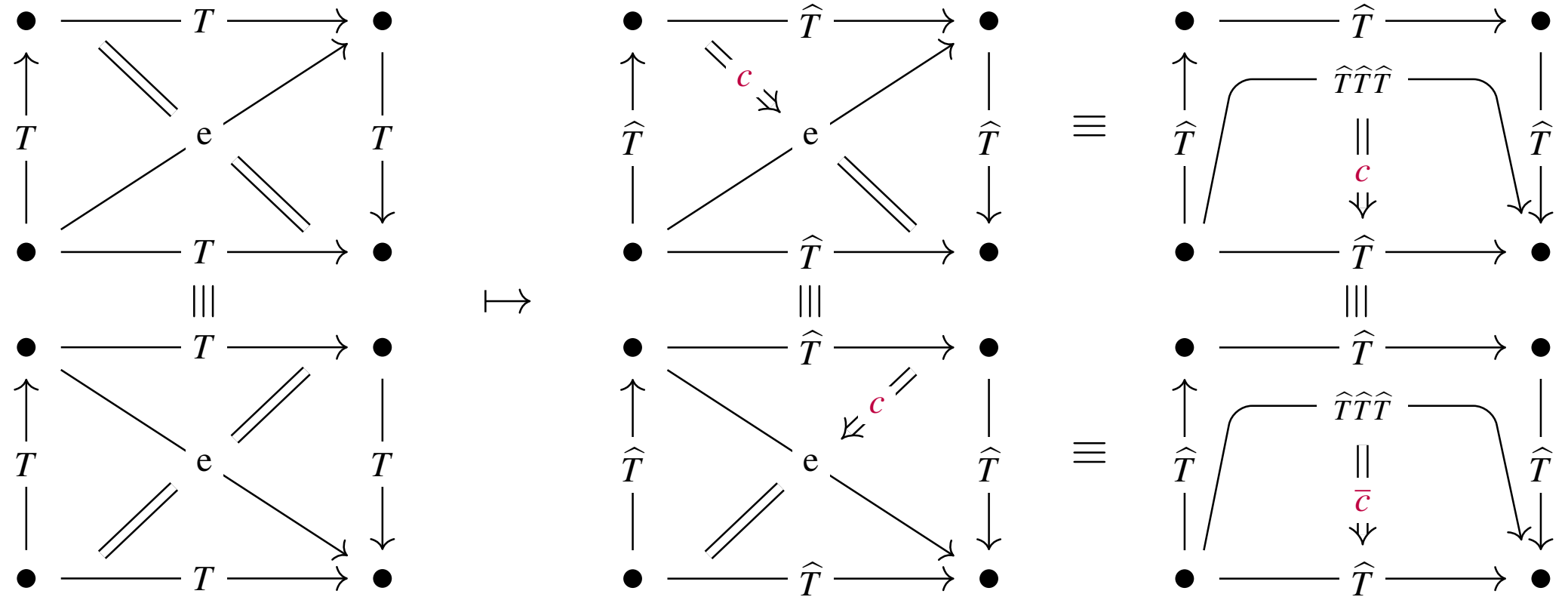
Let's use the previous machinery to compute the possible quantum T-symmetries...

CPT Quantum symmetries.

pure quantum T-symmetry

$$\mathbf{B}(\{e, T\}) \xrightarrow{T \mapsto \hat{T}} \mathbf{B}\left(\frac{\mathbf{U}(\mathcal{H}) \times \mathbf{U}(\mathcal{H})}{\mathbf{U}(1)} \rtimes \{e, T\}\right) \longrightarrow \mathbf{B}(\mathbf{BU}(1) \rtimes \{e, T\})$$

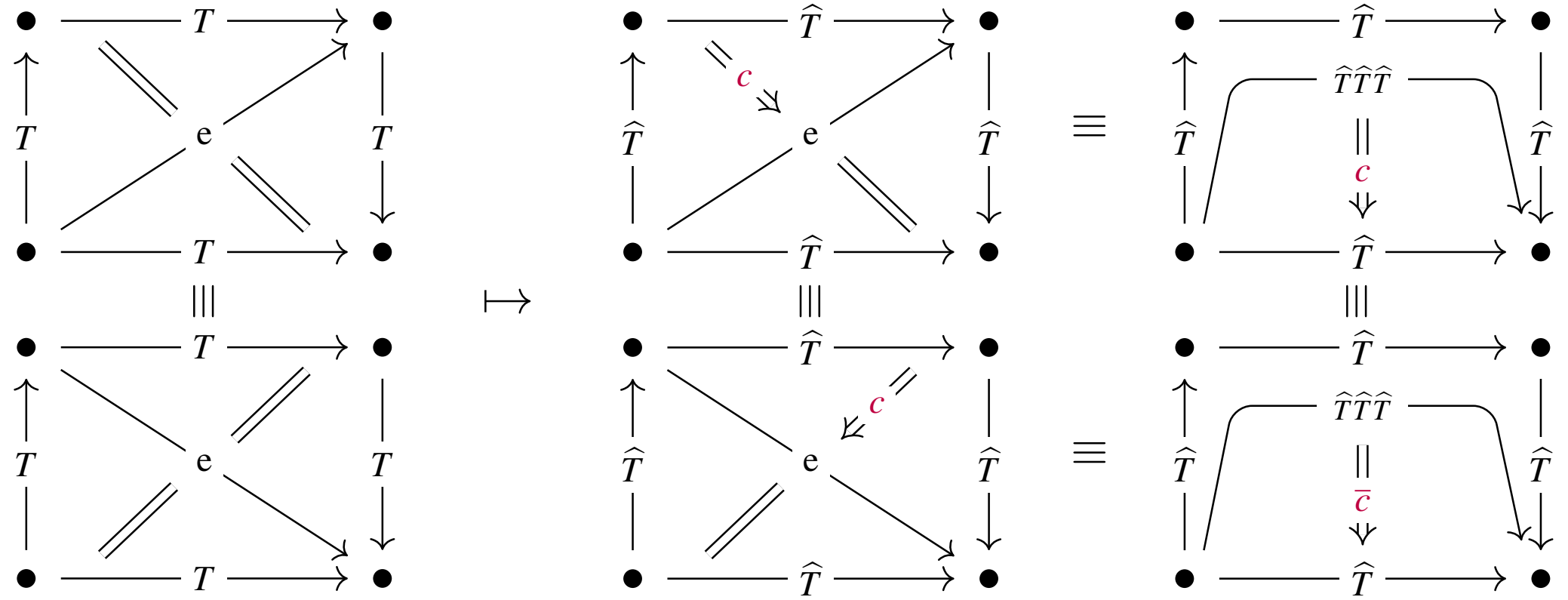
\swarrow \swarrow
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$\swarrow \qquad \searrow$
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So $\bar{c} = c$ and hence there are **two choices for quantum T-symmetry**, up to homotopy:

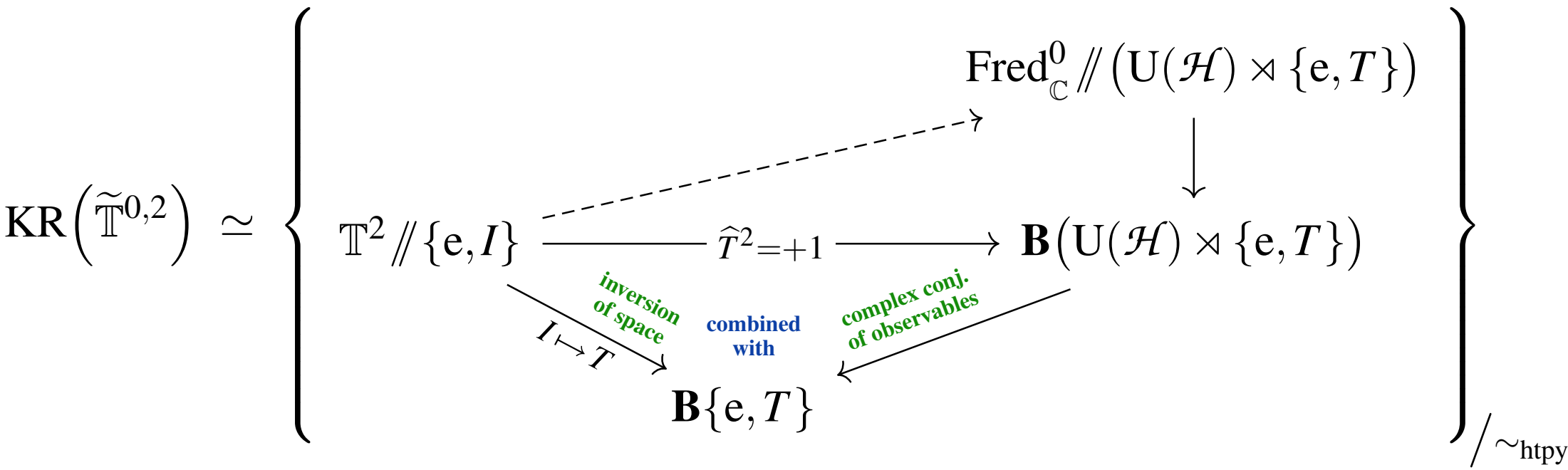
$$\hat{T}^2 = \pm 1 \quad \text{and similarly} \quad \hat{C}^2 = \pm 1.$$

Example – Orientifold KR-theory

Let I be *Inversion* action on 2-torus $\tilde{\mathbb{T}}^2 \simeq \mathbb{R}^2 / \mathbb{Z}^2$ and trivial action on observables

$$\begin{array}{ccc}
 \mathbb{T}^2 & \xrightarrow{I} & \mathbb{T}^2 \\
 k & \longmapsto & -k, \\
 \text{Fred}_{\mathbb{C}}^0 & \xrightarrow{I} & \text{Fred}_{\mathbb{C}}^0 \\
 F & \longmapsto & F.
 \end{array}$$

If T acts as I on \mathbb{T}^2 , then $\text{KR}^{\hat{T}^2 = +1}$ is *Atiyah's Real K-theory* aka *orienti-fold* K-theory:

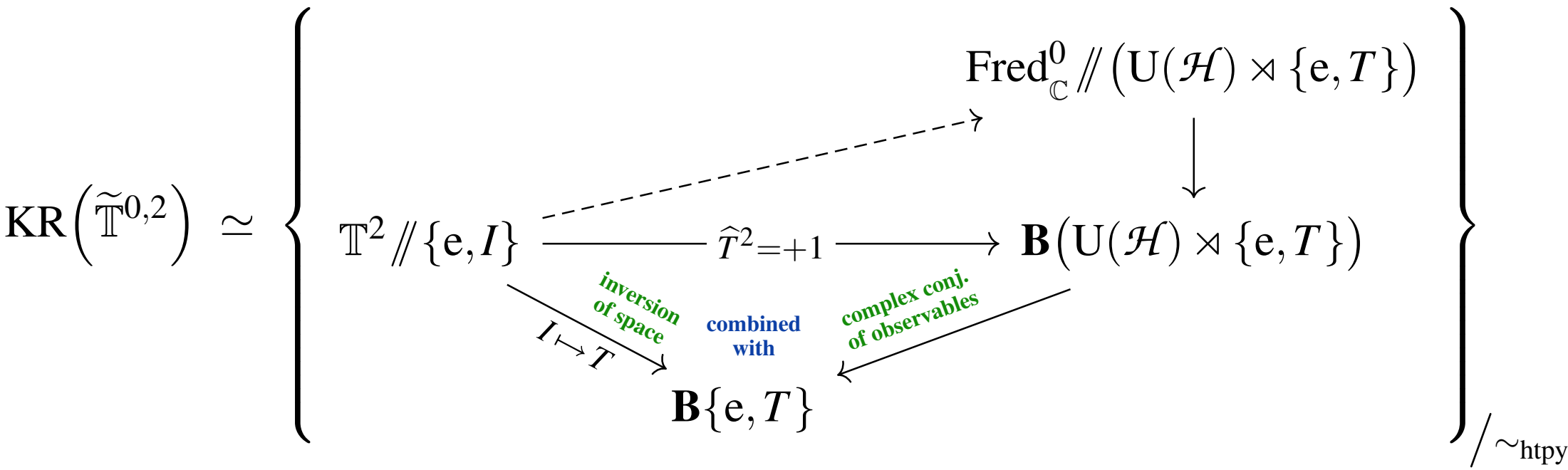


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But what happens on I -fixed loci i.e. on “orientifolds” ? →

CPT Quantum symmetries – 10 global choices.

(following [FM12, Prop. 6.4])

Equivariance group	$G =$	$\{e\}$	$\{e, P\}$	$\{e, T\}$		$\{e, C\}$		$\{e, T\} \times \{e, C\}$			
Realization as quantum symmetry	$\tau:$ $\hat{T}^2 =$			+1	-1			+1	-1	-1	+1
	$\hat{C}^2 =$					+1	-1	+1	+1	-1	-1
Maximal induced Clifford action anticommuting with all G -invariant odd Fredholm operators	$E_{-3} =$								$i\hat{T}\hat{C}\beta$		
	$E_{-2} =$					$i\hat{C}\beta$			$i\hat{C}\beta$		
	$E_{-1} =$		$\hat{P}\beta$			$\hat{C}\beta$		$\hat{C}\beta$	$\hat{C}\beta$		
	$E_{+0} =$	β	β	β	$\begin{pmatrix} \beta & 0 \\ 0 & -\beta \end{pmatrix}$	β	β	β	β	β	β
	$E_{+1} =$				$\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$		$\hat{C}\beta$			$\hat{C}\beta$	$\hat{C}\beta$
	$E_{+2} =$				$\begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}$		$i\hat{C}\beta$			$i\hat{C}\beta$	
	$E_{+3} =$				$\begin{pmatrix} 0 & -\hat{T} \\ \hat{T} & 0 \end{pmatrix}$					$i\hat{T}\hat{C}\beta$	
	$E_{+4} =$				$\begin{pmatrix} 0 & i\hat{T} \\ i\hat{T} & 0 \end{pmatrix}$						
τ -twisted G -equivariant KR-theory of fixed loci	$KR^\tau =$	KU^0	KU^1	KO^0	KO^4	KO^2	KO^6	KO^1	KO^3	KO^5	KO^7

$$\left\{ \begin{array}{ll} \text{bounded ops.} & \widehat{F} : \mathcal{H}^2 \xrightarrow[\mathbb{K}\text{-linear}]{\text{bounded}} \mathcal{H}^2 \\ \text{self-adjoint} & \widehat{F}^* = \widehat{F} := F + F^* \\ \text{Fredholm} & \dim(\ker(\widehat{F})) < \infty \end{array} \right\} \quad \left\{ \begin{array}{ll} \text{graded comm.} & E_0, \dots, E_p : \mathcal{H}^2 \xrightarrow[\mathbb{K}\text{-linear}]{\text{bounded}} \mathcal{H}^2 \\ E_i \circ \widehat{F} = -\widehat{F} \circ E_i & \text{with (anti-)self-adjoint } (E_i)^* = \text{sgn}_i \cdot E_i \\ \text{Clifford gen.} & E_i \circ E_j + E_j \circ E_i = 2\text{sgn}_i \cdot \delta_{ij} \end{array} \right\}$$

$=: \text{Fred}_{\mathbb{C}}^{-p}$

[Karoubi 70]: $\left\{ X \xrightarrow{\text{cnts}} \text{Fred}_{\mathbb{K}}^p \right\} / \sim_{\text{htpy}} = \begin{cases} \text{KU}^p(X) = \text{KU}^{p+2}(X) & | \quad \mathbb{K} = \mathbb{C} \\ \text{KO}^p(X) = \text{KO}^{p+8}(X) & | \quad \mathbb{K} = \mathbb{R} \end{cases}$

Maximal induced Clifford action anticommuting with all G -invariant odd Fredholm operators

$E_{-3} =$									$i\widehat{T}\widehat{C}\beta$		
$E_{-2} =$						$i\widehat{C}\beta$			$i\widehat{C}\beta$		
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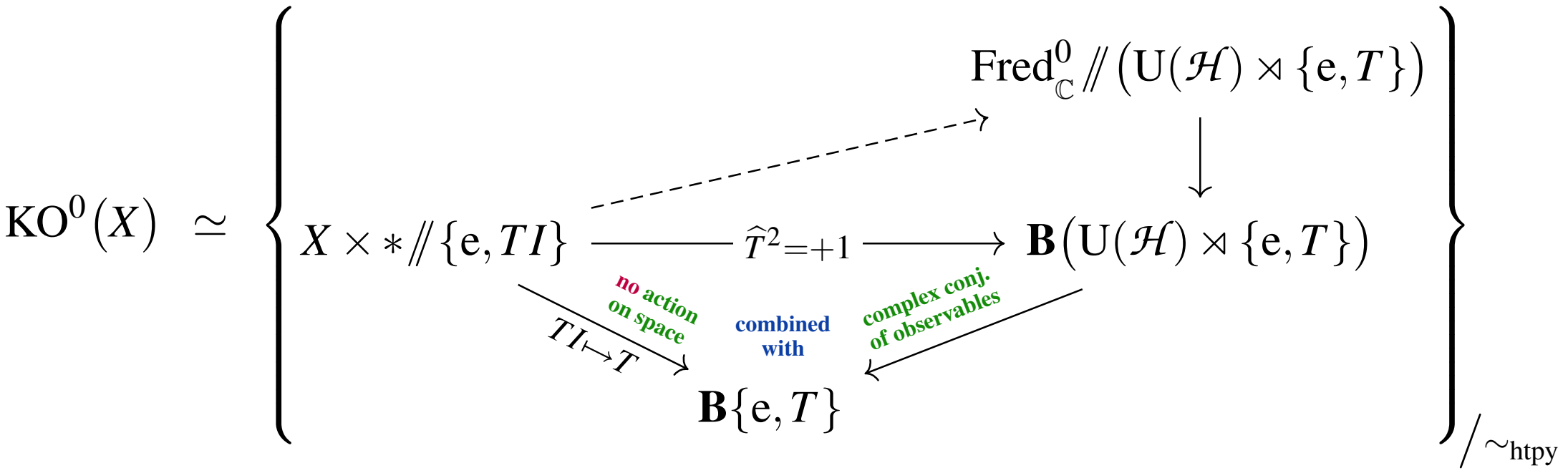
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$\text{KR}^\tau =$	KU^0	KU^1	KO^0	KO^4	KO^2	KO^6	KO^1	KO^3	KO^5	KO^7
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Example – TI -equivariant KR-theory is KO^0 -theory.

The combination $T \cdot I$ acts trivially on the domain space and by complex conjugation on observables.

Hence $(T \cdot I)$ -equivariant $(\widehat{T}^2 = +1)$ -twisted KR-theory is KO^0 -theory:

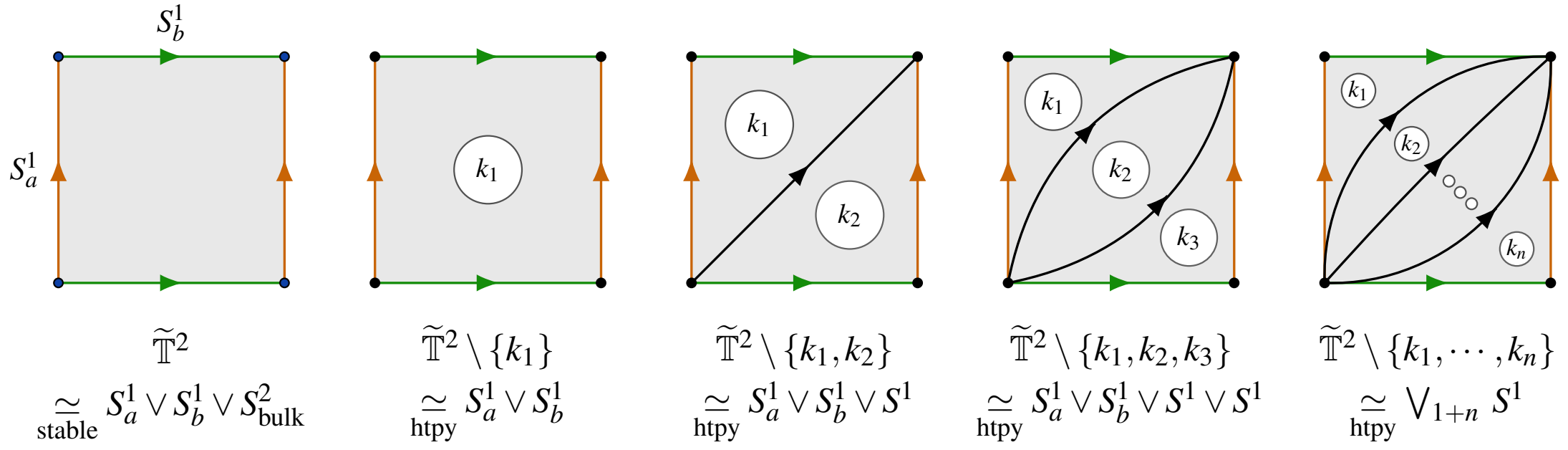


$n =$	0	1	2	3	4	5	6	7	8	9	...
$KO^0(S_*^n) =$	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	0	\mathbb{Z}	0	0	0	\mathbb{Z}	\mathbb{Z}_2	...

Example – TI -equivariant KR-theory of punctured torus.

So the TI -equivariant $(\widehat{T}^2 = +1)$ -twisted KR-theory of the N -punctured torus is

$$\begin{aligned}
 & \text{KR}^{(\widehat{T}^2 = +1)}(\widetilde{\mathbb{T}}^2 \setminus \{k_1, \dots, k_N\}) \\
 & \simeq \text{KO}^0(\widetilde{\mathbb{T}}^2 \setminus \{k_1, \dots, k_N\}) \\
 & \simeq \text{KO}^0\left(\bigvee_{1+N} S_*^1\right) \quad (N \geq 1) \\
 & \simeq \bigoplus_{1+N} \mathbb{Z}_2
 \end{aligned}$$



The B-field twist.

Besides these CPT-quantum symmetries,

K-theory generically admits the famous *twisting by a B-field*:

The homotopy fiber sequence of 2-stacks discussed before

$$\mathbf{BU}(\mathcal{H}) \longrightarrow \mathbf{B}(\mathbf{U}(\mathcal{H})/\mathbf{U}(1)) \xrightarrow{\text{universal Dixmier-Douady class}} \mathbf{B}^2\mathbf{U}(1)$$

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induces a surjection of equivalence classes of equivariant higher bundles

$$\begin{array}{ccc} \text{equivariant projective bundles} & & \text{equivariant bundle gerbes} \\ \pi_0 \text{ Map}\left(\widehat{\mathbf{X}}//G, \mathbf{B}(\mathbf{U}(\mathcal{H})/\mathbf{U}(1))\right) & \xrightarrow{\text{DD}_*} & \pi_0 \text{ Map}\left(\widehat{\mathbf{X}}//G, \mathbf{B}^2\mathbf{U}(1)\right) \end{array}$$

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which has a natural section:

$$\pi_0 \text{ Map}\left(\widehat{\mathbf{X}}//G, \mathbf{B}^2\mathbf{U}(1)\right) \hookrightarrow \pi_0 \text{ Map}\left(\widehat{\mathbf{X}}//G, \mathbf{B}\left(\frac{\mathbf{U}(\mathcal{H}) \times \mathbf{U}(\mathcal{H})}{\mathbf{U}(1)} \rtimes (\{e, C\} \times \{e, P\})\right)\right).$$

equivariant bundle gerbes **“stable twists”** **full quantum-symmetry twists**

The B-field twist – Inner local systems.

On fixed loci (orbi-singularities)

$$X // G \simeq X \times * // G = X \times \mathbf{B}G$$

the B-field twist induces *secondary* twists by “inner local systems”:

stable twists over fixed locus

$$\begin{aligned} \text{Map}(X \times * // G, \mathbf{B}^2\mathbf{U}(1)) &\simeq \text{Map}(X \times \mathbf{B}G, \mathbf{B}^2\mathbf{U}(1)) \\ &\simeq \text{Map}(X, \text{Map}(\mathbf{B}G, \mathbf{B}^2\mathbf{U}(1))) \end{aligned}$$

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Here we are assuming $G \subset_{\text{fin}} \text{SU}(2)$ so that $H_{\text{Grp}}^2(G, \mathbf{U}(1)) = 0$.

And $G^* := \text{Hom}(G, \mathbf{U}(1))$ denotes the Pontrjagin-dual group.

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The B-field twist – Inner local systems – The diagrammatics.

Hence the

inner local system-twisted KU-cohomology of a G-orbi-singularity of shape X

arises as follows:

$$\text{KU}_G^{n+[\omega_1]}(X) = \left\{ \begin{array}{ccc} & & \text{Fred}_{\mathbb{C}}^n // \text{PU}(\mathcal{H}) \\ & \nearrow^{\text{cocycle}} & \downarrow \\ X \times * // G & \xrightarrow[\text{inner local system twist}]{\tau} & \text{BPU}(\mathcal{H}) \end{array} \right\} / \sim_{\text{htpy}}$$

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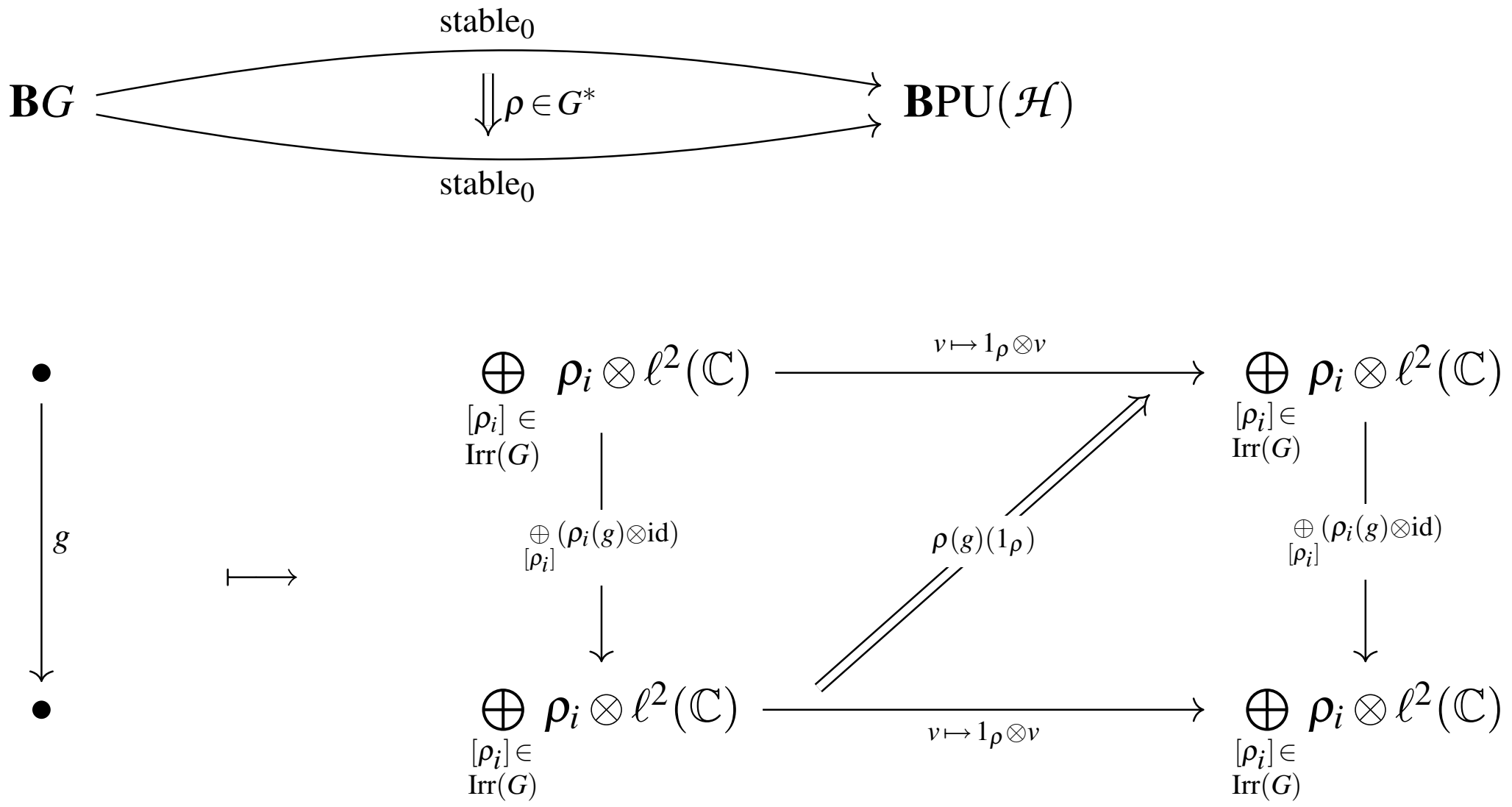
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 \swarrow \text{cocycle} & \downarrow & \downarrow \\
 X & \xrightarrow{\omega_1} & \mathbf{B}G^* \xrightarrow{\text{automorphisms of univ. stable equiv. twist}} \text{Map}(\mathbf{B}G, \mathbf{B}\text{PU}(\mathcal{H})) \\
 \text{inner local system} & & \text{(pb)}
 \end{array} \right\} / \sim_{\text{htpy}}$$

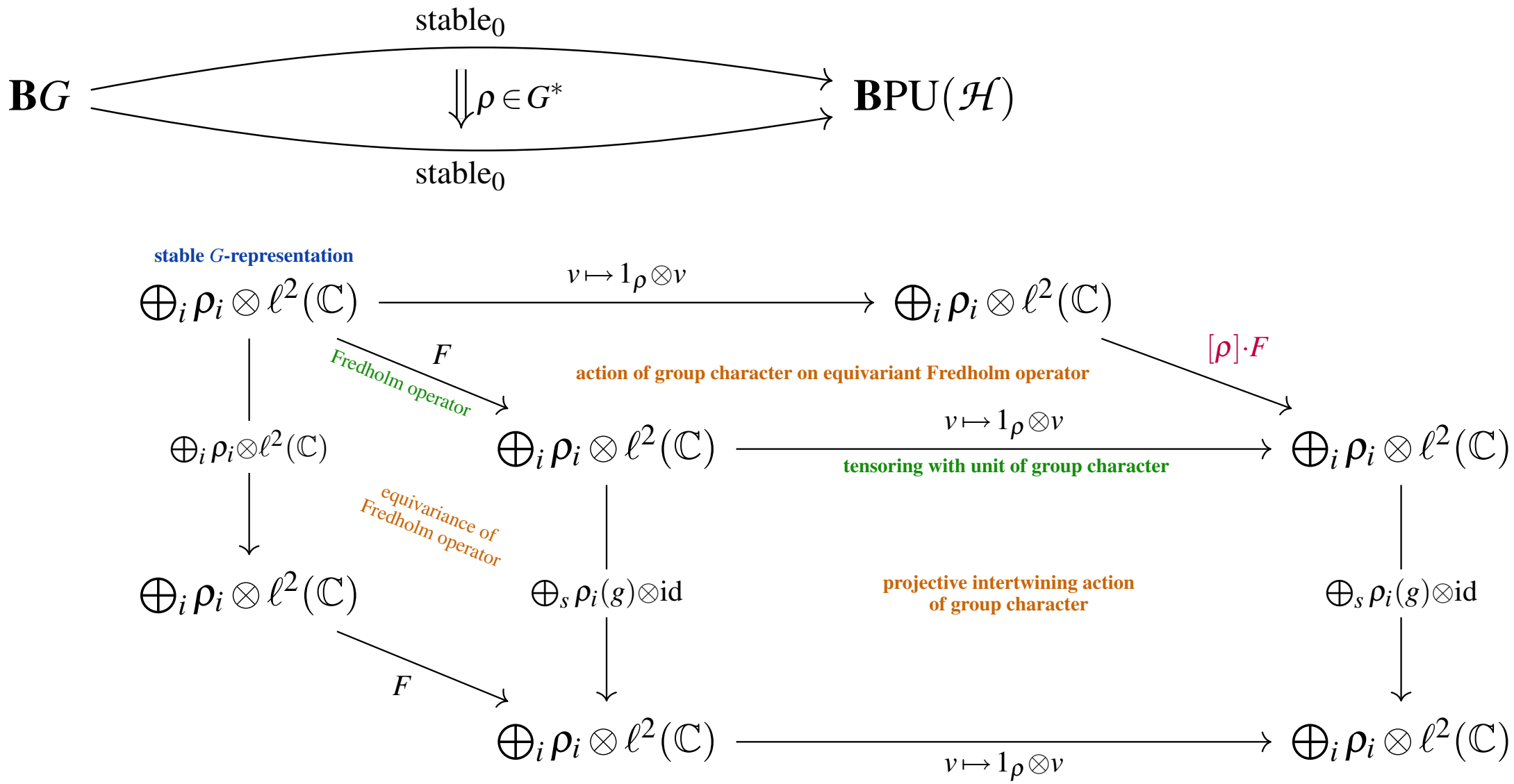
The B-field twist – Inner local systems – The proof.

For the proof we consider the following diagram [SS22-Bun, Ex. 4.1.56][SS22, §3]:



The B-field twist – Inner local systems – The proof.

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The B-field twist – Inner local systems – Chern character.

One aspect of these twistings becomes transparent under the *Chern character*:

complex K-theory

$$\mathrm{KU}^0(X) \xrightarrow{\text{Chern character}} \mathrm{KU}^0(X; \mathbb{C}) \simeq \bigoplus_{d \in \mathbb{N}} H^{2d} \left(\Omega_{\mathrm{dR}}^\bullet(X; \mathbb{C}), d \right)$$

periodic de Rham cohomology

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periodic de Rham cohomology

For twist by B-field \widehat{B}_2 there is a closed differential 3-form H_3 such that:

plain B-field

-twisted K-theory

$$\mathrm{KU}^{n+\widehat{B}_2}(X) \xrightarrow{\text{twisted Chern character}} \mathrm{KU}^{\widehat{B}_2}(X; \mathbb{C}) \simeq \bigoplus_{d \in \mathbb{Z}} H^{n+2d} \left(\Omega_{\mathrm{dR}}^\bullet(X; \mathbb{C}), d + H_3 \wedge \right)$$

3-twisted periodic de Rham cohomology

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plain B-field 3-twisted periodic de Rham cohomology
 -twisted K-theory

$$\text{KU}^{n+\widehat{B}_2}(X) \xrightarrow[\text{Chern character}]{\text{twisted}} \text{KU}^{\widehat{B}_2}(X; \mathbb{C}) \simeq \bigoplus_{d \in \mathbb{Z}} H^{n+2d} \left(\Omega_{\text{dR}}^\bullet(X; \mathbb{C}), d + H_3 \wedge \right)$$

For twist by inner C_κ -local system, there is closed 1-form ω_1 with holon. in $C_\kappa \subset U(1)$ such that:

inner local system 1-twisted periodic de Rham cohomology
 -twisted K-theory

$$\text{KU}_{C_\kappa}^{n+[\omega_1]}(X) \xrightarrow[\text{Chern character}]{\text{twisted equivariant}} \bigoplus_{\substack{d \in \mathbb{Z} \\ 1 \leq r \leq \kappa}} H^{n+2d} \left(\Omega_{\text{dR}}^\bullet(X; \mathbb{C}), d + r \cdot \omega_1 \wedge \right)$$

of A-type singularity

The B-field twist – Inner local systems – Chern character.

One aspect of these twistings becomes transparent under the Chern character:

This is the hidden 1-twisting in TED-K – that we will next relate to anyons. \longrightarrow

<p>inner local system -twisted K-theory</p>		<p>1-twisted periodic de Rham cohomology</p>
$KU_{C_\kappa}^{n+[\omega_1]}(X)$	$\xrightarrow[\text{Chern character}]{\text{twisted equivariant}}$	$\bigoplus_{\substack{d \in \mathbb{Z} \\ 1 \leq r \leq \kappa}} H^{n+2d} \left(\Omega_{\text{dR}}^\bullet(X; \mathbb{C}), \mathbf{d} + r \cdot \omega_1 \wedge \right)$
<p>of A-type singularity</p>		

0 – Cohesive ∞ -Topoi

I – Equivariant ∞ -Bundles

II – TED-K-Theory

III – Anyonic Defect Branes

IV – Quantum Computation

This part is a brief indication
of a few aspects discussed in:

Anyonic Defect Branes in TED-K-Theory [arXiv:2203.11838]

K-Theory classifies non-perturbative vacua.

Solid state physics	K-theory	String theory

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Solid state physics	K-theory	String theory
Single electron state	Line bundle	Single D-brane
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Valence bundle of electron/positron states	Virtual bundle of their kernels and cokernels	D-brane Sen vacuum after tachyon condensation
Topological phase	K-theory class	Stable D-brane charge

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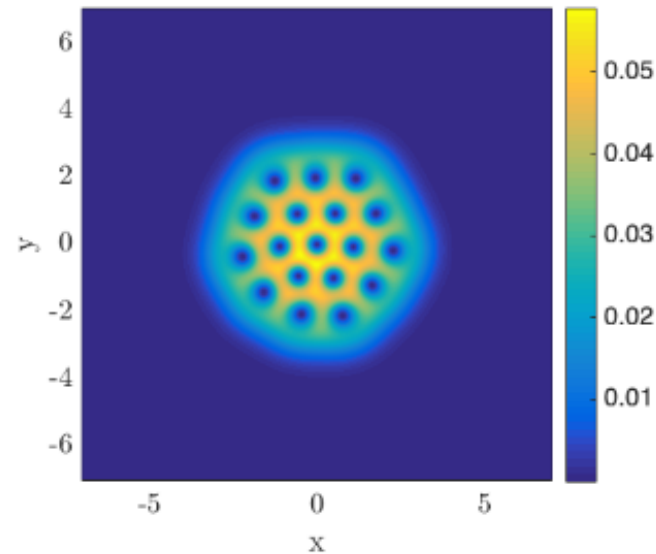
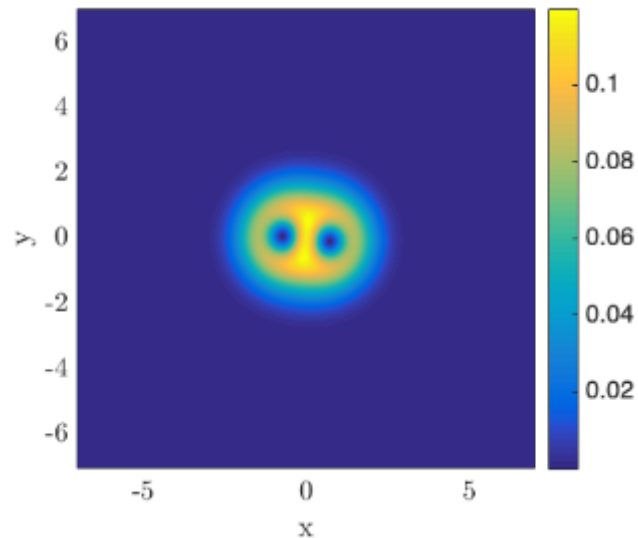
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Internal symmetry	Inner local system-twist	Inside of orbi-singularity
Anyons	Punctures	Defect branes

Anyons in condensed matter & string theory.

In solid state physics

anyons are presumed pointlike defects in gapped topological phases of effectively 2-dimensional materials whose adiabatic dynamics is that of Wilson lines in $\mathfrak{su}(2)$ -CS theory.



(numerical simulation from arXiv:1901.10739)

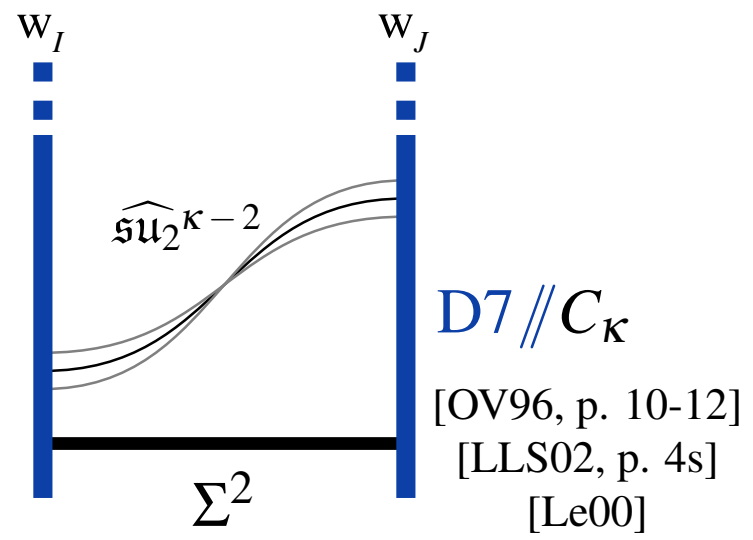
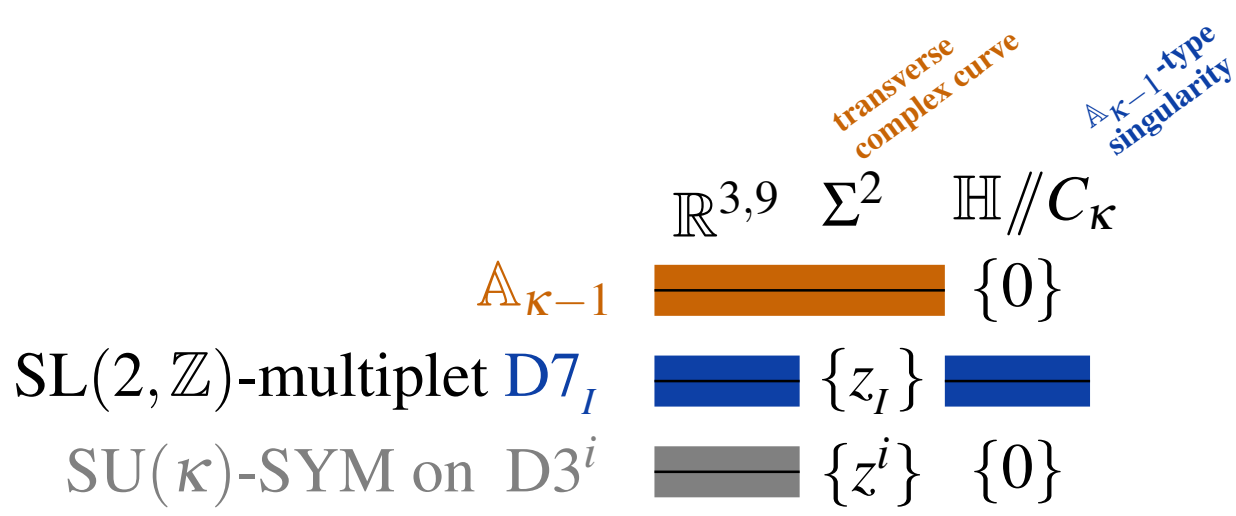
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Concretely, it is expected that:

$$\left\{ \begin{array}{l} \text{ground state wave functions of} \\ \text{spin}=\mathbf{w}_I \quad \widehat{\mathfrak{su}}_2^k\text{-anyons at} \\ \text{positions } z_I \text{ in transverse plane} \end{array} \right\} \simeq \text{space of "conformal blocks"} \\ \simeq \text{ConfBlck}_{\widehat{\mathfrak{sl}}_2^k}(\vec{\mathbf{w}}, \vec{\mathbf{z}})$$

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As the positions z_I move, these spaces constitute braid group representations.

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Yes! \longrightarrow

TED-Cohomological incarnation of Conformal blocks.

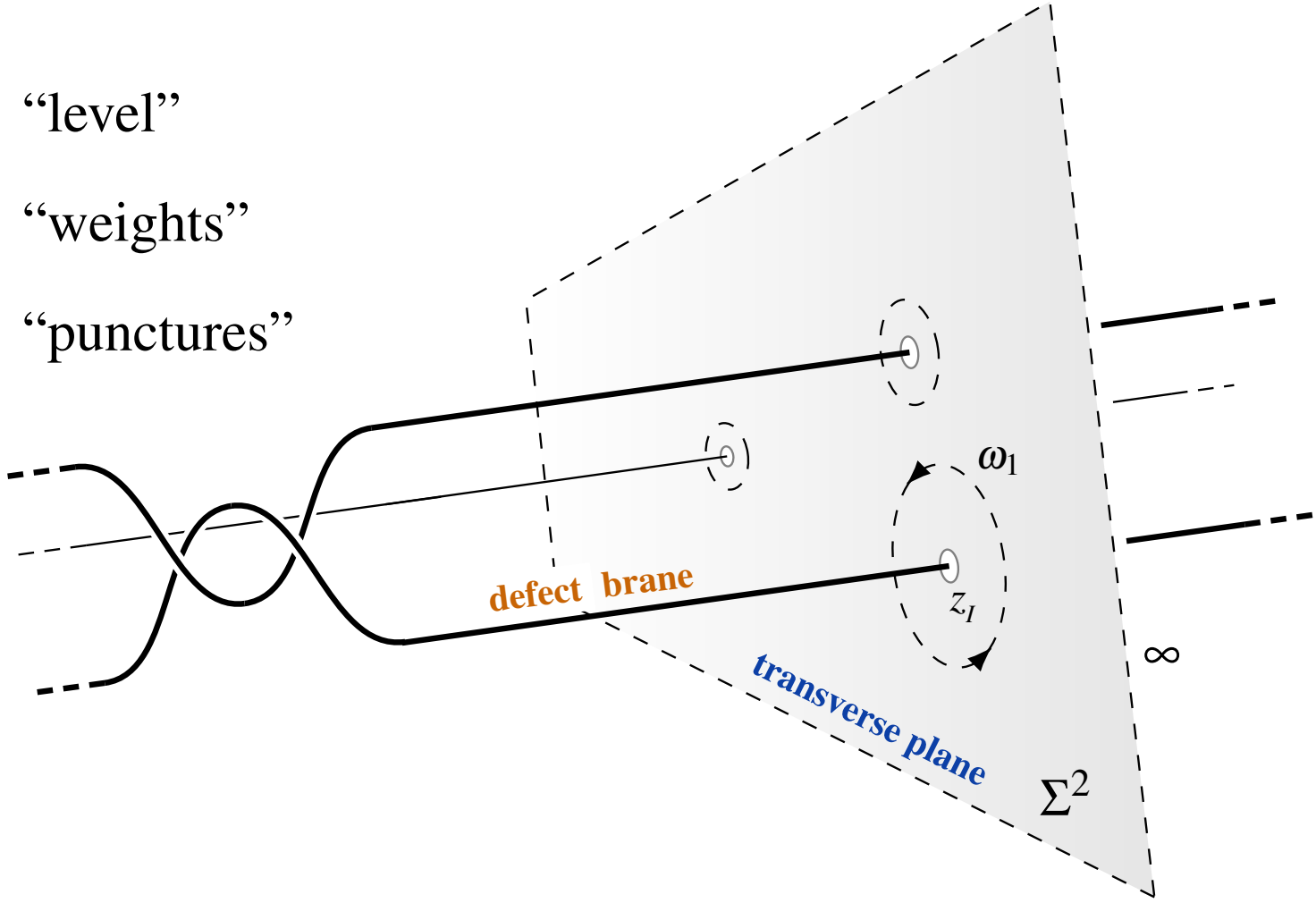
Consider

$\kappa := k + 2$ “level”

$w_I \in \{0, \dots, k\}$ “weights”

$z_I \in \{z_1, \dots, z_N\}$ “punctures”

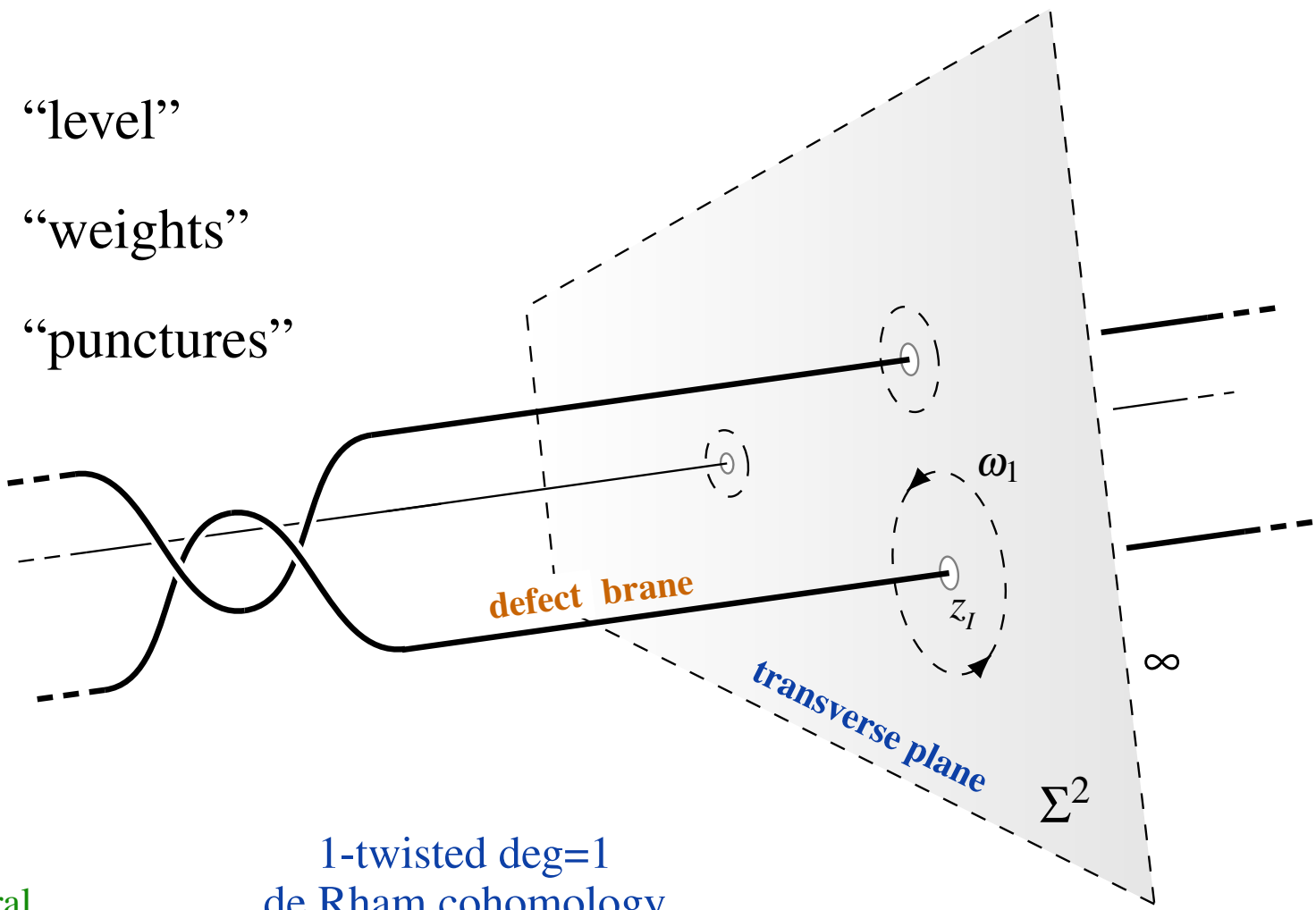
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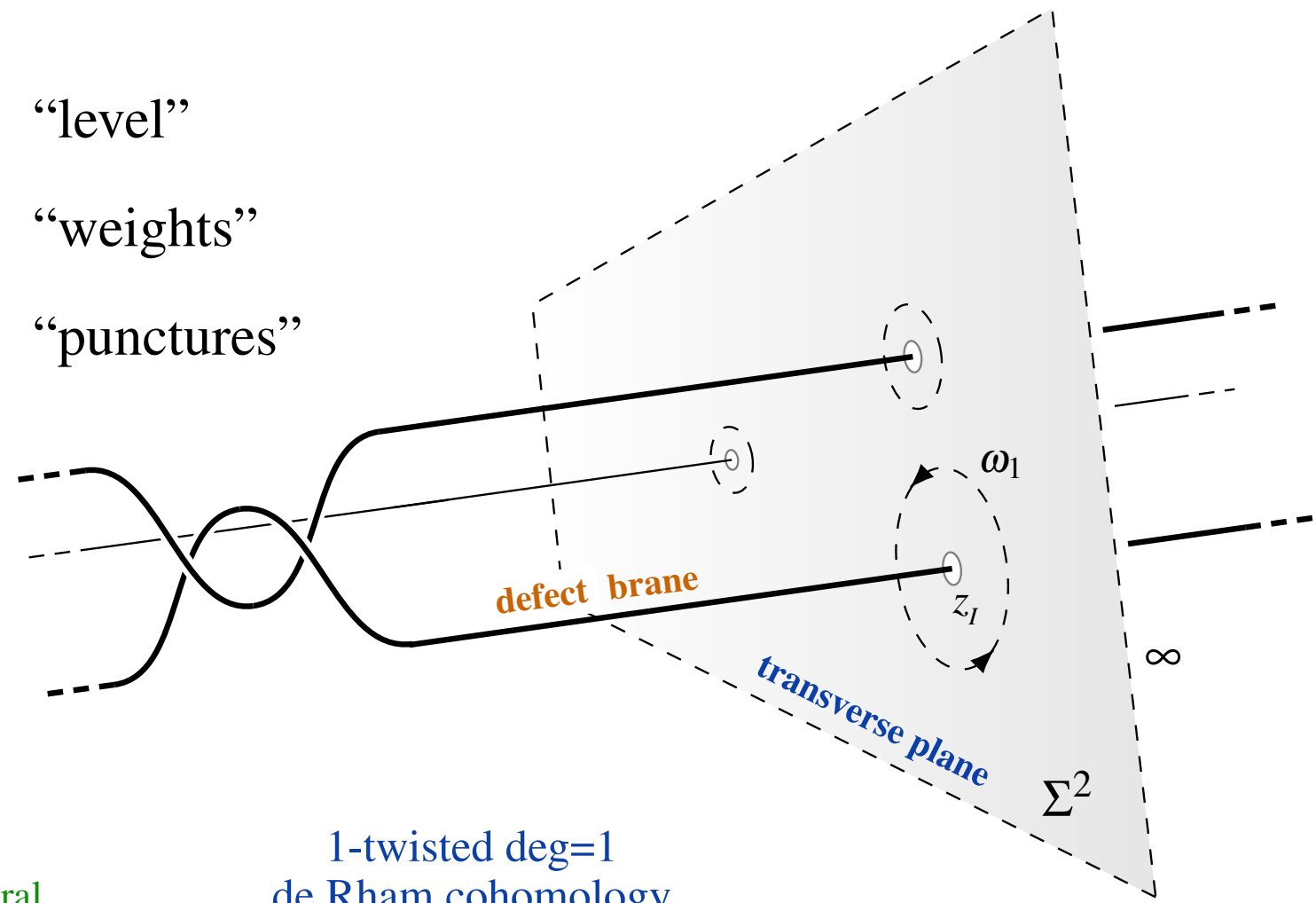
$\mathfrak{su}(2)$ -affine deg=1 conformal blocks $\text{CnfBlck}_{\widehat{\mathfrak{sl}_2^k}}^1(\vec{w}, \vec{z})$
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 $\hookrightarrow H^1\left(\Omega_{\text{dR}}^\bullet(\mathbb{C} \setminus \{\vec{z}\}), d + \omega_1 \wedge\right)$
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1-twisted deg=1 de Rham cohomology

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$\xrightarrow{\text{natural inclusion}}$	$\xrightarrow{\text{natural inclusion}}$	$\text{KU}^{1+\omega_1}\left(\left(\mathbb{C} \setminus \{\vec{z}\}\right) \times * // C_\kappa; \mathbb{C}\right)$	[SS22, Prop. 2.16]
		inner local system-twisted deg=1 K-theory of $\mathbb{A}_{\kappa-1}$ -singularity	(as explained <u>above</u>)

TED-Cohomological incarnation of Conformal blocks.

Generally, consider *configuration spaces of points* (e.g. [SS19, §2.2])

$$\text{Conf}_{\{1, \dots, n\}}(\mathbf{X}) := \left\{ z^1, \dots, z^n \in \mathbf{X} \mid \forall_{i < j} z^i \neq z^j \right\}.$$

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$\mathfrak{su}(2)$ -affine deg= n
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1-twisted deg= n de Rham cohomology
of configuration space of n points

$$\text{CnfBlck}_{\widehat{\mathfrak{sl}}_2^k}^n(\vec{w}, \vec{z}) \hookrightarrow H^n \left(\Omega_{\text{dR}}^\bullet \left(\text{Conf}_{\{1, \dots, n\}}(\mathbb{C} \setminus \{\vec{z}\}) \right), d + \omega_1 \wedge \right) \quad [\text{FSV94, Cor. 3.4.2}]$$

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inner local system-twisted deg= n K-theory of configurations in $\mathbb{A}_{\kappa-1}$ -singularity

The previous statement is subsumed since $\text{Conf}_{\{1\}}(\mathbf{X}) = \mathbf{X}$.

Conclusion.

The commonly expected $\widehat{\mathfrak{su}}_2^k$ -charges of anyons and defect branes are reflected in the TED-K-theory of *configuration spaces of points* in 2-dimensional transverse spaces *inside* \mathbb{A}_{k+1} -*orbi-singularities*.

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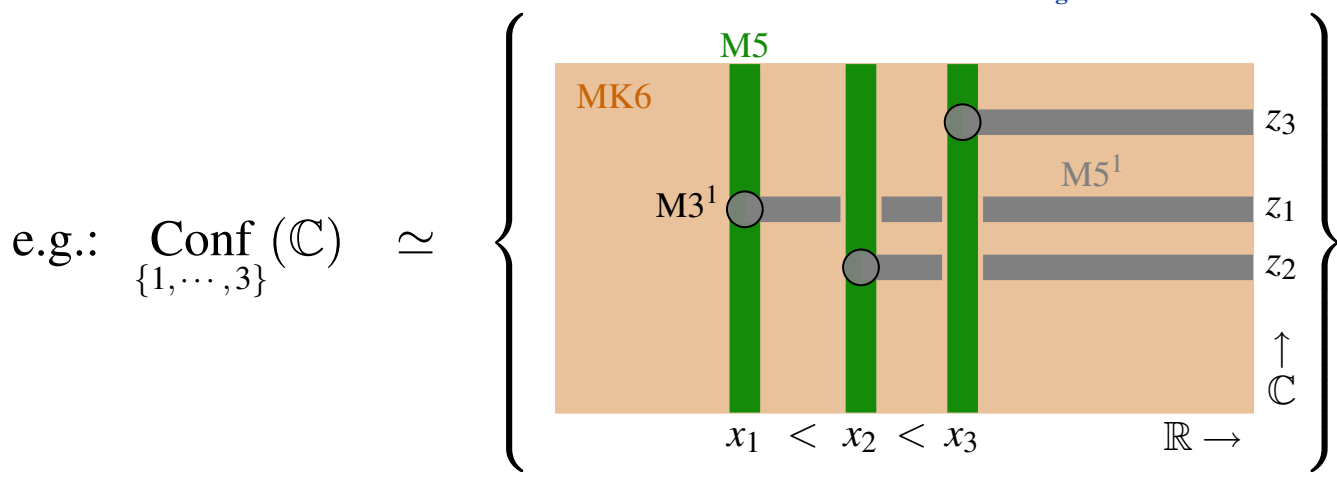
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This is compatible with traditional brane charge quantization (only) in degree 1 while in general degree it is compatible under Hypothesis H, which asserts [SS19] that quantum states of branes are in the generalized cohomology of *Cohomotopy cocycle spaces* of spacetime:

$$\begin{array}{c}
 \text{Configuration space of} \\
 \text{ordered points in the plane} \\
 \coprod_n \text{Conf}_{\{1, \dots, n\}}(\mathbb{C})
 \end{array}
 \simeq
 \begin{array}{c}
 \text{3-Cohomotopy cocycle space} \\
 \text{for codim=1 branes} \\
 \text{Map}^*(\mathbb{R}_+ \wedge \mathbb{C}_{\text{cpt}}, \mathcal{S}^3) \simeq \\
 \overbrace{\coprod_n \text{Conf}_n(\mathbb{C}; \mathbb{R}_{\text{cpt}})} \\
 \times \\
 \overbrace{\coprod_n \text{Conf}_n(\mathbb{R}; \mathbb{C}_{\text{cpt}})} \\
 \text{3-Cohomotopy cocycle space} \\
 \text{for codim=2 branes} \\
 \text{Map}^*(\mathbb{R}_{\text{cpt}} \wedge \mathbb{C}_+, \mathcal{S}^3) \simeq
 \end{array}
 \begin{array}{c}
 \text{Fiber product of respective configuration spaces} \\
 \text{(of un-ordered points escaping to transverse infinity)} \\
 \text{reflecting the brane intersections} \\
 \coprod_n \text{Conf}_n(*; (\mathbb{R} \times \mathbb{C})_{\text{cpt}})
 \end{array}$$



The moduli space of flat M3-branes according to Hypothesis H is the configuration space of ordered points in their transverse plane.

0 – Cohesive ∞ -Topoi

I – Equivariant ∞ -Bundles

II – TED-K-Theory

III – Anyonic Defect Branes

IV – Quantum Computation

For expository outlook on

Topological Quantum Programming in TED-K [arXiv:2209.08331]

see:

ncatlab.org/nlab/files/CQTS-InitialResearcherMeeting-Schreiber-220914.pdf

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Twisted, Equivariant, Differential (TED) K-Theory via Diffeological Stacks

Urs Schreiber on joint work with Hisham Sati

NYU AD Science Division, Program of Mathematics

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Thanks!

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