

Kevin Costello 1

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work with Owen G. Williams

What is the analogue of deformation quantization in QFT instead of quant. mechanics?

Deformation quantization

Classical mechanics.

A^{cl} , comm. algebra,

$\{ , \}$ Poisson bracket.

To quantize this, need to find an assoc. algebra A^q over the ring $\mathbb{R}[[\hbar]]$, such that

1. $A^q / \hbar A^q = A^{cl}$.

2. If $a, b \in A^{cl}$, \tilde{a}, \tilde{b} are lifts to A^q ,

then

$$\{\hat{a}, \hat{b}\} = \frac{1}{\hbar} [\tilde{a}, \tilde{b}] \pmod{\hbar}.$$

These lectures: want to give an analog of this picture for QFT.

- 1) Need to explain what plays the role of commutative, Poisson and associative algebras.
- 2) ~~How to get~~
Explain how classical field theory is encoded in (analog of) commutative + Poisson.
- 3) Explain how to quantize the classical field theory.

Today: the structure that plays the role of assoc. alg is a factorization algebra.

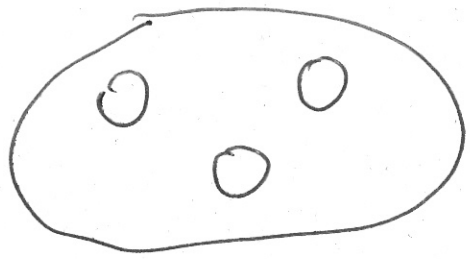
Not so terrifying as in book by Beilinson Drinfeld.

We work with a C^∞ (smooth) analogue of their stuff.

Let M be an manifold ("spacetime") on which we want to do QFT.

Let $B(M) = \left\{ \begin{array}{l} \text{closed} \\ \text{balls in } M \end{array} \right\}$... on ∞ -dim manifold

Let $B_n(M) = \left\{ \begin{array}{l} n \text{ disjoint balls} \\ \text{embedded in a larger ball} \end{array} \right\}$ if we look at Riemannian manifold



$$e \in B_3(\mathbb{R}^2)$$

proj: in practice, locally triv. bundles of ∞ -dim top. vector spaces

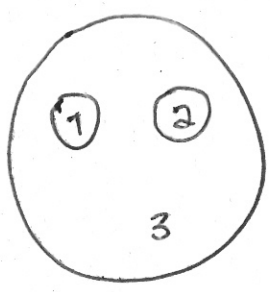
A factorization algebra is a $\hat{}$ vector bundle F on $B(M)$ (S_n-invariant) equipped with maps $p^*(F^{\boxtimes n}) \rightarrow q^*(F)$ satisfying some compatibility.

projection maps

$$B(M) \xleftarrow[\text{q}]{\text{ball on outside}} B_n(M) \xrightarrow{\text{p}} B(M)^n$$

Concretely, $F^{\hat{}}$ smoothly assigns a vector space to every ball $B \subset M$.

If we have some configuration of balls, like



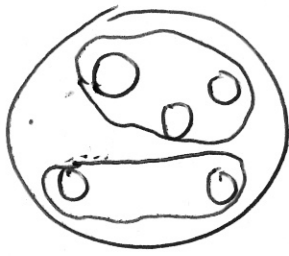
we get a map

$$F(B_1) \otimes F(B_2) \longrightarrow F(B_3) \quad \text{of the balls}$$

These maps must vary smoothly as the configuration varies

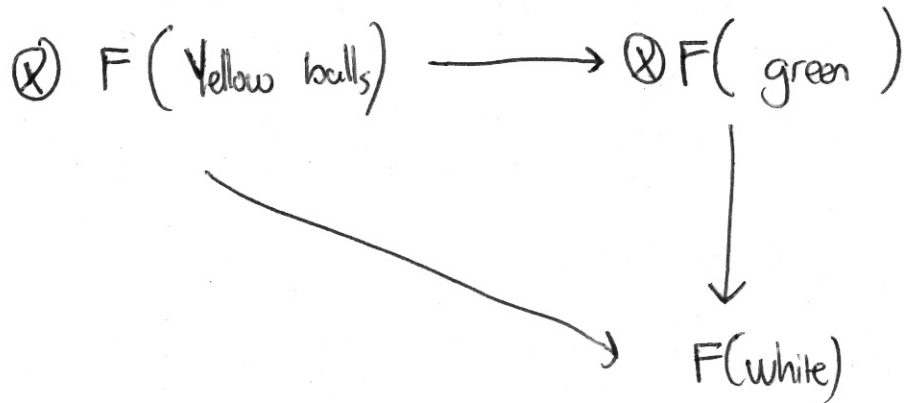
Compatibility condition: like for little discs operad.

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it doesn't matter if you apply
map all in one go or twice.

i.e. the diagram



Commutates.

(ie. its an algebra over the little discs operad (Urs)
Yes... absolutely. But it's a colored operad (K. Costello).)

This is an algebra over colored operad, colours are $B(M)$,

n-ary operations are $B_n(M)$. ~~then~~

with extra condition saying the operations must depend
smoothly, i.e. the vector space we assign to

each color forms a smooth V. bundle and
the maps are compatible with this.

Explanation

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Vector bundles come in 3 flavours:

1. C^∞ \rightsquigarrow factorization algebra
2. holomorphic \rightsquigarrow vertex algebra
3. Locally constant sheaves \rightsquigarrow E_∞ -algebra

~~Even~~ Factorization algebras exist in settings 2+3 also.

Defⁿ ^(vague) A locally constant factorization algebra is like a factorization alg., except F is a locally constant sheaf on $B(M)$, and the structure maps are maps of l.c. sheaves.

Q: sheaves of what?

A: We drink the derived coaloid. Everything is a cochain complex.

Example Let F be a locally constant fact. alg. on \mathbb{R}^n .

$B(\mathbb{R}^n)$ is contractible,

F is quasi-iso to a trivial sheaf, with fibre V , a cochain complex.

If have



get a map
 $V \otimes V \rightarrow V$

As ~~these~~ vary, the conf. of discs vary,
the products change by homotopies.

i.e. a locally constant f. algebra ^{on \mathbb{R}^n} is the same thing as
an E_∞ -algebra.

i.e. the set of colours is contractible.

the set of operations \cong little discs operad.

Next : Holomorphic factorization algebras.

(Here,
For Riemann surfaces).

Let Σ be a Riemann surface.

We know what it means for a map from a
complex manifold to $B(\Sigma)$ to be holomorphic.

(this is why go to 2d, can use
~~Riemann~~ ^{open} mapping theorem)

So, we can talk about hol. objects on $B(\Sigma)$.

(What is a map $U \rightarrow B(\Sigma)$? Its a complex man.
 $M \rightarrow \Sigma$
 \downarrow
 U where fibers of M are open balls)

(i.e. in the sense of a sheaf on
the site of complex manifold)

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André : problem with defn. Everything is holomorphic!

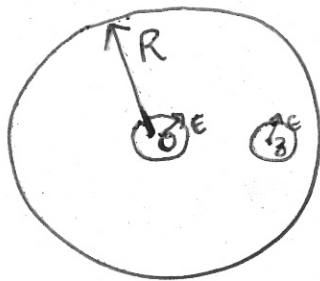
Kevin : you're prob. right. I apologize. Still nontrivial for maps
to be holomorphic.

I think about smooth examples normally.

Let's consider a holomorphic fact. algebra on \mathbb{C} which
is translation invariant, and dilation invariant.

Let $V = F(\text{ony round disc})$.

If we have a configuration of discs like



We get a map

$$m_z : V \otimes V \rightarrow V$$

As z varies, with fixed
radius, this map ~~varies~~ must vary holomorphically.

m_z is a holomorphic map

$$\text{Annulus} \longrightarrow \text{Hom}(V \otimes V, V)$$

So it has a Laurent expansion:

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$$m_3 \sim \sum_{k \in \mathbb{Z}} z^k a_k,$$

$a_k \in$ some completion of $\text{Hom}(V \otimes V, V)$

If I had got defn right, we'd have seen this is reminiscent of a VOA.

In Beilinson + Drinfeld's book, they make this defn in the algebraic setting, call it a chiral algebra.

They show axioms for a chiral algebra on \mathbb{C} are "essentially equivalent" to those of a vertex algebra.

Question: What about finite Laurent condition?

Have described field-state.

Answer: I haven't worked this through. But in B-D book, they do get a correspondence.

Urs: What if you vary radius?

Answer: Get a filtered object, more along the lines of Huang.

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Two situations in QFT where observables well understood

- (1) holomorphic / conformal field theory (incl. observables e.g. Segal)
- (2) quantum mechanics (observables on \mathbb{R}).

In general, want factorization algebra on M .

Note Deformation quantization is a perturbative thing.

We don't know what QFT is nonperturbatively.

