

### Exercise 80

$E$  is closed:

$$\begin{aligned}dE &= \frac{1}{x^2 + y^2}(dx \wedge dy - dy \wedge dx) \\ &= 0\end{aligned}$$

To calculate the integrals, note that the path lies on a 1-dimensional space which can be reparametrised with

$$\begin{aligned}x &= \cos \theta \\ y &= \sin \theta\end{aligned}$$

The 1-form then becomes

$$\frac{\cos \theta \cos \theta - \sin \theta(-\sin \theta)}{\cos^2 \theta + \sin^2 \theta} = 1$$

and the integrals are:

$$\begin{aligned}\int_{\gamma_0} E &= \int_{\pi}^0 1d\theta \\ &= -\pi\end{aligned}$$

and

$$\begin{aligned}\int_{\gamma_1} E &= \int_{-\pi}^0 1d\theta \\ &= \pi\end{aligned}$$

### Exercise 81

Let  $\gamma_0$  and  $\gamma_1$  be two paths from  $p$  to  $q$ . Define the homotopy  $\gamma$  as

$$\gamma(s, t) = (1 - s)\gamma_0(t) + s\gamma_1(t)$$

Then

$$\gamma(0, t) = \gamma_0(t)$$

and

$$\gamma(1, t) = \gamma_1(t)$$

Also

$$\gamma(s, 0) = (1 - s)p + sp = p$$

and similarly  $\gamma(s, 1) = q$ . Finally,  $\gamma$  is smooth.

### Exercise 82

Suppose the 1-form  $E$  is exact. Then there is a function  $\phi$  on  $M$  such that  $E = d\phi$ . Let  $\gamma$  be any loop on  $M$ . Denote  $\gamma(0) = \gamma(1)$  by  $p \in M$ . Then

$$\begin{aligned}\int_{\gamma} E &= \int_{\gamma} d\phi \\ &= \phi(\gamma(1)) - \phi(\gamma(0)) \\ &= \phi(p) - \phi(p) \\ &= 0\end{aligned}$$

Conversely, suppose  $\int_{\gamma} E = 0$  for all loops  $\gamma$  on  $M$ . This implies that

$$\int_{\gamma_0} E = \int_{\gamma_1} E$$

for any two paths  $\gamma_0$  and  $\gamma_1$  both starting at  $p \in M$  and ending at  $q \in M$ :

$$\int_{\gamma_1} E - \int_{\gamma_0} E = \int_{\gamma_0\gamma_1^{-1}} E = 0$$

since  $\gamma_0\gamma_1^{-1}$  is a loop on  $M$ .

But this means that we can follow the exact same steps as on p109 to construct a function  $\phi$  on  $M$  such that  $E = -d\phi$ , showing that  $E$  is exact.

### Exercise 83

Parametrize  $S^1$  by  $\theta \in [0, 1] \subset \mathbb{R}$  such that  $\theta(0) = \theta(1)$ . Choose any two charts, say  $V_1 = (-0.1, 0.6)$  and  $V_2 = (0.4, 1.1)$ . Then construct a 1-form on  $S^1$  which is  $d\theta$  on either of the charts.

If  $\{U_i\}$  is an atlas of charts on  $M$ , then  $\{V_1 \times U_i\} \cup \{V_2 \times U_i\}$  is an atlas for  $S^1 \times M$  and on each chart,  $d\theta$  is still a 1-form. Integrating it around the loop  $S^1 \times \{m\}$  gives a non-zero value.

### Exercise 84

Consider the open subset  $U_1^+ \subset D^n$ , where

$$U_1^+ = \{x \in D^n : x_1 > 0\}$$

Let  $\phi : U_1^+ \rightarrow \mathbb{R}^n$  map  $(x_1, \dots, x_n)$  to

$$\frac{\|x\|}{x_1}(x_1, \dots, x_n)$$

The image of  $\phi$  is

$$\{x \in \mathbb{R}^n : 0 < x_1 \leq 1\}$$

with points with norm 1 mapping to points with  $x_1 = 1$ . We can compose with another mapping from  $\mathbb{R}^n$  to itself that maps  $x_1$  to  $(1 - x_1)$ . This gives a chart which maps  $U_1^+$  to  $\mathbf{H}^n$  with the boundary points being exactly the ones with norm 1.

Finally, charts of the form  $U_i^+$  and  $U_i^-$  cover  $D^n$ , hence forming an atlas.

### Exercise 85

In chapter 3, tangent vectors at  $p \in M$  were defined as functions from  $C^\infty(M)$  to  $\mathbb{R}$  satisfying three properties. Let  $C_p^\infty(M)$  be the germ of smooth functions at  $p$ , and consider the surjective mapping from  $C^\infty(M)$  to  $C_p^\infty(M)$ . The kernel consists of functions defined on a neighbourhood of  $p$  that vanish on this neighbourhood. From the second property ( $v_p(\alpha f) = \alpha v_p(f)$ ) it follows that tangent vectors vanish on this kernel. Hence we can consider them as functions on  $C_p^\infty(M)$ , and could also have defined them like that.

Now let  $p$  be on the boundary of  $M$ . We can assume  $p \in \mathbf{H}^n$ . Let  $\mathbf{H}_\epsilon^n$  be the manifold without boundary defined like  $\mathbf{H}^n$  but with  $x^n > -\epsilon$  instead of  $x^n \geq 0$ . Then the tangent space defined at  $p \in \mathbf{H}_\epsilon^n$  using  $C_p^\infty(\mathbf{H}_\epsilon^n)$  is the usual  $n$ -dimensional vector space. And the tangent space at  $p \in \mathbf{H}^n$  defined using  $C_p^\infty(\mathbf{H}^n)$  is the space that we are interested in (for which we must show that the dimension is also  $n$ ).

But again there is a surjective map from  $C_p^\infty(\mathbf{H}_\epsilon^n)$  to  $C_p^\infty(\mathbf{H}^n)$ . The kernel consists of smooth functions defined on a neighbourhood of  $p \in \mathbb{R}^n$  that vanish for  $x^n \geq 0$ . Such functions do not necessarily vanish on a neighbourhood of  $p$ , but the tangent vector does vanish on such functions. This is essentially because at least one side of any small straight line through  $p$  must lie in  $\mathbf{H}^n$ .

Therefore this last map induces a map between the two tangent spaces which is an isomorphism.

### Exercise 86

(I googled this and found p226 of 'Introduction to manifolds' on Google books which gave most of the proof.)

Let  $\{U_\alpha\}$  be the original atlas with  $\{f_\alpha\}$  the corresponding partition of unity. Let  $\{V_\beta\}$  be another atlas where all the charts have the same orientation as in the original atlas, with  $\{g_\beta\}$  a subordinate (i.e.  $\text{support}(g_\beta) \subset V_\beta$ ) partition of unity. Then

$$\begin{aligned} \sum_\alpha \int_{U_\alpha} f_\alpha \omega &= \sum_\alpha \int_{U_\alpha} f_\alpha \left( \sum_\beta g_\beta \omega \right) \\ &= \sum_\alpha \sum_\beta \int_{U_\alpha} f_\alpha g_\beta \omega \\ &= \sum_\alpha \sum_\beta \int_{U_\alpha \cap V_\beta} f_\alpha g_\beta \omega \end{aligned}$$

By symmetry,

$$\sum_{\beta} \int_{V_{\beta}} g_{\beta} \omega$$

is equal to the same expression, and so the two different ways of defining the integral  $\int_M \omega$  give the same value. Interchanging the summation and integration signs above is allowed due to the second property of partitions of unity, namely that for any point  $p \in M$  there is a neighbourhood where only finitely many of the partition functions  $f_{\alpha}$  do not vanish.

I think the orientations of the charts being the same is used implicitly above when writing integrals without coordinates.

### Exercise 87

Using the same charts as defined in exercise 84, we saw that  $\partial D^n$  is precisely

$$\{x \in D^n : \|x\| = 1\}$$

This is the same set of points as can be used to define  $S^{n-1}$ .

### Exercise 88

Let  $x$  be the implied local coordinate on  $M$ . Then using Stokes theorem with  $\omega = f$ , a function, gives the required result. The different signs follow from the induced orientation on the boundary points, though I don't know how to make this precise.

Perhaps by considering two maps: one from a neighbourhood of 0 to  $0 \in H^1$ , the other from 1 to  $0 \in H^1$ . In the latter case the map would be orientation-reversing. This should imply that the signs of  $f(0)$  and  $f(1)$  are different, but does not show why  $f(1)$  is the positive one.

### Exercise 89

Let

$$f(x) = e^x.$$

Then

$$\int_0^{\infty} f'(x) dx = \int_0^{\infty} e^x dx$$

is not defined, or is  $\infty$ . But if we had applied Stoke's theorem, it should be equal to

$$-f(0) = -1.$$